

QUENTIN BERGER

GIAMBATTISTA GIACOMIN

MAHA KHATIB

DISORDER AND DENATURATION TRANSITION IN THE GENERALIZED POLAND–SCHERAGA MODEL

DÉSORDRE ET TRANSITION DE
DÉNATURATION DANS LE MODÈLE DE
POLAND–SCHERAGA GÉNÉRALISÉ

ABSTRACT. — We investigate the generalized Poland–Scheraga model, which is used in the bio-physical literature to model the DNA denaturation transition, in the case where the two strands are allowed to be non-complementary (and to have different lengths). The homogeneous model was recently studied from a mathematical point of view in [BGK18, GK17], via a 2-dimensional renewal approach, with a *loop exponent* $2 + \alpha$ ($\alpha > 0$): it was found to undergo a localization/delocalization phase transition (which corresponds to the denaturation transition) of order $\nu = \min(1, \alpha)^{-1}$, together with (in general) other phase transitions. In this paper, we turn to the disordered model, and we address the question of the influence of disorder on the denaturation phase transition, that is whether adding an arbitrarily small amount of disorder (i.e. inhomogeneities) affects the critical properties of this transition. Our results are

Keywords: DNA denaturation, disordered polymer pinning model, critical behavior, disorder relevance, two-dimensional renewal processes.

2020 *Mathematics Subject Classification:* 60K35, 82D60, 92C05, 60K05, 60F10.

DOI: <https://doi.org/10.5802/ahl.34>

consistent with Harris' predictions for d -dimensional disordered systems (here $d = 2$). First, we prove that when $\alpha < 1$ (i.e. $\nu > d/2$), then disorder is *irrelevant*: the quenched and annealed critical points are equal, and the disordered denaturation phase transition is also of order $\nu = \alpha^{-1}$. On the other hand, when $\alpha > 1$, disorder is *relevant*: we prove that the quenched and annealed critical points differ. Moreover, we discuss a number of open problems, in particular the smoothing phenomenon that is expected to enter the game when disorder is relevant.

RÉSUMÉ. — Nous étudions ici le modèle de Poland–Scheraga généralisé, qui est utilisé dans la littérature bio-physique pour modéliser la transition de dénaturation de l'ADN, dans le cas où les deux brins ne sont pas forcément complémentaires (et peuvent avoir des longueurs différentes). La version homogène du modèle a été étudiée récemment d'un point de vue mathématique dans [BGK18, GK17], en utilisant un processus de renouvellement bidimensionnel, possédant un *exposant de boucle* $2 + \alpha$ ($\alpha > 0$): une transition de phase de type localisation/délocalisation y a été mise en évidence (correspondant à la transition de dénaturation), d'ordre $\nu = \min(1, \alpha)^{-1}$, ainsi que d'autres transitions de phase (en général). Dans cet article, nous nous tournons vers la version désordonnée du modèle, et nous traitons la question de l'influence du désordre sur la transition de dénaturation, c'est-à-dire de savoir si l'ajout de désordre d'intensité arbitrairement petite (i.e. d'inhomogénéités) possède un effet sur les propriétés critiques de cette transition de phase. Nos résultats sont en accord avec les prédictions de Harris pour les systèmes désordonnés d -dimensionnels (ici, $d = 2$). Tout d'abord, nous montrons que quand $\alpha < 1$ (i.e. $\nu > d/2$), alors le désordre est *non-pertinent* : les points critiques gelés et recuits sont égaux, et la transition de phase du modèle désordonné est aussi d'ordre $\nu = \alpha^{-1}$. D'autre part, quand $\alpha > 1$, le désordre est *pertinent* : nous montrons que les points critiques gelés et recuits diffèrent. De plus, nous commentons un certain nombre de problèmes ouverts, notamment concernant le phénomène de lissage de la transition de phase que l'on s'attend à observer lorsque le désordre est pertinent.

1. Introduction of the model and results

The analysis of the DNA denaturation phenomenon, i.e. the unbinding at high temperature of two strands of DNA, has led to the proposal of a very elementary model, the Poland–Scheraga (PS) model [PS70], that turns out to be relevant not only at a conceptual and qualitative level [Fis84, Gia07], but also at a quantitative level [BBB⁺99, BD98]. This model can naturally embody the inhomogeneous character of the DNA polymer, which is a monomer sequence of four different types (A, T, G and C). The binding energy for A–T pairs is different from the binding energy for G–C pairs. The quantitative analysis is then based on finite length chains with a given sequence of pairs, but in order to analyse general properties of inhomogeneous chains bio-physicists focused on the cases in which the base sequence is the realization of a sequence of random variables, that is often referred to as disorder in statistical mechanics. The PS model is limited to the case in which the two strands are of equal length and the n^{th} base of one strand can only bind with the n^{th} base of the other strand: it does not allow *mismatches* or, more generally, *asymmetric loops*, see Figure 1.1(a). A less elementary model, the generalized Poland–Scheraga model (gPS) [GK17] allows asymmetric loops, and different length strands are allowed too, see Figure 1.1(b).

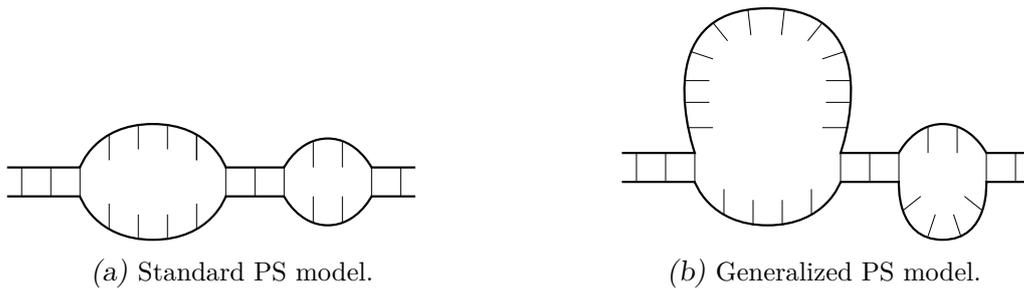


Figure 1.1. Standard v. Generalized Poland–Scheraga models. The figure on the left represents the standard PS model: the two strands of DNA have the same length (there are 14 base pairs in Figure 1.1(a)), loops are symmetric (there are 5 loops of lengths 1, 1 loop of length 3 and 1 loop of length 5). The figure on the right represents the generalized PS model: the two strands may have a different number of bases (22 for the “top” one, and 16 for the “bottom” one), and loops are allowed to be asymmetric and can be encoded by two numbers (n, m) where n is the length the “top” strand and m of the “bottom” strand (the loops in Figure 1.1(b) are from left to right $(1, 1), (1, 1), (13, 5), (1, 1), (1, 1), (3, 5), (1, 1)$).

A remarkable feature of the non disordered PS model (this corresponds to the case in which all the bases are the same: for example a strand AAA... and a second strand TTT...) is its solvable character. Notably, one can show that the model has a *denaturation* transition in the limit of infinite strand length, and one can identify the critical point (the critical temperature) and the critical behavior, i.e. the nature of the singularity of the free energy at the critical value. Somewhat surprisingly, also the gPS model is exactly solvable, in spite of the fact that it is considerably more complex than the PS model. This has been pointed out first in [GO03, GO04, NG06] and a mathematical treatment can be found in [GK17]. Let us stress that the higher complexity level of the gPS model is however reflected in a richer behavior. Notably, in the gPS model, other phase transitions exist, beyond the denaturation transition. Another relevant remark is that PS and gPS models contain a parameter (the *loop exponent*) that, in a mathematical or theoretical physics perspective, can be chosen arbitrarily and on which depends the critical behavior. In fact in this class of models the critical exponent depends on this parameter, and arbitrary critical exponents can be observed by tuning the loop exponent.

Stepping to the disordered model is not (at all) straightforward. One way to attack the problem is by looking at it as a stability issue: is the transition (and we will focus on the denaturation one) still present in the model if we introduce some disorder, for example a small amount? And, if it does, what is the new critical value and is the critical behavior the same as without disorder? We refer to [Gia07, Chapter 5] for an outline on this general very important issue in statistical mechanics and on the renormalization group ideas that lead to the so called *Harris criterion* of disorder irrelevance. We speak of disorder relevance when the disorder, irrespective of its strength, makes the critical behavior of the model different from the one of the non disordered model. Disorder is instead irrelevant if the two critical behaviors coincide for a small disorder strength. In the relevant (resp. irrelevant) case one can argue

that applying a coarse graining procedure makes the disorder stronger (resp. weaker). Harris' idea is that disorder (ir)relevance, can be read out of the critical exponent in the non disordered model.

More precisely, Harris criterion says that, if ν denotes the correlation length exponent of the non-disordered system and d the dimension, $\nu > 2/d$ implies disorder irrelevance, at least if the disorder is not too strong. One also expects disorder relevance if $\nu < 2/d$. The case $\nu = 2/d$ is dubbed marginal and deciding whether disorder is relevant or not is usually a delicate issue, even leaving aside mathematical rigor. The PS and gPS models, with their wide spectra of critical behaviors, therefore become an ideal framework for testing the validity of the physical predictions. In fact, the mathematical activity on the PS model (which is one-dimensional) has been very successful. Results include:

- Very complete understanding of the PS model when disorder is irrelevant [Ale08, Gia11, Lac10a, Ton08];
- Precise estimates on the disorder induced shift of the critical point (with respect to the annealed model) in the relevant disorder case [AZ09, DGLT09], and a proof of the fact that disorder does change the critical exponent [CdH13, GT06b] (without determining the new one: this is an open problem also in the physical literature, even if consensus is starting to emerge about the fact that pinning model in the relevant disorder regime should display a very smooth localization transition, see [BGL19, DR14] and references therein);
- Determination of whether or not there is a disordered induced critical point shift in the marginal case, and precise estimates of this shift: this issue was controversial in the physical literature [BL18, GLT10]. In absence of critical point shift, the critical exponent has also been shown to be unchanged by the noise. Showing that disorder does change the critical behavior when there is a critical point shift at marginality is an open issue, and determining the critical behavior in presence of disorder does not appear to be easier than attacking the same issue in the relevant case [DR14].

Our aim is to analyze the disordered gPS model and to understand the effect of disorder on the denaturation transition for this generalized, 2–dimensional, model.

1.1. The generalized Poland–Scheraga model

Let $\tau = \{\tau_n\}_{n \geq 0} = \{(\tau_n^{(1)}, \tau_n^{(2)})\}_{n \geq 0}$ to be a bivariate renewal process, i.e. $\tau_0 = (0, 0)$ and $\{\tau_n - \tau_{n-1}\}_{n \geq 1}$ are identically distributed \mathbb{N}^2 -valued random vectors. We denote by \mathbf{P} the law of τ , and we assume that it has inter-arrival distribution $\mathbf{P}(\tau_1 = (n, m)) = K(n + m)$, where

$$(1.1) \quad K(n) := \frac{L(n)}{n^{2+\alpha}},$$

for some $\alpha \geq 0$ and some slowly varying function $L(\cdot)$. Let $\mu := \mathbf{E}[\tau_1^{(1)}] = \mathbf{E}[\tau_1^{(2)}] \in (1, \infty]$. Without loss of generality, we assume that τ is persistent, i.e. $\sum_{n,m} K(n+m) = 1$. A set $\tau = \{\tau_n\}_{n \geq 0}$ is then interpreted as a two-strand DNA configuration: the $\tau_n^{(1)}$'s

monomer of the first strand is attached to the $\tau_n^{(2)}$'s monomer of the second strand. Put differently, the n^{th} loop in the double strand is encoded by $(\tau_n^{(1)} - \tau_{n-1}^{(1)}, \tau_n^{(2)} - \tau_{n-1}^{(2)})$, see Figure 1.1a and its caption. We refer to [GK17] for further details.

Let $\omega := \{\omega_{n,m}\}_{n,m \in \mathbb{N}}$ be a sequence of IID centered random variables (the disorder), taking values in \mathbb{R} , with law denoted \mathbb{P} . We assume that the variables $\omega_{n,m}$ are centered, have unit variance and exponential moments of all order, and we set for $\beta \in \mathbb{R}$

$$(1.2) \quad Q(\beta) := \mathbb{E}[\exp(\beta\omega)] < \infty.$$

This choice of disorder is discussed in detail in Section 1.3.

Given $\beta > 0$, $h \in \mathbb{R}$ (the pinning parameter) and $N, M \in \mathbb{N}$, we define $\mathbf{P}_{N,M}^{\beta,h,\omega}$ a measure whose Radon–Nikodym derivative w.r.t. \mathbf{P} is given by

$$(1.3) \quad \frac{d\mathbf{P}_{N,M,\omega}^{\beta,h}}{d\mathbf{P}}(\tau) := \frac{1}{Z_{N,M,\omega}^{\beta,h}} \exp\left(\sum_{n=1}^N \sum_{m=1}^M (\beta\omega_{n,m} + h)\mathbb{1}_{(n,m) \in \tau}\right) \mathbb{1}_{(N,M) \in \tau},$$

where $Z_{N,M,\omega}^{\beta,h}$ is the constrained partition function (the normalization constant)

$$(1.4) \quad Z_{N,M,\omega}^{\beta,h} := \mathbf{E} \left[\exp\left(\sum_{n=1}^N \sum_{m=1}^M (\beta\omega_{n,m} + h)\mathbb{1}_{(n,m) \in \tau}\right) \mathbb{1}_{(N,M) \in \tau} \right].$$

This corresponds to giving a reward $\beta\omega_{n,m} + h$ (or a penalty if it is negative) if the n^{th} monomer of the first strand and the m^{th} monomer of the second strand meet. Note that the presence of $\mathbb{1}_{(N,M) \in \tau}$ in the right-hand side means that we are considering trajectories that are pinned at the endpoint of the system (at a technical level it is more practical to work with the system pinned at the endpoint, see the proof of Theorem 1.1). Note that, by definition, there are at most $\min(N, M)$ renewals in the region $\{1, \dots, N\} \times \{1, \dots, M\}$.

We also define the *free* partition function, where the endpoints are free

$$(1.5) \quad Z_{N,M,\omega}^{f,\beta,h} = \mathbf{E} \left[\exp\left(\sum_{n=1}^N \sum_{m=1}^M (\beta\omega_{n,m} + h)\mathbb{1}_{(n,m) \in \tau}\right) \right],$$

that can be compared to the *constrained* partition function (1.4), see Lemma 2.2. For notational convenience, we will sometimes suppress the β, h from the partition function.

One then defines the *quenched* free energy of the system. We prove the following theorem in Section 2.

THEOREM 1.1. — *For all $\gamma > 0$, $h \in \mathbb{R}$, $\beta \geq 0$ and every choice of $\{M(N)\}_{N=1,2,\dots}$ such that $\lim_{N \rightarrow \infty} M(N)/N = \gamma$ we have*

$$(1.6) \quad \lim_{N \rightarrow \infty} \frac{1}{N} \log Z_{N,M(N),\omega}^{\beta,h} = \lim_{N \rightarrow \infty} \frac{1}{N} \mathbb{E} \log Z_{N,M(N),\omega}^{\beta,h} =: F_\gamma(\beta, h),$$

where the first limit exists $\mathbb{P}(d\omega)$ -almost surely and in $L^1(\mathbb{P})$. The same result holds for the free model, that is $F_\gamma(\beta, h) = \lim_{N \rightarrow \infty} \frac{1}{N} \log Z_{N,M(N),\omega}^{f,\beta,h}$ $\mathbb{P}(d\omega)$ -a.s. and in $L^1(\mathbb{P})$.

The function $(\beta, h) \mapsto F_\gamma(\beta, h)$ is convex, $h \mapsto F_\gamma(\beta, h)$ and $\beta \mapsto F_\gamma(\beta, h)$ are nondecreasing, and $\gamma \mapsto F_\gamma(\beta, h)$ is nondecreasing and continuous.

The homogeneous model corresponds to the case $\beta = 0$: let us drop the β and ω dependence in the partition function that will be simply denoted $Z_{N,M}^h$. The homogeneous model is *exactly solvable* and sharp estimates of $F_\gamma(0, h)$ near criticality are given in [GK17].

THEOREM 1.2 ([GK17]). — *For every $\gamma \geq 1$, when $\beta = 0$ we have $h_c(0) := \sup\{h : F_\gamma(0, h) = 0\} = 0$. Moreover there is a slowly varying function $L_{\alpha,\gamma}(\cdot)$ such that as $h \searrow 0$ one has*

$$(1.7) \quad F_\gamma(0, h) \sim L_{\alpha,\gamma}(h) h^{1/\min(1,\alpha)}.$$

Moreover, if $\sum_n n^2 K(n) < \infty$ (i.e. $\mu < \infty$), then $L_{\alpha,\gamma}(h)^{-1} = c^{-1} := \frac{1}{2} \sum_n n(n-1)K(n)$.

In fact, the disordered system also presents this transition: we define the critical point

$$(1.8) \quad h_c(\beta) := \sup\{h : F_\gamma(\beta, h) = 0\} = \min\{h : F_\gamma(\beta, h) > 0\}.$$

Let us note that $h_c(\beta)$ does not depend on $\gamma > 0$, thanks to (2.3) below.

On the other hand, we define the annealed free energy as

$$(1.9) \quad F_\gamma^a(\beta, h) := \lim_{\substack{N \rightarrow \infty, \\ M(N)/N \rightarrow \gamma}} \frac{1}{N} \log \mathbb{E} Z_{N,M(N),\omega}^{\beta,h} = F_\gamma(0, h + \log Q(\beta)).$$

This link with the homogeneous model and the fact that $h_c(0) = 0$ allow immediately to identify the annealed critical point:

$$(1.10) \quad h_c^a(\beta) := \min\{h : F_\gamma^a(\beta, h) > 0\} = -\log Q(\beta).$$

Now observe that by Jensen's inequality, we have that $\mathbb{E} \log Z_{N,M,\omega} \leq \log \mathbb{E} Z_{N,M,\omega}$ and hence $F_\gamma(\beta, h) \leq F_\gamma^a(\beta, h)$. Moreover, since $\beta \mapsto F_\gamma(\beta, h)$ is non-decreasing, we have that $F_\gamma(0, h) \leq F_\gamma^a(\beta, h)$. Therefore for every β we have

$$(1.11) \quad h_c^a(\beta) \leq h_c(\beta) \leq h_c(0).$$

One can show, by adapting the argument of proof of [Gia11, Theorem 5.2], that the second inequality is strict for every $\beta \neq 0$. The first inequality may or may not be strict and this is an important issue which is directly linked to disorder relevance and irrelevance.

Harris' criterion predicts that disorder is irrelevant if $\nu > 2/d$. Here, Theorem 1.2 suggests that $\nu = 1/\min(1, \alpha)$, if we admit that the correlation length of the non-disordered system can be given by the reciprocal of the free energy, as it is the case for the PS model, see [Gia08]. Since the model is 2-dimensional (contrary to the PS model which is 1-dimensional), it would mean that disorder is irrelevant when $\nu > 1$, that is when $\alpha < 1$.

And in fact our first result states that the first inequality in (1.11) is an equality if $\alpha < 1$ and β is not too large. For the same values of β we can also show that the critical behavior is the same as for the $\beta = 0$ case (disorder irrelevance). Our second result asserts that the inequality is strict for $\alpha > 1$. We interpret this critical point shift, with a certain abuse, as disorder relevance. We however refer to the discussion in Section 1.3 (in particular Conjecture 1.5) regarding the change in the

critical behavior. We therefore prove that disorder is irrelevant if $\alpha < 1$, and relevant (in terms of critical points) if $\alpha > 1$, confirming Harris’ prediction.

1.2. Relevance and irrelevance of disorder

Let us define $\sigma := \tau \cap \tau'$, where τ and τ' are two independent copies of τ : σ is another bivariate renewal process, and Proposition A.3 tells that σ is terminating if $\alpha \in (0, 1)$ and persistent if $\alpha > 1$ (the case $\alpha = 1$ is discussed in Remark A.7).

THEOREM 1.3. — *Assume that σ is terminating (this includes $\alpha < 1$ and excludes $\alpha > 1$). Then there exists $\beta_1 > 0$ (see (3.5)), such that for every $\beta \in (0, \beta_1)$ we have $h_c(\beta) = h_c^a(\beta)$, and moreover*

$$(1.12) \quad \lim_{h \searrow h_c(\beta)} \frac{\log F_\gamma(\beta, h)}{\log(h - h_c(\beta))} = \frac{1}{\alpha}.$$

Hence, the order of the phase transition is unchanged when σ is terminating (which is the case if $\alpha < 1$), at least when β is small enough. We prove Theorem 1.3 in Section 3. We mention that when the disorder distribution is infinitely divisible (for instance Gaussian), one can get sharper bounds regarding the critical behavior of $F_\gamma(\beta, h)$, via a replica-coupling method, as done in [Ton08] or [Wat12]. For a statement and a detailed proof, we refer to [Kha16].

On the other hand, when $\alpha > 1$, we show that the quenched and annealed critical points differ, and we give a lower bound on the critical point shift.

THEOREM 1.4. — *For $\alpha > 1$ we have $h_c(\beta) > h_c^a(\beta)$ for every $\beta > 0$. Moreover, for every $\varepsilon > 0$, there exists $\beta_\varepsilon > 0$ such that for any $\beta \leq \beta_\varepsilon$ we have*

$$(1.13) \quad h_c(\beta) - h_c^a(\beta) \geq \Delta_\beta^\varepsilon := \begin{cases} \beta^{\frac{2\alpha}{\alpha-1} + \varepsilon} & \text{if } \alpha \in (1, 2], \\ \beta^4 |\log \beta|^{-6} & \text{if } \alpha > 2. \end{cases}$$

Moreover, there is a slowly varying function $\tilde{L}(\cdot)$ such that

$$(1.14) \quad h_c(\beta) - h_c^a(\beta) \leq \tilde{L}(1/\beta) \beta^{\frac{2\alpha}{\alpha-1} \vee 4}.$$

We add that $\beta \mapsto h_c(\beta) - h_c^a(\beta)$ is a non decreasing function of β : this result can be proven by the exact same procedure as the one used to prove in [GLT11, Proposition 6.1]. It is to be interpreted that disorder relevance is *non-decreasing* in β .

1.3. On the results, perspectives and related work

1.3.1. On the main theorems

A two replica computation plays a central role in the proof of Theorem 1.3 and in the proof of (1.14) of Theorem 1.4: the intersection renewal σ therefore emerges naturally, like in the PS model. In the PS context, we now know that disorder is

irrelevant (for small values of β) if and only if the intersection renewal is terminating [BL18]. For the gPS model our results go in the same direction, but it is not sharp in the marginal case $\alpha = 1$: we only show disorder irrelevance when the intersection renewal σ is terminating. We refer to Remark A.7 for further discussion on the case $\alpha = 1$, where more technicalities arise.

The proof of (1.13) is based on coarse graining techniques and fractional moment method: we have chosen to adapt the method proposed in [DGLT09] and the difficulties in its generalization come from dealing with the richness of a multidimensional path with respect to the one dimensional structure of the PS models. A keyword for these difficulties is *off-diagonal estimates*. It can certainly be improved in the direction of getting rid of the ε in the exponent for $\alpha \in (1, 2]$ and of the logarithmic term in the case $\alpha > 1$ by using more sophisticated coarse graining techniques (see [Gia11, Chapter 6] and references therein). One could probably aim also for sharp estimates, like in [BL18], but the estimates are technically rather demanding already to obtain (1.13). We have chosen to stick to these simplified non-optimal (but almost optimal) bounds because sharper results would have required a substantially heavier argument of proof. The techniques developed in [BL17, BL18] should transfer to this model: at the expense of a high level of technicality, we expect that, in analogy with the PS model, the necessary and sufficient condition for a critical point shift is the persistence of the intersection renewal $\sigma = \tau \cap \tau'$.

1.3.2. Discussion on the presence of a *smoothing* phenomenon

Of course, a fully satisfactory result on disorder relevance would include showing that the critical exponent is modified by the disorder. We do not have such a result, but let us make one observation and formulate a conjecture.

The observation is that Theorem 1.3 may appear at first surprising in view of the smoothing Inequality [CdH13, GT06b] for PS models that ensures that the free energy exponent cannot be smaller than 2 in presence of disorder: for the gPS model the free energy exponent can go down to 1, since in (1.12) we can choose α arbitrarily close to 1. The reason of the difference is that the PS model is 1-dimensional whereas the gPS model is 2-dimensional: Harris criterion tells that disorder should be irrelevant if $\nu > 2$ for the PS model, and $\nu > 1$ for the gPS model. In the gPS model, the irrelevant disorder regime therefore holds even if ν ($= \min(1, \alpha)^{-1}$) is arbitrarily close to 1: hence one should not hope for a general smoothing inequality valid whatever α is.

It is however worthwhile attempting to sketch the argument in [GT06b], in the simplified set-up of Gaussian charges [Gia07, Chapter 5, Section 4]. This is useful both to understand where the argument fails and because we can realize that a suitable generalization of the argument naturally leads to a conjecture that we state just below.

The argument [GT06b] is based on introducing a coarse graining scale $\ell \in \mathbb{N}$ and considering the environment in terms of ℓ -boxes, see Figure 1.2. We argue for the case $\gamma = 1$ ($M = N$) and we consider the system at criticality, that is $h = h_c(\beta)$:

- (1) A good ℓ -box is a box for which the pinned partition function (i.e. the renewal is pinned at the south-west and north-east corners of the ℓ -box)

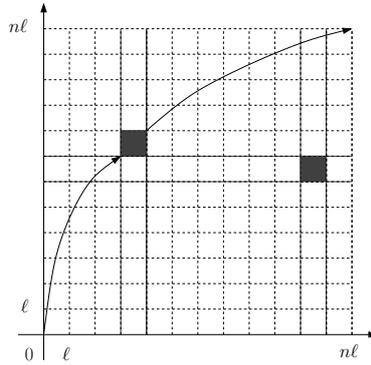


Figure 1.2. Schematic view of the coarse graining procedure proposed for a smoothing inequality. The environment is divided in blocks of size ℓ , called ℓ -boxes. A ℓ -box is good (shadowed in the figure) if the partition function in this block grows at an exponential rate that is larger than the free energy of the system. The good ℓ -boxes will be rare, but we can choose n such that in a system of linear size $n\ell$, with positive probability, there will be at least one good ℓ -box. A lower bound on the partition function follows by the limitation to trajectories that visit only a given good ℓ -box (say, the closest).

is larger than $\exp(\frac{1}{2}\ell F_1(\beta, h + \delta))$, with $\delta > 0$. For $\ell \rightarrow \infty$ this is a rare event. The probability of such an event can be estimated from below by shifting the environment of δ/β , that is $\omega_{i,j}$ is replaced with $\omega_{i,j} + \delta/\beta$, and by performing a relative entropy estimate [Gia11, Chapter 5]. This shows that the probability of such a rare event is at least $\exp(-\delta^2\ell^2/(2\beta^2))$: note the ℓ^2 term, with respect to ℓ in the PS case [GT06b].

- (2) We then make a lower bound on the partition function of the system by discarding renewal trajectories that visit ℓ -boxes that are not good, and keeping only trajectories that enter good ℓ -boxes through the south-west corner and exit through the north-east corner.

The trajectories are therefore alternated jumps to a good box, *visit* of the box, and then a new jump to another good box. Jumps are long because good boxes are rare. The analysis in [GT06b] is ultimately reduced to see what happens in one *jump and visit*: by exploiting super-additivity one can even just choose $N = n\ell$ such that there is (say, with probability at least 1/2), at least one good box in the system (like it is done in [BL11]). We therefore see that we need $n^2 \exp(-\delta^2\ell^2/(2\beta^2)) \approx 1$, so that $n \approx \exp(-\delta^2\ell^2/\beta^2)$: with this level of precision, jumping to enter such a box costs $K(n\ell) = (n\ell)^{-(2+\alpha)}$ (let us consider the case in which $L(\cdot)$ is a constant, but the computation goes through in the same way also in the general case). In the box there will be a contribution $\exp(\frac{1}{2}\ell F_1(\beta, h + \delta))$. The net contribution to the logarithm of the partition function, divided by the size $n\ell$ of the system, is then

$$(1.15) \quad \frac{1}{n\ell} \left(\log K \left(\exp(-\delta^2\ell^2/\beta^2)\ell \right) + \frac{\ell}{2} (F_1(\beta, h + \delta)) \right) \geq \frac{1}{n} \left(-\frac{c}{\beta^2} \delta^2\ell + \frac{1}{2} F_1(\beta, h + \delta) \right),$$

with c a positive constant that we have left implicit (it depends on more accurate computations, and can be in principle just reduced to $2 + \alpha$).

Now let us choose $h = h_c(\beta)$. So the argument we just outlined goes in the direction of saying that

$$(1.16) \quad 0 = F_1(\beta, h_c(\beta)) \geq \frac{1}{n} \left(-\frac{c}{\beta^2} \delta^2 \ell + \frac{1}{2} F_1(\beta, h_c(\beta) + \delta) \right),$$

so that

$$(1.17) \quad F_1(\beta, h_c(\beta) + \delta) \leq \frac{2c}{\beta^2} \delta^2 \ell.$$

At this stage choosing ℓ arbitrarily large is of no help. The steps we have performed up to now require that $\ell\delta$ is large (so that the good boxes we have chosen are really sparse). On the other hand we need to have chosen the size of the boxes so that $Z_{\ell, \ell, \omega}^{\beta, h} \geq \exp(\ell(F_1(\beta, h_c(\beta) + \delta))/2)$. This is a delicate issue, but it definitely appears that for this to hold, $\ell F_1(\beta, h_c(\beta) + \delta)$ needs to be sufficiently large (say, larger than a suitable constant): see for example the discussion on the notion of correlation length given in [Gia11, Chapter 2] and references therein, notably [GT06a], where the correlation length is identified by the reciprocal of the free energy. But if ℓ is (a constant times) $1/F_1(\beta, h_c(\beta) + \delta)$ then from (1.17) we obtain

$$(1.18) \quad F_1(\beta, h_c(\beta) + \delta) \leq C\delta,$$

for some $C > 0$. But such a bound is trivial: it holds with $C = 1$ just because the contact density cannot exceed one! On the other hand, as we have already pointed out, we could not have hoped for a better bound valid for any $\alpha > 0$.

In spite of the fact that it leads to a trivial result, we insist that the argument we have just outlined can be made rigorous: the delicate step is the last one, where one has to use arguments developed in [GT06a]. It can therefore be taken as a starting point to push things further. Indeed, it appears useless to modify the environment in the whole ℓ -box, at least if $\alpha > 1$. In fact if $\alpha > 1$ one can show that for $q > 1/\max(\alpha, 2)$

$$(1.19) \quad \lim_{N \rightarrow \infty} \mathbf{P} \left(\tau \cap [0, N]^2 \subset \{(i, j) \in \mathbb{Z}^2 : |i - j| \leq N^q\} \right) = 1.$$

We can then consider modifying only the environment that is close to the diagonal, that is in a subset of the ℓ -box with $|i - j| \leq \ell^q$. This would improve the lower bound on the probability of a good ℓ -box to $\exp(-c' \delta^2 \ell^{q+1})$, and (1.17) would become

$$F_1(\beta, h_c(\beta) + \delta) \leq \frac{c''}{\beta^2} \delta^2 \ell^q.$$

Taking ℓ a constant times $1/F_1(\beta, h_c(\beta) + \delta)$ as in the argument leading to (1.18), and then taking q arbitrarily close to $1/\min(\alpha, 2)$ supports the following:

CONJECTURE 1.5. — *For every $\alpha > 0$ and every $\beta > 0$*

$$(1.20) \quad \limsup_{\delta \searrow 0} \frac{\log F_1(\beta, h_c(\beta) + \delta)}{\log \delta} \geq \begin{cases} \frac{2\alpha}{\alpha+1} & \text{for } \alpha \in (1, 2), \\ \frac{4}{3} & \text{for } \alpha \geq 2. \end{cases}$$

We stress that a natural concern arises from performing the change of measure only in a subset of the environment, close to the diagonal. One indeed needs to be sure that the trajectories contributing to (a fraction of) $\frac{1}{\ell} \log Z_{\ell, \ell, \omega} \approx F_1(\beta, h_c(\beta) + \delta)$ can be constrained to stay in the region $\{(i, j) \in \mathbb{Z}^2 : |i - j| \leq \ell^q\}$: if it is the case, one can “force” trajectories to visit sites where the environment has indeed been shifted.

1.3.3. An important modeling issue: the choice of the disorder

There is no doubt that the first disorder that comes to mind when thinking of DNA modeling is not the one we have used. One would rather choose $\omega_{i,j} = f(\omega_i, \omega_j)$ for a suitable choice of a function f and a sequence $\{\omega_j\}_{j=1,2,\dots}$ of random variables (let us say IID for simplicity, but if we want to stick to DNA problems very closely it appears that some sort of strongly correlated sequence may be more appropriate [PBG⁺92]). For example, we could choose ω_j taking only two values e_{AT} and e_{GC} and then make a choice for f that reflects the fact that AT bounds are weaker than GC bounds, and that all other possible bounds are even weaker. Even restricting to $\{\omega_j\}_{j=1,2,\dots}$ that is IID, this model is highly non trivial (gPS model with this type of disorder has been considered at a numerical level in [GO03, GO04], see also [EON11, TN08] for related work). But one could also choose to consider the binding of two sequences that are not complementary (the case considered in [NG06] goes in this direction, even if only heuristics and numerics are presented): choose for example two independent sequences $\{\omega_j^{(1)}\}_{j=1,2,\dots}$ and $\{\omega_j^{(2)}\}_{j=1,2,\dots}$ and use $\omega_{i,j} = f(\omega_i^{(1)}, \omega_j^{(2)})$. This is somewhat closer to what we are using (though it can be considered as a one-dimensional disorder), but it is still very difficult to deal with. The problem is in any case due to correlations in the disorder field $\omega_{i,j}$, which can be dealt with in some cases, see e.g. [BL12, BP15] or [AB18, CCP19]. Our choice is in a sense a toy choice, but we stress that it is conceptually similar to the simplification made for example in [BH02] in the RNA context. Moreover it recovers importance once we leave somewhat the DNA context and focus rather on moving toward understanding mathematically Harris’ theory of disorder (ir)relevance—in particular for 2–dimensional systems, compared to the PS model, which is 1-dimensional.

We also point out that this disordered version of the gPS model gives a bridge between pinning model and directed polymers in random environment [Com07, Lac10b], in particular, to the long range directed polymer [Com07, Wei16]. Moreover a different class of two-strand polymer problems (the random walk pinning model) is treated in [BS10, BS11, BT10].

1.3.4. Open questions and perspectives

Several natural issues remain open: let us list some of them.

- (1) Prove a smoothing inequality, thus showing disorder relevance in the original sense of Harris, for $\alpha > 1$ (see Conjecture 1.5).
- (2) What is the effect of disorder on the other phase transitions? Here we have addressed only the denaturation transition, but in [GK17] other transitions

are shown to exist. Do they withstand the introduction of disorder? If so, does the corresponding critical behavior differ from the homogeneous case? This is the question (quickly) addressed [NG06] where a rather bold conjecture is set forth.

- (3) We have dealt only with free energy estimates, but, like for the standard PS model, obtaining precise estimates on the gPS process (i.e. establish properties of trajectories) is very challenging, see [Gia11, Chapter 8] and references therein. The problem comes of course from the inhomogeneous nature of the disorder and the fact that on rare regions atypical disorder behaviors appear (this is ultimately also the problem we face at the free energy level, but it becomes particularly explicit when one analyses the trajectories). A precise analysis of the trajectories of the non disordered gPS model can be found in [BGK18]: this analysis is substantially more demanding than the corresponding one for the PS model.
- (4) Dealing with the marginal case $\alpha = 1$ is open, mostly because of the additional technical difficulties (more complicated coarse-graining procedure, more technical estimates for bivariate renewals, etc.). This appears to be a problem at reach, but a very substantial amount of technical work is certainly needed.

1.3.5. Organization of the rest of the work

The issues of existence and self-averaging of the free energy, i.e. the proof of Theorem 1.1, are treated in Section 2. In Section 3 we prove Theorem 1.3, as well as the upper bound (1.14) of Theorem 1.4. The rest of the Theorem 1.4 is proven in Section 4. We collect in Appendix A a number of statements and proofs about bivariate renewals.

1.4. Some further notations

We stress that τ is symmetric and in the domain of attraction of a $\min(\alpha, 2)$ -stable distribution: we denote $(b_n)_{n \geq 1}$ be the recentering sequence and $(a_n)_{n \geq 0}$ the renormalizing sequence for τ_n , that is such that $\frac{1}{a_n}(\tau_n - (b_n, b_n))$ converges to a $\min(\alpha, 2)$ stable distribution, whose density is denoted $g_\alpha(\cdot, \cdot)$. For b_n , we have $b_n = \mu n$ if $\alpha > 1$, $b_n = n\mathbf{E}[\min(X_1, n)]$ if $\alpha = 1$, and $b_n = 0$ if $\alpha \in (0, 1)$. The asymptotic behavior of a_n is characterized by

$$(1.21) \quad \begin{cases} L(a_n)(a_n)^{-\alpha} \sim 1/n & \text{if } \alpha < 2, \\ \sigma(a_n)(a_n)^{-2} \sim 1/n & \text{if } \alpha \geq 2, \end{cases}$$

where $\sigma(n) := \mathbf{E}[\min(X_1, n)^2]$. If $\alpha = 2$ and $\mathbf{E}[X^2] = +\infty$, then $\sigma(n)$ grows to infinity as a slowly varying function (and verifies $\sigma(n)/L(n) \rightarrow +\infty$), whereas if $\mathbf{E}[X^2] < +\infty$ (in particular when $\alpha > 2$) a_n is proportional to \sqrt{n} .

In any case, there exists some slowly varying function $\psi(\cdot)$ such that

$$(1.22) \quad a_n = \psi(n)n^{1/\min(\alpha, 2)}.$$

We provide some useful results on bivariate renewals in Appendix A, in particular on the renewal mass function $\mathbf{P}((n, m) \in \tau)$.

2. Free Energy: existence and properties

In this section we often assume $\gamma \in \mathbb{Q}$: in this case we write it as $\gamma = p/q$ with p and q relatively prime positive integer numbers.

PROPOSITION 2.1. — *For every $\gamma > 0$ and every $\{M(N)\}_{N=1,2,\dots}$ such that $\lim_{N \rightarrow \infty} M(N)/N = \gamma$ we have that*

$$(2.1) \quad \lim_{N \rightarrow \infty} \frac{1}{N} \log Z_{N,M(N),\omega}^{\beta,h} = \lim_{N \rightarrow \infty} \frac{1}{N} \mathbb{E} \log Z_{N,M(N),\omega}^{\beta,h} =: F_\gamma(\beta, h),$$

where the first limit is meant $\mathbb{P}(d\omega)$ -a.s. and in $L^1(\mathbb{P})$. $F_\gamma(\cdot, \cdot)$ is convex and $F_\gamma(\beta, \cdot)$ is non-decreasing, and also $F_\gamma(\cdot, h)$ is non-decreasing on the positive semi-axis, non-increasing in the negative one. Moreover if $\gamma = \frac{p}{q} \in \mathbb{Q}$

$$(2.2) \quad F_\gamma(\beta, h) = \sup_{N: \frac{N}{q} \in \mathbb{N}} \frac{1}{N} \mathbb{E} \log Z_{N,\gamma N,\omega}^{\beta,h}.$$

Finally we have the bound: for every $\gamma_2 \geq \gamma_1 > 0$

$$(2.3) \quad F_{\gamma_1}(\beta, h) \leq F_{\gamma_2}(\beta, h) \leq \frac{\gamma_2}{\gamma_1} F_{\gamma_1}(\beta, h),$$

which implies that $\gamma \mapsto F_\gamma(\beta, h)$ is locally Lipschitz (hence continuous).

Proof. — The proof is divided into several steps:

- (1) We first show that for $\gamma \in \mathbb{Q}$, along a subsequence with $N/q \in \mathbb{N}$, $\log Z_{N,\gamma N,\omega}$ is super-additive in an ergodic sense, which implies the existence of the free energy limit (2.1) along this subsequence.
- (2) The restriction $\gamma N \in \mathbb{N}$ is then removed by a direct estimate, for what concerns the existence of the free energy limit, still with $\gamma \in \mathbb{Q}$.
- (3) We then prove a comparison estimate between $Z_{N,\gamma_1 N,\omega}$ and $Z_{N,\gamma_2 N,\omega}$ and use it to establish the existence of the free energy limit for $Z_{N,\gamma N,\omega}$, every $\gamma > 0$.
- (4) The same comparison estimate yields also (2.3), and the fact that one can take the limit along an arbitrary sequence satisfying $M(N) \sim \gamma N$, for $N \rightarrow \infty$.
- (5) Finally, we prove the convexity and monotonicity statements.

Step (1). — With $\gamma = p/q$ set $\mathcal{Z}_j(\omega) := Z_{jq,jp,\omega}$. Then one directly sees that

$$(2.4) \quad \mathcal{Z}_{j_1+j_2}(\omega) \geq \mathcal{Z}_{j_1}(\omega) \mathcal{Z}_{j_2}(\Theta_{j_1 q, j_1 p} \omega),$$

where $(\Theta_{q,p}\omega)_{n,m} = \omega_{q+n,p+m}$. Since ω is an IID sequence of L^1 random variables, it is straightforward to see that $\log \mathcal{Z}_j \in L^1(\mathbb{P})$. Also,

$$|\log \mathcal{Z}_j| \leq hn + \beta \sup_{\gamma \in \Gamma} \sum_{(n,m) \in \gamma} |\omega_{n,m}|,$$

where Γ is the set of nearest-neighbors up-right paths: a Last Passage Percolation observation then tells us that a sufficient condition for having $\sup_n \frac{1}{n} \mathbb{E} |\log \mathcal{Z}_j| < +\infty$ is $\mathbb{E}[\omega_{1,1}^2] < +\infty$, see [Mar02]. Hence we see that $\{-\log \mathcal{Z}_j(\omega)\}_{j=1,2,\dots}$ satisfies

the hypotheses of Kingman sub-additive ergodic Theorem (see for example [Gia07, Section A.7]), and we get that $\{\frac{1}{j} \log \mathcal{Z}_j(\omega)\}_{j=1,2,\dots}$ converges $\mathbb{P}(d\omega)$ -a.s. and in $L^1(\mathbb{P})$. Moreover (2.4) directly tells us that $\{\mathbb{E} \log \mathcal{Z}_j(\omega)\}_{j=1,2,\dots}$ is super-additive, so that $\lim_{j \rightarrow \infty} \frac{1}{j} \mathbb{E} \log \mathcal{Z}_j(\omega) = \sup_{j \in \mathbb{N}} \frac{1}{j} \mathbb{E} \log \mathcal{Z}_j(\omega)$. This establishes (2.2), and also (2.1), but only for $M(N) = \gamma N$ with $\gamma \in \mathbb{Q}$ and along the subsequence satisfying $\gamma N \in \mathbb{N}$.

Step (2). — Still with $\gamma = p/q$, the restriction to $\gamma N \in \mathbb{N}$ can be removed by observing that we can write $N = jq + r$, with $r \in \{0, 1, \dots, q - 1\}$, and for $r \neq 0$

$$(2.5) \quad \begin{aligned} Z_{N, \lfloor \gamma N \rfloor, \omega} &\geq Z_{jq, jp, \omega} \exp(\beta \omega_{N, \lfloor \gamma N \rfloor} + h) K\left(r + \left\lfloor jp + \frac{p}{q} r \right\rfloor - jp\right) \\ &\geq c(p, q) \exp(\beta \omega_{N, \lfloor \gamma N \rfloor} + h) Z_{jq, jp, \omega}, \end{aligned}$$

where $c(p, q) > 0$.

In the same way

$$(2.6) \quad \begin{aligned} Z_{(j+1)q, (j+1)p, \omega} &\geq Z_{N, \lfloor \gamma N \rfloor, \omega} \exp(\beta \omega_{(j+1)q, (j+1)p} + h) K\left(q - r + (j + 1)p - \left\lfloor jp + \frac{p}{q} r \right\rfloor\right) \\ &\geq c(p, q) \exp(\beta \omega_{(j+1)q, (j+1)p} + h) Z_{N, \lfloor \gamma N \rfloor, \omega}, \end{aligned}$$

possibly redefining $c(p, q) > 0$. From (2.5) and (2.6) one easily removes the restriction to $\gamma N \in \mathbb{N}$ and establishes (2.1) for $M(N) = \lfloor \gamma N \rfloor$ with $\gamma \in \mathbb{Q}$.

Step (3). — We now establish (2.1) for $M(N) = \lfloor \gamma N \rfloor$ for an arbitrary $\gamma > 0$, by proving the announced comparison bounds, upper and lower.

The upper bound is more general: if $M_2 > M_1$ and if there exists $c > 0$ such that $M_2 \leq cN$ we see that

$$(2.7) \quad \begin{aligned} Z_{N, M_1, \omega} &= \sum_{n=0}^{N-1} \sum_{m=0}^{M_1-1} Z_{n, m, \omega} K(N - n + M_1 - m) \exp(\beta \omega_{N, M_1} + h) \\ &\leq c_K N^{c_K} \exp(\beta(\omega_{N, M_1} - \omega_{N, M_2})) \\ &\quad \times \sum_{n=0}^{N-1} \sum_{m=0}^{M_1-1} Z_{n, m, \omega} K(N - n + M_2 - m) \exp(\beta \omega_{N, M_2} + h) \\ &\leq c_K N^{c_K} \exp(\beta(\omega_{N, M_1} - \omega_{N, M_2})) Z_{N, M_2, \omega}, \end{aligned}$$

where in the first inequality we have used that $K(\cdot)$ is regularly varying and that $M_2 \leq cN$ to see that there exists $c_K > 0$ such that

$$(2.8) \quad \frac{K(N - n + M_1 - m)}{K(N - n + M_2 - m)} \leq c_K N^{c_K},$$

for every N . For the second inequality we have relaxed the constrained $m < M_1$ to $m < M_2$.

On the other hand, we prove a comparison lower bound only for M of the form $\lfloor \gamma N \rfloor$. Let us choose $\gamma_2 > \gamma_1 > 0$. Note that for

$$(2.9) \quad N' := \left\lfloor \frac{\gamma_1}{\gamma_2} N \right\rfloor - \left\lfloor \frac{2}{\gamma_2} \right\rfloor,$$

we have $\lfloor \gamma_2 N' \rfloor + 1 \leq \lfloor \gamma_1 N \rfloor$ so that

$$(2.10) \quad \begin{aligned} Z_{N, \lfloor \gamma_1 N \rfloor, \omega} &\geq K (N - N' + \lfloor \gamma_1 N \rfloor - \lfloor \gamma_2 N' \rfloor) \exp(\beta \omega_{N, \lfloor \gamma_1 N \rfloor} + h) Z_{N', \lfloor \gamma_2 N' \rfloor, \omega} \\ &\geq (c_K N^{c_K})^{-1} \exp(\beta \omega_{N, \lfloor \gamma_1 N \rfloor} + h) Z_{N', \lfloor \gamma_2 N' \rfloor, \omega}, \end{aligned}$$

possibly changing the value of $c_K > 0$.

We now choose $0 < \gamma_1 < \gamma_2 \in \mathbb{Q}$. Then (2.7) implies that $\mathbb{P}(d\omega)$ -a.s.

$$(2.11) \quad \limsup_{N \rightarrow \infty} \frac{1}{N} \log Z_{N, \lfloor \gamma_1 N \rfloor, \omega} \leq F_{\gamma_2}(\beta, h),$$

and (2.10) implies that $\mathbb{P}(d\omega)$ -a.s.

$$(2.12) \quad \liminf_{N \rightarrow \infty} \frac{1}{N} \log Z_{N, \lfloor \gamma_1 N \rfloor, \omega} \geq \lim_{N \rightarrow \infty} \frac{1}{N'} \log Z_{N', \lfloor \gamma_2 N' \rfloor, \omega} = \frac{\gamma_1}{\gamma_2} F_{\gamma_2}(\beta, h),$$

and the proof of (2.1) is achieved in the $\mathbb{P}(d\omega)$ -a.s. sense for $M(N) = \lfloor \gamma_1 N \rfloor$, by choosing a sequence of values for γ_2 converging to γ_1 , defining thus $F_{\gamma_1}(\beta, h)$ also by this limit procedure. Note that a byproduct is that (2.3) holds, hence $\gamma \mapsto F_{\gamma}(\beta, h)$ is non decreasing and (locally) Lipschitz continuous. To upgrade (2.1) to the $L^1(\mathbb{P})$ sense one simply applies the expectation $\mathbb{E}[\cdot]$ to the log of (2.10) and (2.7) so that one obtains $\lim_{N \rightarrow \infty} (1/N) \mathbb{E} \log Z_{N, \lfloor \gamma N \rfloor, \omega} = F_{\gamma}(\beta, h)$ for every $\gamma > 0$, and the first limit in (2.1) holds in the $L^1(\mathbb{P})$ sense by Scheffé’s Lemma.

Step (4). — The generalization to a sequence $M(N) \sim \gamma N$ is just made by observing that given arbitrary $\gamma_1 < \gamma_2$ with $\gamma \in (\gamma_1, \gamma_2)$ for N_0 sufficiently large we have $\lfloor \gamma_1 N \rfloor < M(N) < \lfloor \gamma_2 N \rfloor$ for every $N \geq N_0$. At this point we can apply the comparison bounds like in the previous step and conclude by an approximation procedure.

Step (5). — The function $(\beta, h) \mapsto F_{\gamma}(\beta, h)$ is convex because it is the limit of a sequence of convex functions. Monotonicity in h for β fixed is also evident from the finite N expression. The fact that $\beta \mapsto F_{\gamma}(\beta, h)$ is non increasing for $\beta \leq 0$ and non decreasing for $\beta \geq 0$ follows from convexity and the fact that $\partial_{\beta} \mathbb{E} \log Z_{N, M, \omega}^{\beta, h} = 0$ (by direct computation, since the ω variables are centered), so $\partial_{\beta} F_{\gamma}(\beta, h)|_{\beta=0} = 0$. This completes the proof of Proposition 2.1. \square

We now compare the constrained and the free partition function:

LEMMA 2.2. — For any $\alpha_+ > \alpha$, there exists C such that for every $N, M \in \mathbb{N}$ and

$$(2.13) \quad \begin{aligned} Z_{N, M, \omega}^c &\leq Z_{N, M, \omega}^f \\ &\leq Z_{N, M, \omega}^c \times \left(1 + C(N + M)^{3+\alpha_+} e^{-\beta \omega_{N, M}} \sup_{\substack{1 \leq n \leq N \\ 1 \leq m \leq M}} \{e^{\beta \omega_{n, M}}, e^{\beta \omega_{N, m}}\} \right). \end{aligned}$$

Proof. — The lower bound is trivial: we have $Z_{N, M, \omega}^f \geq Z_{N, M, \omega}^f((N, M) \in \tau) = Z_{N, M, \omega}^c$, where we introduced the notation

$$Z_{N, M, \omega}^f(A) = \mathbf{E} \left[\exp \left(\sum_{n=1}^N \sum_{m=1}^M (\beta \omega_{n, m} + h) \mathbf{1}_{(n, m) \in \tau} \right) \mathbf{1}_A \right].$$

On the other hand, for $N, M \geq 1$, we have

$$(2.14) \quad Z_{N,M,\omega}^f = \sum_{n=0}^N \sum_{m=0}^M Z_{N,M,\omega}^f \left(\tau \cap [n, N] \times [m, M] = \{(n, m)\} \right) \\ \leq Z_{N,M,\omega}^c + \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} Z_{n,m,\omega}^c + \sum_{n=0}^{N-1} Z_{n,M,\omega}^c + \sum_{m=0}^{M-1} Z_{N,m,\omega}^c.$$

Now, observe that for any $n \leq N-1$, $m \leq M-1$,

$$(2.15) \quad Z_{n,m,\omega}^c \leq C_1(N+M)^{2+\alpha_+} e^{-h-\beta\omega_{N,M}} Z_{N,M,\omega}^c K(M+N-n-m) e^{h+\beta\omega_{N,M}}$$

for any $\alpha_+ > \alpha$, so that

$$(2.16) \quad \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} Z_{n,m,\omega}^c \leq C_1(N+M)^{2+\alpha_+} Z_{N,M,\omega}^c e^{-h-\beta\omega_{N,M}}.$$

For $n < N$ and $m = M$, there exists C_2 such that

$$(2.17) \quad Z_{n,M,\omega}^c \leq C_2 N^{2+\alpha_+} Z_{N,M,\omega}^c \exp(\beta\omega_{n,M} - \beta\omega_{N,M}),$$

and we obtain

$$(2.18) \quad \sum_{n=0}^{N-1} Z_{n,M,\omega}^c \leq C_2 N^{3+\alpha_+} \sup_{1 \leq n \leq N} \{\exp(\beta\omega_{n,M})\} e^{-\beta\omega_{N,M}} Z_{N,M,\omega}^c.$$

The analogous holds for the last term in (2.14), and the proof is therefore complete. \square

From Lemma 2.2, and using also that $\lim_{N \rightarrow \infty} \frac{1}{N} \sup_{1 \leq n \leq N} \omega_{n,M} = 0$ \mathbb{P} -a.s. (note that $\mathbb{P}(|\omega_1| > x) = o(1/x)$, since $\mathbb{E}[|\omega_1|] < +\infty$), it follows that Theorem 1.1 also holds for the free model, namely:

$$(2.19) \quad F_\gamma(\beta, h) = \lim_{N \rightarrow \infty} \frac{1}{N} \log Z_{N,M(N),\omega}^{f,\beta,h} \quad \mathbb{P}(d\omega)\text{-a.s. and in } L^1(\mathbb{P}).$$

We now introduce some notation that is used later in the paper: for positive integers $a_1 < a_2$ and $b_1 < b_2$, we define the partition function of the system on $[a_1, a_2] \times [b_1, b_2]$ by

$$(2.20) \quad Z_{(a_1,b_1),(a_2,b_2),\omega} \\ := \mathbf{E} \left[\exp \left(\sum_{n=a_1+1}^{a_2} \sum_{m=b_1+1}^{b_2} (\beta\omega_{n,m} + h) \mathbf{1}_{(n,m) \in \tau} \right) \mathbf{1}_{(a_2,b_2) \in \tau} \middle| (a_1, b_1) \in \tau \right],$$

with the convention that $Z_{(a_1,b_1),(a_1,b_1),\omega} = 1$ and $Z_{(a_1,b_1),(a_1,b_2),\omega} = Z_{(a_1,b_1),(a_2,b_1),\omega} = 0$.

3. Upper bound on the critical point shift

The arguments in this section follow the line of proof of H. Lacoin in [Lac10a], and is mainly based on a second moment computation. We start with some preliminary results.

PROPOSITION 3.1. — *If $\{Z_{N,M(N),\omega}^{f,\beta,h_c^a(\beta)}\}_N$ is uniformly integrable, there exists $\zeta > 0$ such that for every sequence of events $\{A_N\}_{N=1,2,\dots}$ satisfying $\lim_N \mathbf{P}(A_N) = 0$ there is $N_0 \in \mathbb{N}$ such that*

$$(3.1) \quad \inf_{N \geq N_0} \mathbb{P} \left(\mathbf{P}_{N,M(N),\omega}^{f,\beta,h_c^a(\beta)}(A_N) \leq \frac{1}{2} \text{ and } Z_{N,M(N),\omega}^{f,\beta,h_c^a(\beta)} > \frac{1}{2} \right) \geq \zeta.$$

Proof. — We set $h = h_c^a(\beta)$. It is sufficient to prove that there exists $\zeta > 0$ such that

$$(3.2) \quad \inf_N \mathbb{P} \left(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} > \frac{1}{2} \right) \geq 2\zeta.$$

and

$$(3.3) \quad \lim_{N \rightarrow \infty} \mathbb{P} \left(\mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N) > \frac{1}{2} \text{ and } Z_{N,M(N),\omega}^{f,\beta,h_c^a(\beta)} > \frac{1}{2} \right) = 0.$$

Since $\mathbb{E}[Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)}] = 1$, and because $\{Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)}\}_N$ is uniformly integrable, then (3.2) follows immediately from [Gia11, Lemma 4.6].

For (3.3) we observe that the Fubini–Tonelli Theorem implies

$$\begin{aligned} \mathbb{E} \left[Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N) \right] \\ = \mathbb{E} \mathbb{E} \left[\exp \left(\sum_{n=1}^N \sum_{m=1}^M (\beta\omega_{n,m} - \log Q(\beta)) \delta_{n,m} \right) \mathbf{1}_{A_N} \right] = \mathbf{P}(A_N), \end{aligned}$$

with $\delta_{n,m} := \mathbf{1}_{(n,m) \in \tau}$. Hence $\lim_{N \rightarrow \infty} \mathbb{E} \left[Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N) \right] = 0$ and (3.3) follows because

$$\begin{aligned} (3.4) \quad \mathbb{P} \left(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} > \frac{1}{2} \text{ and } \mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N) > \frac{1}{2} \right) \\ = \mathbb{P} \left(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \mathbf{1}_{\{\mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N) > \frac{1}{2}\}} > \frac{1}{2} \right) \leq 2 \mathbb{E} \left[Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \mathbf{1}_{\{\mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N) > \frac{1}{2}\}} \right] \\ \leq 4 \mathbb{E} \left[Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N) \right]. \quad \square \end{aligned}$$

We now prove that $\{Z_{N,M(N),\omega}^{f,\beta,h_c^a(\beta)}\}_N$ is uniformly integrable (and this holds for an arbitrary choice of $M(N)$) provided that the intersection renewal $\sigma = \tau \cap \tau'$ is terminating (τ and τ' are two independent copies of τ) and β is small enough. Let us point out that, since σ is a terminating renewal then the total number $|\sigma|$ of renewal points (except the origin), that is $|\sigma| = \sum_{(n,m) \in \mathbb{N}^2} \tilde{\delta}_{n,m}$ with $\tilde{\delta}_{n,m} = \mathbf{1}_{(n,m) \in \sigma}$, is a geometric random variable of parameter $\mathbf{P}^{\otimes 2}(\sigma_1 < \infty)$, where $\sigma_1 < \infty$ simply means that both components of σ_1 are finite. This in particular implies that $\mathbf{P}^{\otimes 2}(\sigma_1 < \infty) = 1/\mathbf{E}^{\otimes 2}[|\sigma|]$. Moreover it is straightforward to see that $\mathbf{E}^{\otimes 2}[|\sigma|] = \sum_{n,m} \mathbf{P}((n,m) \in \tau)^2$.

LEMMA 3.2. — *If $\sigma := \tau \cap \tau'$ is terminating, then defining*

$$(3.5) \quad 0 < \beta_1 := \sup \left\{ \beta : \log Q(2\beta) - 2 \log Q(\beta) < -\log \mathbf{P}^{\otimes 2}(\sigma_1 < \infty) \right\},$$

we have that for every $\beta \in (0, \beta_1)$ the sequence $\{Z_{N,M(N),\omega}^{f,\beta,h_c^a(\beta)}\}_N$ is bounded in $L^2(\mathbb{P})$, and is therefore uniformly integrable.

Proof. — We write $M = M(N)$ and we compute the second moment of the partition function:

$$(3.6) \quad \begin{aligned} \mathbb{E} \left[\left(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \right)^2 \right] &= \mathbf{E}^{\otimes 2} \left[\mathbb{E} \left[\exp \left(\sum_{n=1}^N \sum_{m=1}^M (\beta \omega_{n,m} - \log Q(\beta)) (\delta_{n,m} + \delta'_{n,m}) \right) \right] \right] \\ &= \mathbf{E}^{\otimes 2} \left[\exp \left(\sum_{n=1}^N \sum_{m=1}^M (\log Q(2\beta) - 2 \log Q(\beta)) \tilde{\delta}_{n,m} \right) \right]. \end{aligned}$$

The sequence $\{Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)}\}_N$ is bounded in $L^2(\mathbb{P})$ if

$$(3.7) \quad \mathbf{E}^{\otimes 2} \left[\exp \left(\sum_{n,m=1}^{\infty} (\log Q(2\beta) - 2 \log Q(\beta)) \tilde{\delta}_{n,m} \right) \right] < \infty.$$

Since $|\sigma|$ is a geometric random variable of parameter $\mathbf{P}^{\otimes 2}(\sigma_1 < \infty)$, (3.7) holds if

$$(3.8) \quad \log Q(2\beta) - 2 \log Q(\beta) < -\log \mathbf{P}^{\otimes 2}(\sigma_1 < \infty). \quad \square$$

Proof of Theorem 1.3. — In view of what we want to prove and of Proposition 2.1, notably the explicit continuity estimate (2.3), it suffices to establish the result for $\gamma \in \mathbb{Q}$ and for $M = \lfloor \gamma N \rfloor$, which we shall assume till the end of the proof, even if this explicit choice is used in full only at the very end.

Because of Lemma 3.2, we have that the sequence $\{Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)}\}_N$ is uniformly integrable for $\beta < \beta_1$. Now for all $0 < \eta < \alpha$ (recall that if σ is terminating, it implies that $\alpha \leq 1$), we set

$$(3.9) \quad A_N := \{|\tau \cap ((0, N] \times (0, M])| \leq N^\eta\}.$$

From Lemma A.2, we have that $\lim_N \mathbf{P}(A_N) = 0$. Observe also that

$$(3.10) \quad \begin{aligned} Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)+h} &= Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \mathbf{E}_{N,M,\omega}^{f,\beta,h_c^a(\beta)} [\exp(h|\tau \cap ((0, N] \times (0, M])|)] \\ &\geq Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N^c) \exp(hN^\eta). \end{aligned}$$

Let us call E_N the event whose probability is estimated from below in (3.1). Then on E_N , whose probability is at least $\zeta > 0$, we have

$$(3.11) \quad Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)+h} \geq \frac{1}{2} \left(1 - \mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N) \right) \exp(hN^\eta) \geq \frac{1}{4} \exp(hN^\eta).$$

Therefore we obtain

$$(3.12) \quad \mathbb{P} \left(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)+h} \geq \frac{1}{4} \exp(hN^\eta) \right) \geq \mathbb{P}(E_N) \geq \zeta.$$

Our aim is to prove that $F_\gamma(\beta, h + h_c^a(\beta)) > 0$ or more precisely give a lower bound for $F_\gamma(\beta, h + h_c^a(\beta))$. We aim at using (2.2), this is why we have chosen $\gamma \in \mathbb{Q}$, and now we choose also N such that $\gamma N \in \mathbb{N}$, so $N = jq$, $j \in \mathbb{N}$ ($\gamma = p/q$). Since the first part of the proof exploits the free partition function, and not the constrained one for which (2.2) holds, we use Lemma 2.2 that guarantees that

$$(3.13) \quad \log Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)+h} \leq \log Z_{N,M,\omega}^{c,\beta,h_c^a(\beta)+h} + c_1 \left(1 + \log(N + M) + \beta |\omega_{N,M}| \right).$$

Since there exists $c_2 > 1$ such that $\beta|\omega_{N,M}| < c_2 \log(N + M)$ with probability at least $1 - \zeta/2$, and recalling that $M \sim \gamma N$, we get that there exists $c_3 > 0$ such that

$$\mathbb{P} \left(\log Z_{N,M,\omega}^{c,\beta,h_c^a(\beta)+h} \leq \log Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)+h} - c_3 \log N \right) \leq \frac{\zeta}{2}.$$

Combining this with (3.12), we get that

$$(3.14) \quad \mathbb{P} \left(\log Z_{N,M,\omega}^{c,\beta,h_c^a(\beta)+h} \geq \frac{1}{2} h N^\eta - c_3 \log N \right) \geq \frac{\zeta}{2}.$$

Now using the uniform bound $Z_{N,M,\omega}^{c,\beta,h_c^a(\beta)} \geq K(N + M)e^{\beta\omega_{N,M} - \log Q(\beta)}$ on the event E_n^c , we arrive at

$$(3.15) \quad \begin{aligned} \mathbb{E} \log Z_{N,M,\omega}^{c,\beta,h_c^a(\beta)+h} &\geq \frac{\zeta}{4} h N^\eta - \frac{c_3 \zeta}{2} \log N + \log K(N + M) - \beta \mathbb{E}[\omega_{1,1}] - \log Q(\beta) \\ &\geq c_4 h N^\eta - c_5 \log N, \end{aligned}$$

for suitably chosen $c_4, c_5 > 0$.

At this point the choice $\gamma = p/q$ and $M = \gamma N \in \mathbb{N}$ enters the game. By (2.2) we have

$$(3.16) \quad F_\gamma(\beta, h_c^a(\beta) + h) \geq \sup_{N=jq: j=j_0, j_0+1, \dots} \left\{ c_4 h N^{\eta-1} - c_5 N^{-1} \log N \right\},$$

and the fact that j has to be chosen larger than a certain j_0 just reflects the fact that the estimates in this proof have been performed for a N larger than a suitable N_0 . We now estimate from below the right-hand side in (3.16) by choosing $N = h^{-\frac{1+\varepsilon}{\eta}}$ (for some $\varepsilon > 0$ fixed): this means that we have chosen $h = (jq)^{-\eta/(1+\varepsilon)}$. With this choice

$$(3.17) \quad F_\gamma(\beta, h_c^a(\beta) + h) \geq c_4 h^{-\varepsilon} N^{-1} - c_5 \frac{1 + \varepsilon}{\eta} N^{-1} \log \frac{1}{h} \geq N^{-1} = h^{\frac{1+\varepsilon}{\eta}},$$

where the last inequality holds provided that h is small enough. This is the estimate we were after since we can choose η arbitrarily close to α and ε close to 0, but we have established it only for h of the form $(jq)^{-\eta/(1+\varepsilon)}$, $j = j_0, j_0 + 1, \dots$. However, we can use that $h \mapsto F_\gamma(\beta, h_c^a(\beta) + h)$ is non decreasing: having demonstrated that $F_\gamma(\beta, h_c^a(\beta) + h) \geq h^{\frac{1+\varepsilon}{\eta}}$ for $h = h_j := (jq)^{-\eta/(1+\varepsilon)}$ implies that $F_\gamma(\beta, h_c^a(\beta) + h) \geq \frac{1}{2} h^{\frac{1+\varepsilon}{\eta}}$ for every sufficiently small h (this can be verified by checking that $\frac{1}{2} h_j^{\frac{1+\varepsilon}{\eta}}$ is smaller than $h_{j+1}^{\frac{1+\varepsilon}{\eta}}$). This completes the proof of Theorem 1.3. \square

The technique used to prove Theorem 1.3 could be adapted for $\alpha > 1$ to deduce the upper bound for the difference between quenched and annealed critical points.

PROPOSITION 3.3. — *Let $\alpha > 1$. There exists a slowly varying function $\tilde{L}(\cdot)$ such that*

$$(3.18) \quad h_c^q(\beta) - h_c^a(\beta) \leq \tilde{L}(1/\beta) \beta^{\frac{2\alpha}{\alpha-1} \vee 4},$$

for $\beta \leq 1$.

Proof. — As in the previous proof, it suffices to work with the case $\gamma = p/q \in \mathbb{Q}$ and $M = \lfloor \gamma N \rfloor$. We set

$$(3.19) \quad N_\beta := \max \left\{ N \in q\mathbb{N} : \mathbb{E} \left[\left(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \right)^2 \right] \leq 2 \right\}.$$

(Note that in view of (3.23) below, $\mathbb{E} \left[\left(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \right)^2 \right]$ is non-decreasing in N).

Using Paley–Zygmund Inequality, we therefore get that $\mathbb{P}(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} > 1/2) \geq 1/8$ for any $N \leq N_\beta$, and we can then adapt the proof of Proposition 3.1.

Let us take $A_N := \{|\tau \cap ((0, N] \times (0, M])| \leq N/2\mu\}$. Since $\lim_{N \rightarrow \infty} \mathbf{P}(A_N) = 0$, and $\mathbb{P}(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} > 1/2) > 1/8$ for $N \leq N_\beta$, we find, exactly as in the proof of Proposition 3.1, that there exists $N_0 \in \mathbb{N}$ such that for every $N_0 \leq N \leq N_\beta$ we have

$$(3.20) \quad \mathbb{P} \left(\mathbf{P}_{N,M,\omega}^{f,\beta,h_c^a(\beta)}(A_N) \leq \frac{1}{2} \text{ and } Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} > \frac{1}{2} \right) \geq \frac{1}{20}.$$

Following the proof of Theorem 1.3 (see (3.16)), provided $N_\beta \geq N_0$, and since $N_\beta \in q\mathbb{N}$, we get that

$$(3.21) \quad F_\gamma(\beta, h) \geq \left\{ c_6(h - h_c^a(\beta)) - c_7 N_\beta^{-1} \log N_\beta \right\}.$$

We therefore observe that if $h - h_c^a(\beta) > c_7/c_6 N_\beta^{-1} \log N_\beta$ then $F_\gamma(\beta, h) > 0$. Hence we get that

$$(3.22) \quad h_c^q(\beta) - h_c^a(\beta) \leq \frac{c_7}{c_6} \cdot \frac{\log N_\beta}{N_\beta}.$$

It therefore boils down to estimating N_β , namely obtaining a lower bound. Recall from (3.6) that

$$(3.23) \quad \mathbb{E} \left[\left(Z_{N,M,\omega}^{f,\beta,h_c^a(\beta)} \right)^2 \right] = \mathbf{E}^{\otimes 2} \left[\exp \left((\log Q(2\beta) - 2 \log Q(\beta)) \mathcal{H}_{N,M}(\sigma) \right) \right],$$

with $\mathcal{H}_{N,M}(\sigma) = \sum_{n=1}^N \sum_{m=1}^M \mathbf{1}_{(n,m) \in \sigma}$, and $\sigma = \tau \cap \tau'$ the intersection renewal. Recall that σ is persisting for $\alpha > 1$, see Proposition A.3.

Note that for $\beta \leq 1$, there exists c_8 such that $\log Q(2\beta) - 2 \log Q(\beta) \leq c_8 \beta^2$, and that $\mathcal{H}_{N,M} \leq \mathcal{H}_{M,M}$ if $M \geq N$. The question is therefore reduced to estimating $\mathbf{E}^{\otimes 2} [\exp(t\mathcal{H}_{M,M}(\sigma))]$, with $t = c_8 \beta^2$. We have

$$(3.24) \quad \begin{aligned} \mathbf{E}^{\otimes 2} \left[\exp(t\mathcal{H}_{M,M}(\sigma)) \right] &= 1 + \sum_{k=1}^M \left(e^{tk} - e^{t(k-1)} \right) \mathbf{P}^{\otimes 2} (\mathcal{H}_{M,M}(\sigma) \geq k) \\ &\leq 1 + (e^t - 1) \sum_{k=1}^M e^{tk} \mathbf{P}^{\otimes 2} (\mathcal{H}_{M,M}(\sigma) \geq k). \end{aligned}$$

In order to obtain an upper bound, we use the following fact

$$(3.25) \quad \mathbf{P}^{\otimes 2} (\mathcal{H}_{M,M}(\sigma) \geq k) = \mathbf{P}^{\otimes 2} (\sigma_k \in (0, M]^2) \leq \mathbf{P}^{\otimes 2} (\sigma_1 \in (0, M]^2)^k.$$

Then we get

$$(3.26) \quad \mathbf{E}^{\otimes 2} [\exp(t\mathcal{H}_{M,M}(\sigma))] \leq 1 + (e^t - 1) \sum_{k=1}^M \exp \left[k \left(t + \log \mathbf{P}^{\otimes 2} (\sigma_1 \in (0, M]^2) \right) \right].$$

Let $\underline{\sigma} := \sigma^{(1)} + \sigma^{(2)}$. An elementary observation is that $\mathbf{P}^{\otimes 2}(\sigma_1 \notin (0, M]^2)$ is of the same order as $\mathbf{P}^{\otimes 2}(\underline{\sigma}_1 > M)$: indeed, for every $M \in \mathbb{N}$ we have

$$(3.27) \quad \mathbf{P}^{\otimes 2}(\underline{\sigma}_1 > 2M) \leq \mathbf{P}^{\otimes 2}(\sigma_1 \notin (0, M]^2) \leq \mathbf{P}^{\otimes 2}(\underline{\sigma}_1 > M).$$

Therefore, using that $\log(1 - x) \leq -x$ for $x \in [0, 1]$, we get that

$$(3.28) \quad \log \mathbf{P}^{\otimes 2}(\sigma_1 \in (0, M]^2) \leq -\mathbf{P}^{\otimes 2}(\underline{\sigma}_1 > 2M) \leq -c_9/U_{M,M},$$

where we used Lemma A.6 to estimate $\mathbf{P}^{\otimes 2}(\underline{\sigma}_1 > 2M)$ (provided that M is large enough), with $U_{N,M} = \sum_{n=0}^N \sum_{m=0}^M \mathbf{P}((n, m) \in \tau)^2$. Since $M \leq \gamma N$ and $U_{N,N}$ is regularly varying, see Proposition A.3, we get that

$$(3.29) \quad \mathbf{E}^{\otimes 2}[\exp(t\mathcal{H}_{M,M}(\sigma))] \leq 1 + (e^t - 1) \sum_{k=1}^N \exp\left(k(t - c_{10}/U_{N,N})\right).$$

We therefore choose N such that $c_{10}/U_{N,N} \geq 3t = 3c_8\beta^2$. By Proposition A.3, for $\alpha > 1$, we can choose

$$(3.30) \quad N = \tilde{\psi}(1/\beta)\beta^{-(\frac{2\alpha}{\alpha-1}V^4)},$$

for some slowly varying function $\tilde{\psi}(\cdot)$. For this choice of N , we therefore get that

$$(3.31) \quad \mathbb{E}\left[\left(Z_{N,M,\omega}^{f,\beta,h_c^\alpha(\beta)}\right)^2\right] \leq 1 + (e^{c_8\beta^2} - 1) \sum_{k=1}^N \exp(-2c_8\beta^2 k) \leq 1 + \frac{e^{c_8\beta^2} - 1}{1 - e^{-2c_8\beta^2}},$$

which is smaller than 2 provided that β is small enough. It therefore implies that there exists some $\beta_1 > 0$ such that

$$(3.32) \quad N_\beta \geq \tilde{\psi}(1/\beta)\beta^{-(\frac{2\alpha}{\alpha-1}V^4)} \quad \text{for } \beta \leq \beta_1.$$

The proof is therefore complete by putting (3.32) in (3.22). □

4. Lower bound on the critical point shift

From now on, $L_i(\cdot)$ will denote slowly varying functions and C_i positive constants for $i = 1, 2, \dots$. Also, we sometimes treat certain large quantities as if they were integers, simply to avoid the integer-part notation; in all cases these can be treated as if the integer-part notation were in use.

Our proof is based on combining the fractional moment method and a change of measure argument, following the same strategy adopted in [DGLT09]. Let

$$(4.1) \quad z_{n,m} := \exp(\beta\omega_{n,m} + h).$$

Choose $k \leq N$ and M such that $M \sim \gamma N$ and decompose the partition function (1.4) as follows, see Figure 4.1:

$$(4.2) \quad Z_{N,M,\omega} = Z_{N,M,\omega}^1 + Z_{N,M,\omega}^2 + Z_{N,M,\omega}^3,$$

with (recall the notation (2.20))

$$\begin{aligned} Z_{N,M,\omega}^1 &= \sum_{n=k}^N \sum_{m=k}^M Z_{N-n,M-m,\omega} \sum_{i=0}^{k-1} \sum_{j=0}^{k-1} K(n-i+m-j) z_{N-i,M-j} Z_{(N-i,M-j),(N,M),\omega}, \\ Z_{N,M,\omega}^2 &= \sum_{n=1}^{k-1} \sum_{m=k}^M Z_{N-n,M-m,\omega} \sum_{i=0}^{n-1} \sum_{j=0}^{k-1} K(n-i+m-j) z_{N-i,M-j} Z_{(N-i,M-j),(N,M),\omega}, \\ Z_{N,M,\omega}^3 &= \sum_{n=k}^N \sum_{m=1}^{k-1} Z_{N-n,M-m,\omega} \sum_{i=0}^{k-1} \sum_{j=0}^{m-1} K(n-i+m-j) z_{N-i,M-j} Z_{(N-i,M-j),(N,M),\omega}. \end{aligned}$$

Note that $Z_{(N-i,M-j),(N,M),\omega}$ has the same law as $Z_{i,j,\omega}$ and that $Z_{N-n,M-m,\omega}$, $z_{N-i,M-j}$ and $Z_{(N-i,M-j),(N,M),\omega}$ are independent for $i < n$ and $j < m$.

Let $\delta \in (0, 1)$ (that will be chosen close to 1 later in the proof), and define

$$(4.3) \quad \mathcal{A}_{N,M} := \mathbb{E} \left[(Z_{N,M,\omega})^\delta \right] \quad \text{for every } N, M \in \mathbb{N}^2,$$

with $\mathcal{A}_{0,0} = 1$, and $\mathcal{A}_{i,0} = \mathcal{A}_{0,i} = 0$ for every $i \geq 1$. We apply the inequality $(\sum a_i)^\delta \leq \sum a_i^\delta$ (which holds for any finite and countable collection of positive real numbers) to the decomposition (4.2) to get

$$(4.4) \quad \mathcal{A}_{N,M} \leq \mathcal{A}_{N,M}^1 + \mathcal{A}_{N,M}^2 + \mathcal{A}_{N,M}^3,$$

where

$$(4.5) \quad \begin{aligned} \mathcal{A}_{N,M}^1 &\leq \mathbb{E}[z_{1,1}^\delta] \sum_{n=k}^N \sum_{m=k}^M \mathcal{A}_{N-n,M-m} \sum_{i=0}^{k-1} \sum_{j=0}^{k-1} K(n-i+m-j)^\delta \mathcal{A}_{i,j}, \\ \mathcal{A}_{N,M}^2 &\leq \mathbb{E}[z_{1,1}^\delta] \sum_{n=1}^{k-1} \sum_{m=k}^M \mathcal{A}_{N-n,M-m} \sum_{i=0}^{n-1} \sum_{j=0}^{k-1} K(n-i+m-j)^\delta \mathcal{A}_{i,j}, \\ \mathcal{A}_{N,M}^3 &\leq \mathbb{E}[z_{1,1}^\delta] \sum_{n=k}^N \sum_{m=1}^{k-1} \mathcal{A}_{N-n,M-m} \sum_{i=0}^{k-1} \sum_{j=0}^{m-1} K(n-i+m-j)^\delta \mathcal{A}_{i,j}. \end{aligned}$$

The key idea of the proof is to the following proposition.

PROPOSITION 4.1. — *For fixed β and h , if there exists $k \in \mathbb{N}$ such that $\rho_1 + \rho_2 + \rho_3 \leq 1$ with*

$$(4.6) \quad \rho_1 + \rho_2 + \rho_3 := \mathbb{E}[z_{1,1}^\delta] \left(\sum_{n=k}^\infty \sum_{m=k}^\infty \sum_{i=0}^{k-1} \sum_{j=0}^{k-1} + \sum_{n=1}^{k-1} \sum_{m=k}^\infty \sum_{i=0}^{n-1} \sum_{j=0}^{k-1} + \sum_{n=k}^\infty \sum_{m=1}^{k-1} \sum_{i=0}^{k-1} \sum_{j=0}^{m-1} \right) K(n-i+m-j)^\delta \mathcal{A}_{i,j},$$

then $F_\gamma(\beta, h) = 0$.

Proof. — Define $\bar{A} := \max \{ \{ \mathcal{A}_{i,j}, \mathcal{A}_{i,s}, \mathcal{A}_{t,j} \}, 1 \leq i, j \leq k-1, s, t \geq k \}$. Note that by Jensen's inequality we have $\mathcal{A}_{i,j} \leq \mathbb{E}[Z_{i,j,\omega}^\delta] \leq \exp(\delta h \min\{i, j\})$, since there are at most $\min\{i, j\}$ renewals in the region $\{1, \dots, i\} \times \{1, \dots, j\}$: we get that $\bar{A} \leq e^{hk}$.

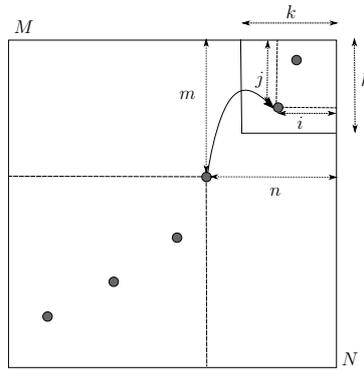


Figure 4.1. Fixing a value k , the partition function is decomposed by summing over the values of the last renewal epoch outside the corner block $(N - k, N] \times (M - k, M]$, and the first one inside that block. We distinguish three cases: either the last renewal epoch is in $[0, N - k] \times [0, M - k]$ (which is the case represented in the figure, giving $Z_{N,M,\omega}^1$), or it is in $(N - k, N] \times [0, M - k]$ ($Z_{N,M,\omega}^2$) or in $[0, N - k] \times (M - k, M]$ ($Z_{N,M,\omega}^3$).

Then from (4.4) and the fact that $\rho_1 + \rho_2 + \rho_3 \leq 1$, we deduce (by induction) that $\mathcal{A}_{N,M} \leq \bar{A} \leq e^{kh}$ for all N, M . Then by Jensen’s inequality

$$(4.7) \quad F_\gamma^q(\beta, h) = \lim_{\substack{N \rightarrow \infty \\ M/N \rightarrow \gamma}} \frac{1}{\delta N} \mathbb{E} \log (Z_{N,M,\omega})^\delta \leq \lim_{\substack{N \rightarrow \infty \\ M/N \rightarrow \gamma}} \frac{1}{\delta N} \log \mathcal{A}_{N,M} = 0. \quad \square$$

Our aim is therefore to prove that for $h = h_c^a(\beta) + \Delta_\beta^\varepsilon$ (where Δ_β^ε is defined in Theorem 1.4) we have that $F_1^q(\beta, h) = 0$ (provided that β is small enough), by showing that ρ_1, ρ_2, ρ_3 are smaller than $1/3$ for such h , for some $k = k_\beta$ wisely chosen. For the choice of k , we pick k proportional to the correlation length of the annealed system, that is $k \propto F(0, \Delta_\beta^\varepsilon)^{-1}$, and in view of Theorem 1.2 (here $\alpha > 1$), we can take

$$(4.8) \quad k = k_\beta = \frac{1}{\Delta_\beta^\varepsilon} = \begin{cases} \beta^{-(1+\varepsilon)\frac{2\alpha}{\alpha-1}} & \text{if } \alpha \in (1, 2], \\ \beta^{-4} |\log \beta|^6 & \text{if } \alpha > 2. \end{cases}$$

Note that, in view of (4.6) and (1.1), provided that δ is close to 1 so that $(2 + \alpha)\delta > 2$, we have

$$(4.9) \quad \rho_1 \leq \sum_{i=0}^{k-1} \sum_{j=0}^{k-1} \frac{L_1(2k - i - j)}{(2k - i - j)^{(2+\alpha)\delta - 2}} \mathcal{A}_{i,j},$$

and

$$(4.10) \quad \rho_2 \leq \sum_{i=0}^{k-1} \sum_{j=0}^{k-1} \frac{L_2(k - j)}{(k - j)^{(2+\alpha)\delta - 2}} \mathcal{A}_{i,j}.$$

The ρ_3 case being symmetric to ρ_2 , we can therefore focus on ρ_1 and ρ_2 .

4.1. Finite-volume fractional moment estimate

To estimate (4.9) and (4.10), we need a good control over the fractional moment $\mathcal{A}_{i,j}$ for any $i, j \leq k$, and we provide estimates in this section.

First of all, using Jensen’s Inequality, we have that $\mathcal{A}_{i,j} \leq (\mathbb{E}Z_{i,j,\omega})^\delta$. Moreover, because $h = h_c^\alpha(\beta) + \Delta_\beta^\varepsilon$, we get that for any $i, j \leq k$

$$\mathbb{E}Z_{i,j,\omega} = \mathbf{E} \left[\exp \left(\Delta_\beta^\varepsilon |\tau \cap \{1, \dots, i\} \times \{1, \dots, j\}| \right) \mathbf{1}_{(i,j) \in \tau} \right] \leq e \mathbf{P}((i, j) \in \tau),$$

since $|\tau \cap \{1, \dots, i\} \times \{1, \dots, j\}| \leq k$ and thanks to our choice of $k = (\Delta_\beta^\varepsilon)^{-1}$. We therefore get that,

$$(4.11) \quad \mathcal{A}_{i,j} \leq e^\delta \mathbf{P}((i, j) \in \tau)^\delta,$$

and $\mathbf{P}((i, j) \in \tau)$ can be estimated thanks to Theorems A.4–A.5.

However, this estimate is rather rough, especially when i, j is close to the diagonal (that is for example $i \leq j \leq i + a_i$ where $(a_n)_{n \geq 0}$ is the scaling sequence for τ_n , defined in Section 1.4). We therefore prove the following proposition:

PROPOSITION 4.2. — *Let $h = h_c^\alpha(\beta) + \Delta_\beta^\varepsilon$ and $k = (\Delta_\beta^\varepsilon)^{-1}$. Then, define also*

$$(4.12) \quad \ell_i := \begin{cases} i^{(1+\varepsilon^3)/\alpha} & \text{if } \alpha \in (1, 2], \\ C\sqrt{i \log i} & \text{if } \alpha > 2, \end{cases}$$

so in any case $\ell_i \gg a_i$. There exists some k_0 such that, provided that $k \geq k_0$ then for all $\sqrt{k} \leq i \leq k$ and $i \leq j \leq i + \ell_i$ we have that

$$(4.13) \quad \mathcal{A}_{i,j} \leq L_{10}(i) \left(i^{\delta(1-\frac{1}{\alpha\lambda^2})} (\ell_i)^{-\delta\alpha} + \frac{i^{-\delta(1+\alpha)} \ell_i^\delta}{\beta^{2\delta}} + i^{-\frac{\delta}{\alpha\lambda^2}} e^{-c\delta(\beta^2 i / \ell_i)^{1/2}} \right).$$

This result is the core of the proof, and is based on a change of measure argument. With this result in hand, we are able to show that ρ_1 and ρ_2 are small, for $\alpha > 2$ in Section 4.2 and for $\alpha \in (1, 2]$ in Section 4.3. Let us apply this proposition to get bounds on $\mathcal{A}_{i,j}$ in the different cases.

Case $\alpha > 2$. — We get that uniformly for $k/2 \leq i \leq k$ and $i \leq j \leq i + C'\sqrt{k \log k}$

$$(4.14) \quad \begin{aligned} \mathcal{A}_{i,j} &\leq L_{11}(k) k^{\frac{\delta}{2}(1-\alpha)} + L_{12}(k) \frac{k^{-\delta(1+\alpha)} (k \log k)^{\delta/2}}{\beta^{2\delta}} + L_{13}(k) k^{-\delta/2} e^{-c(\beta^4 k / \log k)^{1/4}} \\ &\leq L_{14}(k) \left(k^{\frac{\delta}{2}(1-\alpha)} + k^{-\delta\alpha} + k^{-\delta/2} e^{-c(\log k)^{5/4}} \right) \leq L_{15}(k) k^{-\frac{\delta}{2}(\alpha-1)}, \end{aligned}$$

where the choice (4.8) of $k = (\Delta_\beta^\varepsilon)^{-1} = \beta^{-4} |\log \beta|^6$ is crucial, in order to get that $\beta^4 k / \log k \geq c(\log k)^5$. For the last inequality, we observe that the first term dominates.

Case $\alpha \in (1, 2]$. — We use also the choice (4.8) of $k = (\Delta_\beta^\varepsilon)^{-1} = \beta^{-(1+\varepsilon)2\alpha/(\alpha-1)}$ to get that uniformly for $k^{1-\varepsilon^2} \leq i \leq k$, we have provided that ε is small enough

$$(4.15) \quad \beta^2 i / \ell_i = k^{-\frac{\alpha-1}{(1+\varepsilon)\alpha}} i^{\frac{\alpha-1-\varepsilon^3}{\alpha}} \geq k^{\frac{1}{(1+\varepsilon)\alpha} (\varepsilon(\alpha-1) + O(\varepsilon^2))} \geq k^{\varepsilon(\alpha-1)/2\alpha}.$$

Therefore, using also that for $i \leq k$, $\beta^{-2}\ell_i \leq k^{\frac{\alpha-1}{(1+\varepsilon)\alpha}} k^{(1+\varepsilon^3)/\alpha} \leq k$ (if ε has been fixed small enough), we have that uniformly for $k^{1-\varepsilon^2} \leq i \leq k$ and $i \leq j \leq i + \ell_i$,

$$(4.16) \quad \begin{aligned} \mathcal{A}_{i,j} &\leq L_{10}(i)i^{-\delta(\frac{1}{\alpha}+\varepsilon^3)} + L_{10}(i)i^{-\delta\alpha}k^\delta + L_{10}(i)i^{-\delta/\alpha}e^{-ck^\varepsilon(\alpha-1)/4\alpha} \\ &\leq L_{16}(k)k^{-\delta(1-\varepsilon^2)(\frac{1}{\alpha}+\varepsilon^3)} \leq L_{16}(k)k^{-\frac{\delta}{\alpha}(1+\varepsilon^2/2)} \end{aligned}$$

where again, for the second to last inequality, we observe that the first term dominates, since $1/\alpha > \alpha$ and ε can be fixed arbitrarily small.

Proof of Proposition 4.2. — The idea is to use a change of measure argument. We define a strip $J_{i,j}$ in which we will tilt the environment by some quantity λ (to be chosen wisely):

$$(4.17) \quad J_{i,j} := \left\{ (n, m) \in \llbracket 0, i \rrbracket \times \llbracket 0, j \rrbracket : |n - m| \leq 2\ell_i \right\},$$

and hence $\#J_{i,j} \leq 2i\ell_i$. The width $2\ell_i$ of the strip is chosen because of the scaling of the bivariate renewal: it is very unlikely that the renewal deviates from the diagonal by more than ℓ_i , see Theorem A.5.

Now, for $\lambda \in \mathbb{R}$ and $i, j \in \mathbb{N}$, we define a new probability measure $\mathbb{P}_{i,j,\lambda}$, under which the $\omega_{n,m}$ are still independent variables, but tilted by λ in the strip $J_{i,j}$:

$$(4.18) \quad \frac{d\mathbb{P}_{i,j,\lambda}}{d\mathbb{P}}(\omega) = \frac{1}{Q(-\lambda)^{\#J_{i,j}}} \exp\left(-\lambda \sum_{(n,m) \in J_{i,j}} \omega_{n,m}\right),$$

where $Q(\cdot)$ is defined in (1.2). Observe now that by Hölder inequality

$$(4.19) \quad \mathcal{A}_{i,j} = \mathbb{E}_{i,j,\lambda} \left[(Z_{i,j,\omega})^\delta \frac{d\mathbb{P}}{d\mathbb{P}_{i,j,\lambda}}(\omega) \right] \leq \mathbb{E}_{i,j,\lambda} [Z_{i,j,\omega}]^\delta \mathbb{E}_{i,j,\lambda} \left[\left(\frac{d\mathbb{P}}{d\mathbb{P}_{i,j,\lambda}}(\omega) \right)^{1/(1-\delta)} \right]^{1-\delta}.$$

The second term in the right-hand side of (4.19) is equal to

$$(4.20) \quad \mathbb{E}_{i,j,\lambda} \left[\left(\frac{d\mathbb{P}}{d\mathbb{P}_{i,j,\lambda}}(\omega) \right)^{1/(1-\delta)} \right]^{1-\delta} = (Q(-\lambda)^\delta Q(\lambda/(1-\delta))^{1-\delta})^{\#J_{i,j}}.$$

Observe that there exists $c_{11} > 0$ such that $0 \leq \log Q(x) \leq c_{11}x^2$ for $|x| \leq 1$. Therefore for $|\lambda| \leq 1 - \delta \leq 1$ and by (4.19) and (4.20), we get

$$(4.21) \quad \mathcal{A}_{i,j} \leq \mathbb{E}_{i,j,\lambda} [Z_{i,j,\omega}]^\delta \exp\left(c_{11} \left(\frac{\delta(1-\delta) + 1}{1-\delta} \right) \lambda^2 \#J_{i,j}\right).$$

Now, we choose $\lambda := (i\ell_i)^{-1/2}$, so that $\lambda^2 \#J_{i,j} \leq 2$, and

$$(4.22) \quad \mathcal{A}_{i,j} \leq e^{4c_{11}/(1-\delta)} \mathbb{E}_{i,j,\lambda} [Z_{i,j,\omega}]^\delta,$$

so that we are left with estimating $\mathbb{E}_{i,j,\lambda} [Z_{i,j,\omega}]$ for $\lambda := (i\ell_i)^{-1/2}$.

Recall (1.10) and the Definition (4.18) of $\mathbb{P}_{i,j,\lambda}$. Using that $\mathbb{E}_{i,j,\lambda}[e^{\beta\omega_{n,m}}]$ equals $Q(\beta) = e^{-h_c^\alpha(\beta)}$ if $(n, m) \notin J_{i,j}$ and $Q(\beta - \lambda)/Q(-\lambda)$ if $(n, m) \in J_{i,j}$, we have that

for every β, h, λ and (i, j)

$$(4.23) \quad \mathbf{E}_{i,j,\lambda} [Z_{i,j,\omega}] = \mathbf{E} \left[e^{(h-h_c^a(\beta))|\tau \cap \llbracket 0, i \rrbracket \times \llbracket 0, j \rrbracket|} \left(\frac{Q(\beta - \lambda)}{Q(\beta)Q(-\lambda)} \right)^{|\tau \cap J_{i,j}|} \mathbf{1}_{(i,j) \in \tau} \right] \\ \leq e \mathbf{E} \left[\left(\frac{Q(\beta - \lambda)}{Q(\beta)Q(-\lambda)} \right)^{|\tau \cap J_{i,j}|} \mathbf{1}_{(i,j) \in \tau} \right],$$

where we used that $|\tau \cap \llbracket 0, i \rrbracket \times \llbracket 0, j \rrbracket| \leq k$ and $h - h_c^a(\beta) = \Delta_\beta^\varepsilon = k^{-1}$.

Now, observe that $\frac{Q(\beta - \lambda)}{Q(\beta)Q(-\lambda)} = 1 - \lambda\beta + O(\lambda^2 + \beta^2)$ as $\lambda, \beta \downarrow 0$. Here, because of our choice (4.8) of $k = (\Delta_\beta^\varepsilon)^{-1}$, we have that $k \geq \beta^{-4}$. Since we are considering $i \geq \sqrt{k}$, and using that $\ell_i \geq \sqrt{i}$, we have that $\lambda := (i\ell_i)^{-1/2} \leq i^{-3/4} \leq k^{-3/8}$, and hence we have that $\lambda \leq \beta^{3/2} \leq \beta$. Therefore, there exists a constant $c_{12} > 0$ such that provided that β is small enough (or k is large enough) we have

$$(4.24) \quad \frac{Q(\beta - \lambda)}{Q(\beta)Q(-\lambda)} \leq \exp(-c_{12}\beta\lambda),$$

and we end up with

$$(4.25) \quad e^{-1} \mathbf{E}_{i,j,\lambda} [Z_{i,j,\omega}] \leq \mathbf{E} \left[e^{-c_{12}\beta\lambda|\tau \cap J_{i,j}|} \mathbf{1}_{(i,j) \in \tau} \right] \\ \leq \mathbf{P}(\exists s, \tau_s \notin J_{i,j}, (i, j) \in \tau) + \mathbf{E} \left[e^{-c_{12}\lambda\beta|\tau \cap \{1, \dots, i\} \times \{1, \dots, j\}|} \mathbf{1}_{(i,j) \in \tau} \right],$$

where in the last term we dropped the indicator function that all renewals occur in the strip $J_{i,j}$. We now estimate these two terms separately.

LEMMA 4.3. — *There exists a slowly varying function L_4 such that, for every $1 \leq i \leq j \leq i + \ell_i$ we have*

$$(4.26) \quad \mathbf{P}(\exists s, \tau_s \notin J_{i,j}, (i, j) \in \tau) \leq L_4(i) i^{1-1/\alpha \wedge 2} (\ell_i)^{-\alpha}.$$

Proof. — Let us first observe that by symmetry, we get that

$$(4.27) \quad \mathbf{P}(\exists s, \tau_s \notin J_{i,j}, (i, j) \in \tau) = 2 \mathbf{P}(\exists s, \tau_s \notin J_{i,j}, \tau_s^{(1)} \leq i/2, (i, j) \in \tau) \\ \leq 2 \sum_{(a,b) \in J_{i,j}} \sum_{\substack{(k,l); (a+k,b+l) \notin J_{i,j} \\ a+k \leq i/2}} \mathbf{P}((a, b) \in \tau) K(k+l) \mathbf{P}((i-a-k, j-b-l) \in \tau).$$

From Theorem A.5, we see that the last term in the double sum of (4.27) is bounded above by c_7/a_i (since $i - a - k \geq i/2$). We get that (4.27) is bounded above by

$$(4.28) \quad \frac{c_{13}}{a_i} \sum_{a=1}^{i/2} \sum_{r=0}^{\ell_i} \mathbf{P}((a, a+r) \in \tau) \sum_{k=1}^{i/2-a} \sum_{l \geq \ell_i - a - r} K(k+l) \\ \leq \frac{c_{14}}{a_i} \mathbf{P}(\exists t \leq i/2, \tau_t \notin \bar{J}_i),$$

with $\bar{J}_i := \{(a, b) : |a - b| \leq \ell_i\}$.

Let us now define

$$(4.29) \quad S_k = \tau_k^{(1)} - \tau_k^{(2)},$$

Then we see that

$$(4.30) \quad \mathbf{P} \left(\exists t \leq i/2, \tau_t \notin \bar{J}_i \right) = \mathbf{P} \left(\max_{t \leq i/2} |S_t| \geq \ell_i \right).$$

Observe that $\{S_k\}$ is a centred random walk in the domain of attraction of a stable law of index $\alpha > 1$. From the Lemma in [SW11] for the case $\alpha \in (1, 2]$ (and infinite variance) and [Pin81, Corollary 1] or [BB01, Equation (12)] for the case $\alpha > 2$ (or $\alpha = 2$ and finite variance) we get that

$$(4.31) \quad \mathbf{P} \left(\max_{t \leq i/2} |S_t| \geq \ell_i \right) \leq i L_3(i) (\ell_i)^{-\alpha}.$$

Therefore by (4.28), (4.30) and (4.31), we obtain (4.26). □

LEMMA 4.4. — *Assume that $i \leq j$ and $\alpha > 1$. There exist constants $c_{15}, c_{16} > 0$ such that, for any sequence $u_i \leq 1$ (we may take $u_i \rightarrow 0$ as $i \rightarrow +\infty$), we have*

$$Z_{i,j}(-u_i) := \mathbf{E} \left[e^{-u_i |\tau \cap \{1, \dots, i\} \times \{1, \dots, j\}|} \mathbb{1}_{(i,j) \in \tau} \right] \leq c_{15} \frac{K(i+j)}{u_i^2} + \mathbf{P}((i, j) \in \tau) e^{-c_{16} i u_i}.$$

In particular, we always have

$$Z_{i,j}(-u_i) \leq \frac{L_5(i) i^{-(2+\alpha)}}{u_i^2} + L_6(i) i^{-1/\alpha \wedge 2} e^{-c_{16} i u_i}$$

Proof. — The last inequality comes from the fact that for $i \leq j$ we have $K(i+j) \leq cL(i) i^{-(2+\alpha)}$, and the fact that Theorem A.5 give $\mathbf{P}((i, j) \in \tau) \leq c_{14}/a_i$ with $a_i = \psi(i) i^{1/\alpha \wedge 2}$.

We write

$$Z_{i,j}(-u_i) = \sum_{k=1}^i e^{-ku_i} \mathbf{P}(\tau_k = (i, j)) = \left(\sum_{k=1}^{i/2\mu} + \sum_{k=i/2\mu}^i \right) e^{-ku_i} \mathbf{P}(\tau_k = (i, j)).$$

For the first sum, we use Theorem A.1 to get that $\mathbf{P}(\tau_k = (i, j)) \leq c_{15} k K(i+j)$ for $k \leq i/2\mu$ ($\leq j/2\mu$), so

$$(4.32) \quad \sum_{k=1}^{i/2\mu} e^{-ku_i} \mathbf{P}(\tau_k = (i, j)) \leq c_{15} \frac{K(i+j)}{u_i^2} \sum_{k=1}^{+\infty} u_i k u_i e^{-ku_i} \leq c_{15} \frac{K(i+j)}{u_i^2},$$

where for the last inequality we bounded the sum by a constant times $\int_{\mathbb{R}_+} x e^{-x} dx$ (thanks to a Riemann–sum approximation for sequences $u_i \rightarrow 0$).

For the second sum we simply bound k by $i/2\mu$ to get that it is smaller than

$$(4.33) \quad e^{-iu_i/2\mu} \sum_{k=i/2\mu}^i \mathbf{P}(\tau_k = (i, j)) \leq \mathbf{P}((i, j) \in \tau) e^{-c_{16} i u_i}.$$

Combining (4.32) and (4.33), we obtain Lemma 4.4. □

Using Lemmas 4.3 and 4.4 in (4.25), and with $u_i = c_{17} \lambda \beta = c_{17} \beta (i \ell_i)^{-1/2} \leq 1$,

$$(4.34) \quad \mathbb{E}_{i,j,\lambda} [Z_{i,j,\omega}] \leq L_7(i) i^{1-1/\alpha \wedge 2} (\ell_i)^{-\alpha} + L_8(i) \frac{i^{-(1+\alpha)} \ell_i}{\beta^2} + L_9(i) i^{-1/\alpha \wedge 2} e^{-c_{17} \beta i^{1/2} \ell_i^{-1/2}}.$$

Finally, this concludes the proof of Proposition 4.2 thanks to (4.22), using that $(a + b + c)^\delta \leq a^\delta + b^\delta + c^\delta$ for $\delta \in (0, 1)$. □

4.2. Conclusion of the proof of Theorem 1.4 in the case $\alpha > 2$

Let $\delta < 1$ be sufficiently close to 1 to have

$$(4.35) \quad (2 + \alpha)\delta > 4,$$

which implies that $\delta(\alpha - 1) > 1$.

We start by estimating ρ_1 . Let R be a large constant and split the sum in (4.9) as

$$(4.36) \quad S_1 + S_2 := \left(\sum_{i,j=0}^{k-R-1} + \sum_{i,j=k-R}^{k-1} \right) \frac{L_1(2k - i - j)}{(2k - i - j)^{(2+\alpha)\delta-2}} \mathcal{A}_{i,j},$$

and

$$(4.37) \quad S_3 + S_4 := \left(\sum_{i=0}^{k-R-1} \sum_{j=k-R}^{k-1} + \sum_{i=k-R}^{k-1} \sum_{j=0}^{k-R-1} \right) \frac{L_1(2k - i - j)}{(2k - i - j)^{(2+\alpha)\delta-2}} \mathcal{A}_{i,j}.$$

Using the fact that $\mathcal{A}_{i,j} \leq e^\delta$ from (4.11), we get

$$(4.38) \quad S_1 \leq \frac{L_{17}(R)}{R^{(2+\alpha)\delta-4}},$$

and the right-hand side of (4.38) can be made small by (4.35) and because R is large.

For S_2 , there exists C_4 such that $S_2 \leq C_4 \max_{k-R \leq i,j < k} \mathcal{A}_{i,j}$, and from (4.11), combined with Theorem A.5, there exists C_5 such that

$$(4.39) \quad \max_{k-R \leq i,j < k} \mathcal{A}_{i,j} \leq e^\delta \max_{k-R \leq i,j < k} \mathbf{P}((i, j) \in \tau)^\delta \leq \frac{C_5}{k^{\delta/2}},$$

then S_2 is arbitrarily small for k large.

Since S_3 and S_4 are the same quantity, we just focus on S_3 . Since $\mathcal{A}_{i,j} \leq e^\delta$ from (4.11), we obtain

$$(4.40) \quad S_3 \leq \frac{L_{18}(R)}{R^{(2+\alpha)\delta-4}},$$

which again can be made small in view of the condition (4.35) and because R is large. Hence ρ_1 can be made arbitrarily small by choosing R large and k large (i.e. β small).

Let us now look at ρ_2 in (4.10). We split the sum to:

$$(4.41) \quad S_5 + S_6 = \left(\sum_{i=0}^{k-1} \sum_{j=0}^i + \sum_{i=0}^{k-1} \sum_{j=i+1}^{k-1} \right) \mathcal{A}_{i,j} \frac{L_2(k - j)}{(k - j)^{(2+\alpha)\delta-2}}.$$

Let us first study S_5 :

$$(4.42) \quad \begin{aligned} S_5 &= \sum_{j=0}^{k-1} \sum_{i=j}^{k-1} \mathcal{A}_{i,j} \frac{L_2(k - j)}{(k - j)^{(2+\alpha)\delta-2}} \\ &= \left(\sum_{j=0}^{k/2} \sum_{i=j}^{k-1} + \sum_{j=k/2+1}^{k-1} \sum_{i=j}^{k-1} \right) \mathcal{A}_{i,j} \frac{L_2(k - j)}{(k - j)^{(2+\alpha)\delta-2}} := S_{5a} + S_{5b}. \end{aligned}$$

Using that $\mathcal{A}_{i,j} \leq e^\delta$ from (4.11), we get

$$(4.43) \quad S_{5a} \leq \frac{L_{19}(k)}{k^{(2+\alpha)\delta-4}}.$$

For S_{5b} , we use (4.11) and Theorem A.5 which gives that if $i \geq j$ $\mathcal{A}_{i,j} \leq \text{cst} \cdot j^{-\delta/2}$ to get

$$(4.44) \quad S_{5b} \leq \sum_{j=k/2+1}^{k-1} \sum_{i=j}^{k-1} \frac{C_6}{j^{\delta/2}} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}} \leq \frac{C_7}{k^{\delta/2}}.$$

Then S_5 can be made small for k large and from condition (4.35).

Now we split S_6 as

$$(4.45) \quad S_{6a} + S_{6b} = \left(\sum_{i=0}^{k/2} + \sum_{i=k/2+1}^{k-1} \right) \sum_{j=(i+\ell_i+1) \wedge (k-1)}^{k-1} \mathcal{A}_{i,j} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}}.$$

and

$$(4.46) \quad S_{6c} + S_{6d} = \left(\sum_{i=0}^{k/2} + \sum_{i=k/2+1}^{k-1} \right) \sum_{j=i+1}^{(i+\ell_i) \wedge (k-1)} \mathcal{A}_{i,j} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}}.$$

Using (4.11) and Theorem A.5, we see that

$$(4.47) \quad \begin{aligned} S_{6a} &\leq \sum_{i=0}^{k/2} \sum_{j=1}^{3k/4} \frac{L_{20}(k-j)}{(k-j)^{(2+\alpha)\delta-2}} + \sum_{i=0}^{k/2} \sum_{j=3k/4}^{k-1} \frac{C_8 i^\delta L(j-i)^\delta}{(j-i)^{(1+\alpha)\delta}} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}} \\ &\leq \frac{L_{21}(k)}{k^{(2+\alpha)\delta-4}} + \frac{L_{22}(k)}{k^{\alpha\delta-1}}, \end{aligned}$$

and

$$\begin{aligned} S_{6b} &\leq \sum_{i=k/2+1}^{k-1} \sum_{j=(i+c\sqrt{k \log k}) \wedge (k-1)}^{k-1} \frac{C_8 i^\delta L(j-i)^\delta}{(j-i)^{(1+\alpha)\delta}} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}} \\ &\leq C_8 k^\delta \sum_{j=k/2}^{k-1} \sum_{x \geq c\sqrt{k \log k}} \frac{L(x)^\delta}{x^{(1+\alpha)\delta}} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}} \leq \frac{L_{23}(k) k^\delta}{k^{\frac{1}{2}((1+\alpha)\delta-1)}} = \frac{L_{23}(k)}{k^{\frac{1}{2}((\alpha-1)\delta-1)}}. \end{aligned}$$

Hence, both S_{6a} and S_{6b} are arbitrarily small for k large, by the condition (4.35).

By (4.11), and since provided that k is large enough we have $i + \sqrt{i \log i} \leq 3k/4$ for $i \leq k/2$, we obtain

$$(4.48) \quad S_{6c} \leq C_9 \sum_{i=0}^{k/2} \sqrt{i \log i} \frac{L_2(k)}{k^{(2+\alpha)\delta-2}} \leq \frac{L_{24}(k)}{k^{(2+\alpha)\delta-7/2}},$$

which is arbitrarily small for k large.

For the term S_{6d} , since for every $j \in \{k/2 + 2, \dots, k - 1\}$ there are at most $C_{10} \sqrt{k \log k}$ corresponding terms in the sum over i , we have

$$(4.49) \quad S_{6d} \leq C_{11} \sqrt{k \log k} \max_{\substack{k/2 \leq i \leq k \\ i \leq j \leq i + \sqrt{i \log i}}} \mathcal{A}_{i,j}.$$

Then we use Proposition 4.2, and more precisely (4.14), to get that

$$(4.50) \quad S_{6d} \leq L_{25}(k) k^{-\frac{1}{2}(\delta(\alpha-1)-1)}$$

In view of the condition (4.35), S_{6d} can be made arbitrarily small for k large. This completes the proof of (1.13) in the case $\alpha > 2$.

4.3. Conclusion of the proof of Theorem 1.4 in the case $\alpha \in (1, 2]$

Fix $\varepsilon > 0$ small and let $0 < \delta < 1$ such that

$$(4.51) \quad \delta \left[(2 + \alpha) + (1 - \varepsilon^2)/\alpha \right] > 4 - \varepsilon^2.$$

which implies in particular that $\delta(2 + \alpha) > 3$. We also assume that

$$(4.52) \quad \delta\alpha > 1, \quad \delta(1 + (1 + \alpha)\varepsilon^4) > 1 + \varepsilon^4, \quad \text{and} \quad \delta > (1 + \varepsilon^4)/(1 + \varepsilon^2/2).$$

Let us start with showing that ρ_1 is small: we split the sum in (4.9) to

$$(4.53) \quad T_1 + T_2 := \left(\sum_{i,j=0}^{k^{1-\varepsilon^2}} + \sum_{i,j=k^{1-\varepsilon^2}+1}^{k-1} \right) \frac{L_1(2k - i - j)}{(2k - i - j)^{(2+\alpha)\delta-2}} \mathcal{A}_{i,j},$$

and

$$(4.54) \quad T_3 + T_4 := \left(\sum_{i=0}^{k^{1-\varepsilon^2}} \sum_{j=k^{1-\varepsilon^2}+1}^{k-1} + \sum_{i=k^{1-\varepsilon^2}+1}^{k-1} \sum_{j=0}^{k^{1-\varepsilon^2}} \right) \frac{L_1(2k - i - j)}{(2k - i - j)^{(2+\alpha)\delta-2}} \mathcal{A}_{i,j}.$$

For $\alpha \leq 2$, we know that there exists a slowly varying function $\psi(\cdot)$ such that $a_i = \psi(i)i^{1/\alpha}$. For T_1 , using (4.11) and Theorem A.5, we get

$$(4.55) \quad T_1 \leq \frac{L_{26}(k)}{k^{(2+\alpha)\delta-2}} \sum_{i,j=0}^{k^{1-\varepsilon^2}} \frac{1}{(a_{\min(i,j)})^\delta} \leq \frac{L_{27}(k)}{k^{(1-\varepsilon^2)(\delta/\alpha-2)+(2+\alpha)\delta-2}},$$

and from the Condition (4.51), T_1 can be made small for k large.

For T_2 , since $(2 + \alpha)\delta - 2 \in (1, 2)$, we have

$$(4.56) \quad T_2 \leq L_{28}(k)k^{-(2+\alpha)\delta+4} \max_{k^{1-\varepsilon^2} \leq i,j \leq k} \mathcal{A}_{i,j} < \frac{L_{29}(k)}{k^{(1-\varepsilon^2)\delta/\alpha+(2+\alpha)\delta-4}},$$

where for the last inequality we used (4.11) and Theorem A.5. Then T_2 is small for k large thanks to (4.51).

For T_3 (which is equal to T_4), since for the range of i, j considered we have $2k - i - j \geq k/2$, we get using (4.11) and Theorem A.5

$$(4.57) \quad T_3 \leq \sum_{i=0}^{k^{1-\varepsilon^2}} \frac{1}{(a_i)^\delta} \frac{L_{30}(k)}{k^{(2+\alpha)\delta-3}} \leq \frac{L_{31}(k)}{k^{(2+\alpha)\delta-3+(1-\varepsilon^2)(\delta/\alpha-1)},$$

which can be made small by taking k large, thanks to (4.51). In the end, we get that ρ_1 is bounded from above by a small constant for k large.

As far as ρ_2 is concerned, we split the right-hand side of (4.10) to

$$(4.58) \quad T_5 + T_6 = \sum_{i=0}^{k-1} \sum_{j=i+1}^{k-1} \mathcal{A}_{i,j} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}} + \sum_{i=0}^{k-1} \sum_{j=0}^i \mathcal{A}_{i,j} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}}.$$

Recall the definition of ℓ_i in Proposition 4.2, and define $\bar{\ell}_i = i^{(1+\varepsilon^4)/\alpha} \gg \ell_i$. We split T_5 as

$$(4.59) \quad T_{5a} + T_{5b} = \left(\sum_{i=0}^{k/2} + \sum_{i=k/2+1}^{k-1} \right) \sum_{j=(i+\bar{\ell}_i) \wedge (k-1)}^{k-1} \mathcal{A}_{i,j} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}}.$$

and

$$(4.60) \quad T_{5c} + T_{5d} = \left(\sum_{i=0}^{k^{1-\varepsilon^2}} + \sum_{i=k^{1-\varepsilon^2}+1}^{k-1} \right) \sum_{j=i+1}^{(i+\bar{\ell}_i)\wedge(k-1)} \mathcal{A}_{i,j} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}}.$$

From (4.11) and Theorem A.5, we get that (using that $(2 + \alpha)\delta - 2 > 1$ for the second line)

$$(4.61) \quad \begin{aligned} T_{5a} &\leq \sum_{i=0}^{k/2} \left(\sum_{j=i+\bar{\ell}_i}^{3k/4} + \sum_{j=3k/4+1}^{k-1} \right) \frac{C_{12} i^\delta L(j-i)^\delta}{(j-i)^{(1+\alpha)\delta}} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}} \\ &\leq \frac{L_{32}(k)}{k^{(2+\alpha)\delta-2}} \sum_{i=1}^{k/2} i^\delta \frac{L_{33}(\bar{\ell}_i)}{(\bar{\ell}_i)^{(1+\alpha)\delta-1}} + L_{34}(k) k^{\delta+1} \frac{1}{k^{(1+\alpha)\delta}} \\ &\leq \frac{L_{35}(k) k^{1+\delta}}{k^{(2+\alpha)\delta-2+\frac{1+\varepsilon^4}{\alpha}((1+\alpha)\delta-1)}} + \frac{L_{34}(k)}{k^{\delta\alpha-1}} \leq \frac{L_{35}(k)}{k^{\frac{1+\varepsilon^4}{\alpha}((1+\alpha)\delta-1)-\delta}} + \frac{L_{34}(k)}{k^{\delta\alpha-1}}, \end{aligned}$$

and also (using also here that $(2 + \alpha)\delta - 2 > 1$ for the third line)

$$(4.62) \quad \begin{aligned} T_{5b} &\leq \sum_{i=k/2+1}^{k-1} \sum_{j=(i+c\bar{\ell}_k)\wedge(k-1)}^{k-1} \frac{C_{12} i^\delta L(j-i)^\delta}{(j-i)^{(1+\alpha)\delta}} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}} \\ &\leq C_{12} k^\delta \sum_{j=k/2}^{k-1} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}} \sum_{x \geq c\bar{\ell}_k} \frac{L(x)}{x^{(1+\alpha)\delta}} \\ &\leq L_{36}(k) k^\delta (\bar{\ell}_k)^{1-(1+\alpha)\delta} = \frac{L_{36}(k)}{k^{\frac{1+\varepsilon^4}{\alpha}((1+\alpha)\delta-1)-\delta}}. \end{aligned}$$

Then the condition (4.52) guarantees that T_{5a}, T_{5b} can be made arbitrarily small by choosing k large.

Using (4.11) and Theorem A.5, we get

$$(4.63) \quad \begin{aligned} T_{5c} &\leq \sum_{i=0}^{k^{1-\varepsilon^2}} \sum_{j=i+1}^{(i+\bar{\ell}_i)\wedge(k-1)} \frac{C_{13}}{a_i^\delta} \frac{L_2(k-j)}{(k-j)^{(2+\alpha)\delta-2}} \leq \frac{L_{37}(k)}{k^{(2+\alpha)\delta-2}} \sum_{i=0}^{k^{1-\varepsilon^2}} L_{38}(i) i^{(1+\varepsilon^4)/\alpha-\delta/\alpha} \\ &\leq \frac{L_{39}(k)}{k^{(2+\alpha)\delta-2}} k^{(1-\varepsilon^2)((1+\varepsilon^4-\delta)/\alpha+1)}. \end{aligned}$$

Again, (4.51) insures that T_{5c} can be made arbitrarily small by choosing k large.

Finally, it remains to bound T_{5d} . As for (4.49), there are at most $\bar{\ell}_k$ terms in the sum over i (and $(2 + \alpha)\delta - 2 > 1$), so that

$$(4.64) \quad T_{5d} \leq C_{14} \bar{\ell}_k \max_{\substack{k^{1-\varepsilon^2} \leq i < k-1 \\ i \leq j \leq i+\bar{\ell}_i}} \mathcal{A}_{i,j}.$$

Then we use Proposition 4.2, and more precisely (4.16), to get that

$$(4.65) \quad T_{5d} \leq L_{40}(k) k^{(1+\varepsilon^4)/\alpha} k^{-\delta(1+\varepsilon^2/2)/\alpha},$$

which can be made arbitrarily small by choosing k large, because of condition (4.52).

For T_6 , we have

$$(4.66) \quad T_6 \leq \sum_{j=0}^{k-1} \sum_{i=j}^{k-1} \mathcal{A}_{i,j} \frac{L_2(k-i)}{(k-i)^{(2+\alpha)\delta-2}}.$$

By following the same procedure adopted for T_5 , T_6 is bounded above by a small term when k is large. The proof of (1.13) in the case $\alpha \in (1, 2]$ is therefore complete and, with it, also the proof of the lower bound part of Theorem 1.4.

Appendix A. Bivariate renewal theory, important estimates

We present here some results on the bivariate renewal process τ defined in Section 1.1, and in particular Proposition A.3 which gives some conditions on the transience/recurrence of the intersection renewal $\sigma = \tau \cap \tau'$. Recall the notations of Section 1.4 for the recentering sequence $(b_n)_{n \geq 0}$ and for the scaling sequence $(a_n)_{n \geq 0}$.

A.1. Local large deviations and a useful Lemma

We first present some local large deviation estimate, which is used in the proof of Lemma 4.4.

THEOREM A.1 ([Ber19b, Theorem 2.4]). — *Assume that $\mu < +\infty$. We have that there exists a constant $C > 0$ such that uniformly for (n, m) such that $n - \mu k \geq a_k \wedge (C\sqrt{k \log k})$,*

$$\mathbf{P}(\tau_k = (n, m)) \leq C k K(|n - \mu k| + |m - \mu k|).$$

We give another useful lemma, that controls the number of renewals in $(0, N] \times (0, M]$, in the case $\alpha \in (0, 1)$.

LEMMA A.2. — *Assume $\alpha \in (0, 1)$. Given $\delta > 0$ there exists $\varepsilon > 0$ such that for N sufficiently large and $M \sim \gamma N$ we have*

$$(A.1) \quad \mathbf{P}\left(|\tau \cap (0, N] \times (0, M]| \geq \varepsilon N^\alpha / L(N)\right) \geq 1 - \delta.$$

Proof. — Set $n = n(\varepsilon, N) = \varepsilon N^\alpha / L(N)$ and $B_{N,M} := (0, N] \times (0, M]$. We want to prove

$$(A.2) \quad \mathbf{P}(\tau_n \notin B_{N,M}) \leq \delta.$$

Let us define $\tilde{\tau}_n = (\tilde{\tau}_n^{(1)}, \tilde{\tau}_n^{(2)})$ with

$$(A.3) \quad \tilde{\tau}_n^{(1)} := \sum_{i=1}^n (\tau_i^{(1)} - \tau_{i-1}^{(1)}) \mathbf{1}_{\{\tau_i^{(1)} - \tau_{i-1}^{(1)} \leq N\}}, \quad \tilde{\tau}_n^{(2)} := \sum_{i=1}^n (\tau_i^{(2)} - \tau_{i-1}^{(2)}) \mathbf{1}_{\{\tau_i^{(2)} - \tau_{i-1}^{(2)} \leq M\}}.$$

Then we have

$$(A.4) \quad \mathbf{P}(\tau_n \notin B_{N,M}) \leq \mathbf{P}(\tilde{\tau}_n \notin B_{N,M}) + \mathbf{P}(\exists i \leq n; (\tau_i - \tau_{i-1}) \notin B_{N,M}).$$

Note that the marginals $\tau_n^{(1)}$ and $\tau_n^{(2)}$ have the same distribution: as $N \rightarrow \infty$ we have

$$(A.5) \quad \mathbf{P}(\tau_1^{(1)} = N) = \mathbf{P}(\tau_1^{(2)} = N) \sim \frac{1}{(1 + \alpha)} L(N) N^{-(1+\alpha)}.$$

Therefore, we can bound the second term in (A.4) by

$$(A.6) \quad n \left(\mathbf{P}(\tau_1^{(1)} > N) + \mathbf{P}(\tau_1^{(2)} > M) \right) \leq n \left(C_\alpha L(N) N^{-\alpha} + C_\alpha L(M) M^{-\alpha} \right) \leq \varepsilon C_{\alpha, \gamma},$$

which is smaller than $\delta/2$ if $\varepsilon \leq \delta/(2C_{\alpha, \gamma})$.

The first term in (A.4) is bounded by $\mathbf{P}(\tilde{\tau}_n^{(1)} > N) + \mathbf{P}(\tilde{\tau}_n^{(2)} > M)$. Observe that for every choice of $\lambda_1 > 0$ we have

$$(A.7) \quad \mathbf{P}(\tilde{\tau}_n^{(1)} > N) \leq e^{-\lambda_1 N} \mathbf{E} \left[e^{\lambda_1 \tilde{\tau}_n^{(1)}} \right] \leq e^{n \log \mathbf{E} [e^{\lambda_1 \tilde{\tau}_1^{(1)}}] - \lambda_1 N}.$$

Using the fact that $\tilde{\tau}_1^{(1)} = \tau_1^{(1)} \mathbf{1}_{\{\tau_1^{(1)} \leq N\}} \leq N$, we get that for any $s \geq 1$

$$(A.8) \quad \mathbf{E}[(\tilde{\tau}_1^{(1)})^s] \leq N^{s-1} \mathbf{E} \left[\tau_1^{(1)} \mathbf{1}_{\{\tau_1^{(1)} \leq N\}} \right] \leq c_\alpha L(N) N^{-\alpha} N^s,$$

where we used that $\mathbf{P}(\tau_1^{(1)} = N) \leq \text{cst.} L(N) N^{-(1+\alpha)}$ with $\alpha \in (0, 1)$ to estimate the second expectation. In the end, expanding the exponential and using the above bound, we get

$$(A.9) \quad \log \mathbf{E} \left[e^{\lambda_1 \tilde{\tau}_1^{(1)}} \right] \leq \log \left(1 + c_\alpha L(N) N^{-\alpha} (e^{\lambda_1 N} - 1) \right) \leq c_\alpha L(N) N^{-\alpha} e^{\lambda_1 N}.$$

We pick k_0 such that $\delta^{k_0-1} \leq e^{-1}/4$, and choose $\lambda_1 = k_0 N^{-1} \log(1/\delta)$, then with the definition of $n = \varepsilon L(N)^{-1} N^\alpha$, we get $n \log \mathbf{E} [e^{\lambda_1 \tilde{\tau}_1^{(1)}}] \leq c_\alpha \varepsilon \delta^{-k_0}$, and choosing $\varepsilon \leq c_\alpha^{-1} \delta^{k_0}$ we get from (A.7)

$$(A.10) \quad \mathbf{P}(\tilde{\tau}_n^{(1)} > N) \leq e \delta^{k_0} \leq \delta/4.$$

Using the same reasoning and choosing $\lambda_2 = k_0 M^{-1} \log(1/\delta)$, we have that if $\varepsilon \leq c_{\alpha, \gamma}^{-1} \delta^{k_0}$ (for some constant $c_{\alpha, \gamma}$),

$$(A.11) \quad \mathbf{P}(\tilde{\tau}_n^{(2)} > M) \leq e \delta^{k_0} \leq \delta/4.$$

The proof is therefore complete by taking $\varepsilon = \min\{\delta/(2C_{\alpha, \gamma}), c_\alpha^{-1} \delta^{k_0}, c_{\alpha, \gamma}^{-1} \delta^{k_0}\}$. \square

A.2. Renewal theorems, and the intersection of two independent copies

The goal of this section is to estimate the mean overlap of two copies τ and τ' in the region $(0, N] \times (0, M]$. We leave aside the case $\alpha = 1$ which is more technical (in particular if $\mu = +\infty$): we refer to Remark A.7 for more comments on this case. We define

$$(A.12) \quad U_{N, M} := \mathbf{E} [|\sigma \cap ([0, N] \times [0, M])|] = \sum_{n=0}^N \sum_{m=0}^M \mathbf{P}((n, m) \in \tau)^2,$$

and for any $\lambda > 0$

$$(A.13) \quad \hat{U}(\lambda) := \sum_{n, m=0}^{+\infty} e^{-\lambda(n+m)} \mathbf{P}((n, m) \in \tau)^2.$$

PROPOSITION A.3. — *If $\alpha < 1$, then $\sup_{N,M \in \mathbb{N}} U_{N,M} < +\infty$. If $\alpha > 1$, then set $\rho := 1 - \min(\alpha, 2)^{-1} \in [0, 1/2]$. We have*

$$(A.14) \quad \sum_{n=1}^N \frac{1}{a_n} \sim \varphi(N)N^\rho \rightarrow +\infty \quad \text{as } N \rightarrow \infty,$$

for some slowly varying function $\varphi(\cdot)$. Moreover,

$$(A.15) \quad \begin{aligned} U_{N,N} &\sim 2c_\alpha \varphi(N)N^\rho && \text{as } N \rightarrow \infty, \\ \widehat{U}(\lambda) &\sim \frac{2^{1-\rho}c_\alpha}{\Gamma(1+\rho)}\varphi(1/\lambda)\lambda^{-\rho} && \text{as } \lambda \downarrow 0, \end{aligned}$$

with $c_\alpha = \int_0^\infty c_\alpha(t)^2 dt$, $c_\alpha(t)$ being the constant appearing in Theorem A.5. As a consequence, $\sigma = \tau \cap \tau'$ is terminating if $\alpha < 1$, and persistent if $\alpha > 1$.

This proposition is based on renewal theorems (see Theorems A.4–A.5 below), that can be found in [Ber19b] (in a more general setting), giving sharp asymptotics along the favorite direction, and general upper bounds away from it. The case $\alpha = 1$ can also be found in [Ber19b] but we do not include it here, see Remark A.7.

THEOREM A.4. — *If $\alpha \in (0, 1)$, then for $n \rightarrow +\infty$ and r such that $r/n \rightarrow t \in \mathbb{R}_+$, we have*

$$(A.16) \quad \mathbf{P}((n, n+r) \in \tau) \stackrel{n \rightarrow \infty}{\sim} C_\alpha(t)L(n)^{-1}n^{-(2-\alpha)},$$

with $C_\alpha(t) := \alpha \int_0^{+\infty} x^{1-\alpha} g_\alpha(x, (1+t)x) dx$. Moreover, for any $\delta > 0$ there is a constant $C_\delta > 0$ such that for any $r \geq n$,

$$(A.17) \quad \mathbf{P}((n, n+r) \in \tau) \leq C_\delta L(n)^{-1}n^{-(2-\alpha)} \times \left(\frac{r}{n}\right)^{-(1+\alpha)+\delta}.$$

THEOREM A.5. — *If $\alpha > 1$, for $n \rightarrow \infty$ and r such that $r/a_n \rightarrow t \in \mathbb{R}_+$, we have that*

$$(A.18) \quad \mathbf{P}((n, n+r) \in \tau) \stackrel{n \rightarrow \infty}{\sim} c_\alpha(t) \frac{1}{a_n},$$

where $c_\alpha(t) = \mu_\alpha \int_{-\infty}^{+\infty} g_\alpha(x, x + \mu_\alpha t) dx$ with $\mu_\alpha := \mu^{1/\min(\alpha, 2)}$ and g_α the density of the limiting distribution of $(\tau_n - \mu n)/a_n$ (α -stable if $\alpha \in (1, 2)$ or normal if $\alpha \geq 2$). Moreover, for any $\delta > 0$ there exists a constant $C_\delta > 0$ such that for any $r \geq a_n$,

$$(A.19) \quad \mathbf{P}((n, n+r) \in \tau) \leq \frac{C}{a_n} \left(\frac{r}{a_n}\right)^{-(1+\alpha)+\delta}.$$

Theorems A.4 and A.5 are extracted from [Ber19b, Theorems 3.1, 4.1 and Theorems 3.3, 4.2 respectively]: we refer to [Ber19b, Equations (3.4), (4.2), (3.7), (4.5)] respectively, for a statement in the symmetric setting we are considering here.

Proof of Proposition A.3.

Case $\alpha < 1$. — Notice that by symmetry, for $M \geq N$,

$$U_{N,M} \leq U_{M,M} \leq 2 \sum_{n=1}^M \sum_{r=0}^{M-n} \mathbf{P}((n, n+r) \in \tau)^2.$$

We therefore need to control the last sum. Let us denote

$$W_n := \sum_{r \geq 0} \mathbf{P}((n, n+r) \in \tau)^2.$$

Using Theorem A.4 (and properties of slowly varying functions), we get that there is a constant c such that for all $n \geq 1$

$$W_n \leq c \sum_{r=1}^n L(n)^{-2} n^{-2(2-\alpha)} + c \sum_{r \geq n} L(n)^{-2} n^{-2(2-\alpha)} (r/n)^{-2} \leq C' L(n)^{-2} n^{2\alpha-3},$$

where we used that $\sum_{r \geq n} (r/n)^{-2} \sim n \int_1^\infty x^{-2}$ as $n \rightarrow \infty$. Therefore, since $\alpha < 1$, we get

$$\sup_{N,M} U_{N,M} \leq \sum_{n=1}^{+\infty} W_n < +\infty.$$

Case $\alpha > 1$. — First of all, it is immediate that $\sum_{n=1}^N 1/a_n$ diverges as a regularly varying function with exponent ρ , since $a_n \sim \psi(n)n^{\frac{1}{\min(\alpha,2)}}$, see (1.22): it directly gives (A.14).

We now prove (A.15). We fix $\varepsilon > 0$, and denote, in complement to the definition of W_n above

$$W_n^{(\varepsilon)} := \sum_{r=0}^{\lfloor \frac{1}{\varepsilon} a_n \rfloor} \mathbf{P}((n, n+r) \in \tau)^2.$$

As a preliminary, we show that there exists some n_ε such that, provided that $n \geq n_\varepsilon$

$$(A.20) \quad (1 - \varepsilon) \frac{c_\alpha}{a_n} \leq W_n^{(\varepsilon)} \leq W_n \leq (1 + \varepsilon) \frac{c_\alpha}{a_n}.$$

Note that we also have that $W_n \leq \sum_{r=0}^{+\infty} \mathbf{P}((n, n+r) \in \tau) \leq 1$ for any n , since there is at most one renewal in the column $\{n\} \times \mathbb{N}$.

To prove (A.20), we use Theorem A.5 to get that uniformly for $0 \leq r \leq \frac{1}{\varepsilon} a_n$, we have $\mathbf{P}((n, n+r) \in \tau)^2 \sim (a_n)^{-2} c_\alpha (r/a_n)^2$ as $n \rightarrow \infty$. Hence, provided that n is large enough, we get that

$$a_n W_n^{(\varepsilon)} \geq (1 - \varepsilon^2) \frac{1}{a_n} \sum_{r=0}^{\lfloor \frac{1}{\varepsilon} a_n \rfloor} c_\alpha (r/a_n)^2 \geq (1 - 2\varepsilon^2) \int_0^{1/\varepsilon} c_\alpha(t)^2 dt,$$

the last inequality holding by Riemann–sum approximation. Note that a similar upper bound, with $1 - 2\varepsilon^2$ replaced with $1 + 2\varepsilon^2$ holds. Now, thanks to (A.19) (and since $1 + \alpha - \delta \geq 3/2$), there exists a constant $c > 0$ such that

$$a_n(W_n - W_n^{(\varepsilon)}) = a_n \sum_{r > \frac{1}{\varepsilon} a_n}^{+\infty} \mathbf{P}((n, n+k) \in \tau)^2 \leq c \frac{1}{a_n} \sum_{r > \frac{1}{\varepsilon} a_n} \left(\frac{r}{a_n}\right)^{-3} \leq c' \varepsilon^2,$$

where the last inequality also comes from a Riemann–sum approximation. Finally, note that $c_\alpha - \int_0^{1/\varepsilon} c_\alpha(t)^2 dt$ is positive, and thanks to (A.19), it is smaller than $\int_{1/\varepsilon}^\infty ct^{-3} dt \leq c'' \varepsilon^2$. In the end, we get that, provided that n is large enough,

$$(A.21) \quad (1 - 2\varepsilon^2)(c_\alpha - c'' \varepsilon^2) \leq a_n W_n^{(\varepsilon)} \leq a_n W_n \leq (1 + 2\varepsilon^2)c_\alpha + c' \varepsilon^2,$$

which gives (A.20) provided that ε has been fixed small enough.

We are now ready to estimate $U_{N,N}$. We write

$$U_{N,N} = 2 \sum_{n=0}^N \sum_{r=0}^{N-n} \mathbf{P}((n, n+r) \in \tau)^2 - \sum_{n=0}^N \mathbf{P}((n, n) \in \tau)^2.$$

The second sum is negligible compared to $\sum_{n=1}^N \frac{1}{a_n} \sim \varphi(N)N^\rho$, since $\mathbf{P}((n, n) \in \tau)^2 \sim c(a_n)^{-2}$, with $a_n \rightarrow +\infty$. We therefore focus on the first sum.

An upper bound is simply

$$\sum_{n=0}^N \sum_{r=0}^{N-n} \mathbf{P}((n, n+r) \in \tau)^2 \leq \sum_{n=0}^N W_n,$$

and since we have that $W_n \sim c_\alpha/a_n$ together with (A.14), we get that for n large enough

$$\sum_{n=0}^N \sum_{r=0}^{N-n} \mathbf{P}((n, n+r) \in \tau)^2 \leq (1+2\varepsilon)c_\alpha \sum_{n=1}^N \frac{1}{a_n} \leq (1+3\varepsilon)\varphi(N)N^\rho.$$

For a lower bound, because $a_N \leq \varepsilon N$ provided that N is large enough, we have

$$\sum_{n=0}^N \sum_{k=0}^{N-n} \mathbf{P}((n, n+k) \in \tau)^2 \geq \sum_{n=n_\varepsilon}^{(1-\varepsilon)N} W_n^{(\varepsilon)} \geq (1-2\varepsilon)c_\alpha \sum_{n=1}^{(1-\varepsilon)N} \frac{1}{a_n} \geq (1-c\varepsilon)c_\alpha\varphi(N)N^\rho,$$

where we used the lower bound (A.20) valid for n large enough, together with (A.14) for the last inequality.

We now turn to estimating $\widehat{U}(\lambda)$ as $\lambda \downarrow 0$. By symmetry, we can write that

$$\widehat{U}(\lambda) = 2 \sum_{n=0}^{+\infty} e^{-2\lambda n} \sum_{r=0}^{+\infty} e^{-\lambda r} \mathbf{P}((n, n+r) \in \tau)^2 - \sum_{n=0}^{+\infty} e^{-2\lambda n} \mathbf{P}((n, n) \in \tau)^2.$$

The second term is negligible compared to $\varphi(1/\lambda)\lambda^{-\rho}$ as $\lambda \downarrow 0$, since $\sum_{n=0}^N \mathbf{P}((n, n) \in \tau)^2$ is negligible compared to $\varphi(N)N^\rho$, by standard properties of Laplace transforms, and we again focus on the first term.

First of all, an upper bound is

$$\sum_{n=0}^{+\infty} e^{-2\lambda n} \sum_{r=0}^{+\infty} e^{-\lambda r} \mathbf{P}((n, n+r) \in \tau)^2 \leq \sum_{n=0}^{+\infty} e^{-2\lambda n} W_n.$$

Since $\sum_{n=0}^N W_n \sim c_\alpha\varphi(N)N^\rho$, we get by standard properties of Laplace transforms (see [BGT87, Corollary 1.7.3]) that

$$\sum_{n=0}^{+\infty} e^{-2\lambda n} W_n \sim \frac{c_\alpha}{\Gamma(1+\rho)} \varphi(1/2\lambda)(2\lambda)^{-\rho} \quad \text{as } \lambda \downarrow 0.$$

For a lower bound, we get that

$$\sum_{n=0}^{+\infty} e^{-2\lambda n} \sum_{r=0}^{+\infty} e^{-\lambda r} \mathbf{P}((n, n+r) \in \tau)^2 \geq \sum_{n=0}^{+\infty} e^{-2\lambda n} e^{-\lambda a_n/\varepsilon} W_n^{(\varepsilon)}.$$

Now, we use that there is some n_ε such that for $n \geq n_\varepsilon$ we have that $W_n^{(\varepsilon)} \geq (1 - \varepsilon)c_\alpha/a_n$ (see (A.20)), and that $a_n/\varepsilon \leq \varepsilon n$. We therefore get that

$$\begin{aligned} & \sum_{n=0}^{+\infty} e^{-2\lambda n} \sum_{r=0}^{+\infty} e^{-\lambda r} \mathbf{P}((n, n+r) \in \tau)^2 \\ & \geq (1 - \varepsilon)c_\alpha \sum_{n=n_\varepsilon}^{+\infty} e^{-2(1+\varepsilon)\lambda n} \frac{1}{a_n} \stackrel{\lambda \rightarrow 0}{\sim} \frac{(1 - \varepsilon)c_\alpha}{\Gamma(1 + \rho)} \varphi(1/\lambda) (2(1 + \varepsilon)\lambda)^{-\rho}, \end{aligned}$$

where we used again [BGT87, Corollary 1.7.3] for the last asymptotics.

By letting $\varepsilon \downarrow 0$, we obtain matching upper and lower bound, so that (A.15) is proved. \square

We now use Proposition A.3, and in particular the estimate of the Laplace transform $\widehat{U}(\lambda)$, to obtain estimates on the tail probability of the intersection renewal $\sigma = \tau \cap \tau'$. More precisely, we define $\underline{\sigma} := \sigma^{(1)} + \sigma^{(2)}$ and estimate $\mathbf{P}^{\otimes 2}(\underline{\sigma}_1 > N)$.

LEMMA A.6. — Assume that $\alpha > 1$. Then recalling that $\rho = 1 - \min(\alpha, 2)^{-1} \in [0, 1/2]$, we get that

$$(A.22) \quad \mathbf{P}^{\otimes 2}(\underline{\sigma}_1 > N) \stackrel{N \rightarrow \infty}{\sim} \frac{2^\rho \sin(\pi\rho)}{\pi\rho} (U_{N,N})^{-1} \stackrel{N \rightarrow \infty}{\sim} C_{\alpha,\rho} \varphi(N)^{-1} N^{-\rho}.$$

Proof. — Recall the definition of $\widehat{U}(\lambda) = \sum_{n,m \geq 0} e^{-\lambda(n+m)} \mathbf{P}^{\otimes 2}((n, m) \in \sigma)$. We also set, for any $\lambda > 0$,

$$(A.23) \quad \widehat{K}(\lambda) := \sum_{n,m \geq 1} e^{-\lambda(n+m)} \mathbf{P}^{\otimes 2}(\sigma_1 = (n, m)) = \sum_{k \geq 2} e^{-\lambda k} \mathbf{P}^{\otimes 2}(\underline{\sigma}_1 = k).$$

The key idea of this proof is the following identity

$$(A.24) \quad \widehat{U}(\lambda) = 1 + \widehat{K}(\lambda)\widehat{U}(\lambda) \iff 1 - \widehat{K}(\lambda) = \frac{1}{\widehat{U}(\lambda)},$$

which is obtained from the identity

$$(A.25) \quad \mathbf{P}^{\otimes 2}((n, m) \in \sigma) = \mathbf{1}_{\{n=m=0\}} + \sum_{i=1}^n \sum_{j=1}^m \mathbf{P}^{\otimes 2}(\sigma_1 = (i, j)) \mathbf{P}^{\otimes 2}((n-i, n-j) \in \sigma).$$

Now, since we know the behavior of $\widehat{U}(\lambda)$ as $\lambda \downarrow 0$, we get the behavior of $\widehat{K}(\lambda)$, from which we should be able to infer that of $\mathbf{P}^{\otimes 2}(\underline{\sigma}_1 > N)$. Let us develop here how we proceed: we use [BGT87, Corollary 1.7.3 and Theorem 8.7.3]. We can view $\underline{\sigma}$ as a renewal process with inter-arrival distribution $\mathbf{P}^{\otimes 2}(\underline{\sigma}_1 = k) = \mathbf{P}^{\otimes 2}(\sigma_1^{(1)} + \sigma_1^{(2)} = k)$, and we set $u_n := \mathbf{P}^{\otimes 2}(n \in \underline{\sigma})$ its renewal mass function, so we have $\widehat{U}(\lambda) = \sum_{n=0}^\infty e^{-\lambda n} u_n$ (and (A.24) is standard from the one-dimensional renewal equation). Now, [BGT87, Corollary 1.7.3] tells that since $\widehat{U}(\lambda)$ is regularly varying with exponent $-\rho$ (recall $\rho = 1 - \min(\alpha, 2)^{-1}$), we have that $\sum_{n=0}^N u_n \sim \Gamma(1 + \rho)\widehat{U}(1/N) \sim 2^{-\rho}U_{N,N}$ (where we used (A.15)). In turn [BGT87, Theorem 8.7.3] gives that

$$\mathbf{P}^{\otimes 2}(\underline{\sigma}_1 > N) \stackrel{N \rightarrow \infty}{\sim} \frac{(2^{-\rho}U_{N,N})^{-1}}{\Gamma(1 + \rho)\Gamma(1 - \rho)},$$

and we are done. \square

Remark A.7. — The case $\alpha = 1$ has been left aside, mostly to avoid technicalities. Denote $\mu(n) := \mathbf{E}[\min(\tau_1^{(1)}, n)]$ the truncated first moment of $\tau_1^{(1)}$. It is shown in [Ber19b, Theorem 3.4] (or (3.11) in the symmetric context) that along the favorite direction, for $n \rightarrow \infty$ and r with $r/a_{n/\mu(n)} \rightarrow t \in \mathbb{R}_+$, we have

$$(A.26) \quad \mathbf{P}((n, n+r) \in \tau) \sim \frac{c_1(t)}{\mu(n)a_{n/\mu(n)}},$$

with $c_1(t) := \int_{-\infty}^{+\infty} g_\alpha(x, (1+t)x) dx$. Notice that $n/\mu(n)$ is the typical number of steps to reach distance n . Again, estimates away from the favorite direction are provided in [Ber19b, Theorem 4.2] (or (4.6) in the symmetric case): for any $\delta > 0$, there is a constant C_δ such that for any $r \geq a_{n/\mu(n)}$,

$$(A.27) \quad \mathbf{P}((n, n+r) \in \tau) \leq \frac{C_\delta}{\mu(n)a_{n/\mu(n)}} \left(\frac{r}{\mu(n)a_{n/\mu(n)}} \right)^{-2+\delta}.$$

This shows that the main contribution to $U_{N,N}$ comes also here from the terms close to the diagonal, that is

$$U_{N,N} \asymp 2 \sum_{n=1}^N \sum_{r=0}^{a_{n/\mu(n)}} \mathbf{P}((n, n+r) \in \tau)^2 \asymp \sum_{n=1}^N \frac{1}{\mu(n)^2 a_{n/\mu(n)}}.$$

(We denoted $x_n \asymp y_n$ if x_n/y_n is bounded away from 0 and $+\infty$.) Let us stress that we have $\mu(n) \sim \mu(a_{n/\mu(n)})$ (this comes from [Ber19a, Lemma 4.3]): by a change of variable $x = n/\mu(n)$ (comparing the sum to an integral, and considering $\mu(n), a_n$ as functions of positive real numbers), we get that

$$U_{N,N} \asymp \int_1^{N/\mu(N)} \frac{dx}{a_x \mu(a_x)} \asymp \int_1^{a_{N/\mu(N)}} \frac{du}{uL(u)\mu(u)},$$

where we used another change of variables $u = a_x$ ($dx \sim L(u)^{-1} du$, since $n \sim a_n/L(a_n)$). As a conclusion, we expect to have the following criterion:

$$\sigma = \tau \cap \tau' \text{ is persistent} \iff \sum_{n \geq 1} \frac{1}{a_n \mu(a_n)} = +\infty \iff \sum_{n \geq 1} \frac{1}{nL(n)\mu(n)} = +\infty.$$

As an example, if $L(n) = (\log n)^\kappa$ with $\kappa \geq -1$, then $\mu(n) \sim c_\kappa \max(1, (\log n)^{1+\kappa})$ and hence $\sigma = \tau \cap \tau'$ should be persistent if and only if $\kappa > 0$.

BIBLIOGRAPHY

[AB18] Kenneth S. Alexander and Quentin Berger, *Pinning of a renewal on a quenched renewal*, Electron. J. Probab. **23** (2018), 6. †309

[Ale08] Kenneth S. Alexander, *The effect of disorder on polymer depinning transitions*, Commun. Math. Phys. **279** (2008), no. 1, 117–146. †302

[AZ09] Kenneth S. Alexander and Nikos Zygouras, *Quenched and annealed critical points in polymer pinning models*, Commun. Math. Phys. **291** (2009), no. 3, 659–689. †302

[BB01] Aleksandr Alekseevich Borovkov and Konstantin A. Borovkov, *On probabilities of large deviations for random walks. I. Regularly varying distribution tails*, Theory Probab. Appl. **46** (2001), no. 2, 193–213. †325

- [BBB⁺99] R. D. Blake, J. W. Bizzaro, J. D. Blake, G. R. Day, Scott G. Delcourt, J. Knowles, K. A. Marx, and J. Jr SantaLucia, *Statistical mechanical simulation of polymeric DNA melting with MELTSIM*, *Bioinformatics* **15** (1999), no. 5, 370–375. ↑300
- [BD98] R. D. Blake and Scott G. Delcourt, *Thermal stability of DNA*, *Nucleic Acids Res.* **26** (1998), no. 14, 3323–3332. ↑300
- [Ber19a] Quentin Berger, *Notes on random walks in the Cauchy domain of attraction*, *Probab. Theory Relat. Fields* **175** (2019), no. 1-2, 1–44. ↑336
- [Ber19b] ———, *Strong renewal theorems and local large deviations for multivariate random walks and renewals*, *Electron. J. Probab.* **24** (2019), 46. ↑330, 332, 336
- [BGK18] Quentin Berger, Giambattista Giacomin, and Maha Khatib, *Dna melting structures in the generalized Poland–Scheraga model*, *ALEA, Lat. Am. J. Probab. Math. Stat.* **15** (2018), no. 2, 993–1025. ↑299, 300, 310
- [BGL19] Quentin Berger, Giambattista Giacomin, and Hubert Lacoin, *Disorder and critical phenomena: the $\alpha = 0$ copolymer model*, *Probab. Theory Relat. Fields* **174** (2019), no. 3-4, 787–819. ↑302
- [BGT87] Nicolas Hugh Bingham, Charles M. Goldie, and Jozef L. Teugels, *Regular variations*, *Encyclopedia of Mathematics and Its Applications*, vol. 27, Cambridge University Press, 1987. ↑334, 335
- [BH02] Ralf Bundschuh and Terence Hwa, *Statistical mechanics of secondary structures formed by random RNA sequences*, *Phys. Rev. E* **65** (2002), no. 3, 031903. ↑309
- [BL11] Quentin Berger and Hubert Lacoin, *The effect of disorder on the free-energy for the Random Walk Pinning Model: smoothing of the phase transition and low temperature asymptotics*, *J. Stat. Phys.* **142** (2011), no. 2, 322–341. ↑307
- [BL12] ———, *Sharp critical behavior for pinning models in a random correlated environment*, *Stochastic Processes Appl.* **122** (2012), no. 4, 1397–1436. ↑309
- [BL17] ———, *The high-temperature behavior of the directed polymer in dimension $1 + 2$* , *Ann. Inst. Henri Poincaré, Probab. Stat.* **53** (2017), no. 1, 430–450. ↑306
- [BL18] ———, *Pinning on a defect line: characterization of marginal disorder relevance and sharp asymptotics for the critical point shift*, *J. Inst. Math. Jussieu* **17** (2018), no. 2, 305–346. ↑302, 306
- [BP15] Quentin Berger and Julien Poisat, *On the critical curve of the pinning and copolymer models in correlated Gaussian environment*, *Electron. J. Probab.* **20** (2015), 71. ↑309
- [BS10] Matthias Birkner and Rongfeng Sun, *Annealed vs quenched critical points for a random walk pinning model*, *Ann. Inst. Henri Poincaré, Probab. Stat.* **46** (2010), no. 2, 414–441. ↑309
- [BS11] ———, *Disorder relevance for the random walk pinning model in dimension 3*, *Ann. Inst. Henri Poincaré, Probab. Stat.* **47** (2011), no. 1, 259–293. ↑309
- [BT10] Quentin Berger and Fabio L. Toninelli, *On the critical point of the Random Walk Pinning Model in dimension $d = 3$* , *Electron. J. Probab.* **15** (2010), 21. ↑309
- [CCP19] Dimitris Cheliotis, Yuki Chino, and Julien Poisat, *The random pinning model with correlated disorder given by a renewal set*, *Ann. Henri Lebesgue* **2** (2019), 281–329. ↑309
- [CdH13] Francesco Caravenna and Frank den Hollander, *A general smoothing inequality for disordered polymers*, *Electron. Commun. Probab.* **18** (2013), 76. ↑302, 306
- [Com07] Francis Comets, *Weak disorder for low dimensional polymers: the model of stable laws*, *Markov Process. Relat. Fields* **13** (2007), no. 4, 681–696. ↑309
- [DGLT09] Bernard Derrida, Giambattista Giacomin, Hubert Lacoin, and Fabio L. Toninelli, *Fractional moment bounds and disorder relevance for pinning models*, *Commun. Math. Phys.* **287** (2009), no. 3, 867–887. ↑302, 306, 319

- [DR14] Bernard Derrida and Martin Retaux, *The depinning transition in presence of disorder: a toy model*, J. Stat. Phys. **156** (2014), no. 2, 268–290. ↑302
- [EON11] T. R. Einert, Henri Orland, and Roland R. Netz, *Secondary structure formation of homopolymeric single-stranded nucleic acids including force and loop entropy: implications for DNA hybridization*, Eur. Phys. J. E **34** (2011), no. 6, 55. ↑309
- [Fis84] Michael E. Fisher, *Walks, walls, wetting, and melting*, J. Stat. Phys. **34** (1984), no. 5-6, 667–730. ↑300
- [Gia07] Giambattista Giacomin, *Random polymer models*, Imperial College Press; World Scientific, 2007. ↑300, 301, 306, 312
- [Gia08] ———, *Renewal convergence rates and correlation decay for homogeneous pinning models*, Electron. J. Probab. **13** (2008), 18. ↑304
- [Gia11] ———, *Disorder and critical phenomena through basic probability models. École d'Été de Probabilités de Saint-Flour XL – 2010*, Lecture Notes in Mathematics, vol. 2025, Springer, 2011. ↑302, 304, 306, 307, 308, 310, 315
- [GK17] Giambattista Giacomin and Maha Khatib, *Generalized Poland–Sheraga denaturation model and two dimensional renewal processes*, Stochastic Processes Appl. **127** (2017), no. 2, 526–573. ↑299, 300, 301, 303, 304, 309
- [GLT10] Giambattista Giacomin, Hubert Lacoïn, and Fabio L. Toninelli, *Marginal relevance of disorder for pinning models*, Commun. Pure Appl. Math. **63** (2010), no. 2, 233–265. ↑302
- [GLT11] ———, *Disorder relevance at marginality and critical point shift*, Ann. Inst. Henri Poincaré, Probab. Stat. **47** (2011), no. 1, 148–175. ↑305
- [GO03] Thomas Garel and Henri Orland, *On the role of mismatches in DNA denaturation*, <https://arxiv.org/abs/cond-mat/0304080>, 2003. ↑301, 309
- [GO04] ———, *Generalized Poland–Scherraga model for DNA hybridization*, Biopolymers **75** (2004), no. 6, 453–467. ↑301, 309
- [GT06a] Giambattista Giacomin and Fabio L. Toninelli, *The localized phase of disordered copolymers with adsorption*, ALEA, Lat. Am. J. Probab. Math. Stat. **1** (2006), 149–180. ↑308
- [GT06b] ———, *Smoothing effect of quenched disorder on polymer depinning transitions*, Commun. Math. Phys. **266** (2006), no. 1, 1–16. ↑302, 306, 307
- [Kha16] Maha Khatib, *Le modèle de Poland–Scherraga généralisé/une approche de renouvellement bidimensionnel pour la dénaturation de l'ADN*, Ph.D. thesis, Université Paris Diderot, France, 2016. ↑305
- [Lac10a] Hubert Lacoïn, *The martingale approach to disorder irrelevance for pinning models*, Electron. Commun. Probab. **15** (2010), 418–427. ↑302, 314
- [Lac10b] ———, *New bounds for the free energy of directed polymers in dimension $1 + 1$ and $1 + 2$* , Commun. Math. Phys. **294** (2010), no. 2, 471–503. ↑309
- [Mar02] James B. Martin, *Linear growth for greedy lattice animals*, Stochastic Processes Appl. **98** (2002), no. 1, 43–66. ↑311
- [NG06] Richard A. Neher and Ulrich Gerland, *Intermediate phase in DNA melting*, Phys. Rev. E **73** (2006), no. 4, 030902R. ↑301, 309, 310
- [PBG⁺92] Chung-Kang Peng, Sergey V. Buldyrev, Ary L. Goldberger, Shlomo Havlin, Francesco Sciortino, M. Simons, and H. Eugene Stanley, *Long-range correlations in nucleotide sequences*, Nature **356** (1992), 168–170. ↑309
- [Pin81] Iosif Froimovich Pinelis, *A problem on large deviations in a space of trajectories*, Theory Probab. Appl. **26** (1981), 69–84. ↑325
- [PS70] Douglas Poland and Harold A. Scherraga, *Theory of helix-coil transitions in biopolymers. Statistical mechanical theory of order-disorder transitions in biological macromolecules*, Academic Press Inc., 1970. ↑300

- [SW11] Seva Shneer and Vitali I. Wachtel, *A unified approach to the heavy-traffic analysis of the maximum of random walks*, Theory Probab. Appl. **55** (2011), no. 2, 332–341. ↑325
- [TN08] Mikail V. Tamm and Serguei K. Nechaev, *Unzipping of two random heteropolymers: Ground-state energy and finite-size effects*, Phys. Rev. E **78** (2008), 011903. ↑309
- [Ton08] Fabio L. Toninelli, *A replica-coupling approach to disordered pinning models*, Commun. Math. Phys. **280** (2008), no. 2, 389–401. ↑302, 305
- [Wat12] Frédérique Watbled, *Sharp asymptotics for the free energy of 1 + 1 dimensional directed polymers in an infinitely divisible environment*, Electron. Commun. Probab. **17** (2012), 53. ↑305
- [Wei16] Ran Wei, *On the long-range directed polymer model*, J. Stat. Phys. **165** (2016), no. 2, 320–350. ↑309

Manuscript received on 31st July 2018,
revised on 6th May 2019,
accepted on 21st May 2019.

Recommended by Editor H. Lacoïn.
Published under license CC BY 4.0.



This journal is a member of Centre Mersenne.



Quentin BERGER
Sorbonne Université,
Laboratoire de Probabilités Statistique
et Modélisation, UMR 8001,
75205 Paris, France

quentin.berger@sorbonne-universite.fr

Giambattista GIACOMIN
Université de Paris,
Laboratoire de Probabilités Statistique
et Modélisation, UMR 8001,
75205 Paris, France

giambattista.giacomin@univ-paris-diderot.fr

Maha KHATIB
Université Libanaise,
Laboratoire de Mathématiques-EDST,
Beyrouth, Liban
m.khatib@ul.edu.lb