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LIEB–THIRRING INEQUALITIES FOR WAVE FUNCTIONS VANISHING ON THE DIAGONAL SET

INÉGALITÉS DE LIEB–THIRRING POUR
LES FONCTIONS D’ONDES NULLES SUR
LA DIAGONALE

ABSTRACT. — We propose a general strategy to derive Lieb–Thirring inequalities for scale-covariant quantum many-body systems. As an application, we obtain a generalization of the Lieb–Thirring inequality to wave functions vanishing on the diagonal set of the configuration space, without any statistical assumption on the particles.

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RÉSUMÉ. — Nous proposons une stratégie générale pour démontrer les inégalités de Lieb–Thirring pour des systèmes quantiques à beaucoup de particules, invariant par changement d’échelle. Comme application, nous obtenons une généralisation de l’inégalité de Lieb–Thirring pour les fonctions d’ondes qui s’annulent sur la diagonale de l’espace des configurations, sans aucune hypothèse statistique sur les particules.

1. Introduction

The celebrated Lieb–Thirring inequality states that the expected kinetic energy of a free Fermi gas is bounded from below by its semiclassical approximation up to a universal factor, namely

$$(1.1) \quad \left\langle \Psi_N, \sum_{i=1}^N (-\Delta_{\mathbf{x}_i})^s \Psi_N \right\rangle \geq K \int_{\mathbb{R}^d} \varrho_{\Psi_N}(\mathbf{x})^{1+2s/d} d\mathbf{x}.$$

Here Ψ_N is an N -particle wave function in $L^2((\mathbb{R}^d)^N)$, normalized so that $\|\Psi_N\|_{L^2(\mathbb{R}^{dN})} = 1$ and thus encoding in its squared amplitude a probability distribution for particle positions $\mathbf{x} = (\mathbf{x}_1, \dots, \mathbf{x}_N)$, $\mathbf{x}_j \in \mathbb{R}^d$, with one-body density

$$\varrho_{\Psi_N}(\mathbf{x}) := \sum_{j=1}^N \int_{\mathbb{R}^{d(N-1)}} |\Psi(\mathbf{x}_1, \dots, \mathbf{x}_{j-1}, \mathbf{x}, \mathbf{x}_{j+1}, \dots, \mathbf{x}_N)|^2 \prod_{i \neq j} d\mathbf{x}_i,$$

and, crucially, subject to the anti-symmetry

$$(1.2) \quad \begin{aligned} \Psi_N(\mathbf{x}_1, \dots, \mathbf{x}_i, \dots, \mathbf{x}_j, \dots, \mathbf{x}_N) \\ = -\Psi_N(\mathbf{x}_1, \dots, \mathbf{x}_j, \dots, \mathbf{x}_i, \dots, \mathbf{x}_N), \quad \forall i \neq j. \end{aligned}$$

This is Pauli’s exclusion principle for fermions⁽¹⁾. Replacing the minus sign in (1.2) by a plus sign defines bosonic particles, while if the particles are non-identical, i.e. distinguishable, no exchange symmetry may be imposed.

The inequality (1.1) was first proved by Lieb and Thirring in 1975 for the case $s = 1$ relevant to non-relativistic particles [LT75, LT76], and extended by Daubechies in 1983 to general $s > 0$, thus also including the relativistic case $s = 1/2$ [Dau83]. The constant $K = K(d, s) > 0$ is independent of N and Ψ_N (see [FHJN18] for the best known value of K).

The Lieb–Thirring inequality is a beautiful combination of the uncertainty and exclusion principles of quantum mechanics, and has also been very actively studied in the mathematical literature from the dual perspective of estimation of eigenvalues of one-body Schrödinger operators (see e.g. [Lap12, LS09] for reviews). Historically, the Lieb–Thirring inequality was invented to give a short, elegant proof of the stability of ordinary non-relativistic matter with Coulomb forces [LT75]. In that context it is well known that stability of the *first* kind, i.e. that the ground state energy of the Coulomb system is finite, follows easily from some sort of the uncertainty principle (e.g. Sobolev’s inequality). On the other hand, the stability of the *second* kind, that the ground state energy does not diverge faster than the number of particles, is much

⁽¹⁾Here we ignore the spin of particles for simplicity (in our analysis the effect of the spin is mathematically trivial).

more subtle: for this the fermionic nature of particles is crucial. In fact, the stability of the second kind fails for bosonic (or distinguishable) charged systems [Dys67].

Without the anti-symmetry condition (1.2), the Lieb–Thirring inequality (1.1) fails and the best one can get is the Gagliardo–Nirenberg–Sobolev inequality

$$(1.3) \quad \left\langle \Psi_N, \sum_{i=1}^N (-\Delta_{\mathbf{x}_i})^s \Psi_N \right\rangle \geqslant K N^{-2s/d} \int_{\mathbb{R}^d} \varrho_{\Psi_N}(\mathbf{x})^{1+2s/d} d\mathbf{x}$$

(see e.g. [LNP16]). The emergence of the factor $N^{-2s/d}$ can be seen by considering the bosonic trial state $\Psi_N = u^{\otimes N}$ (whose density is $\varrho_{\Psi_N}(x) = N|u(x)|^2$). This factor is small when N becomes large, making (1.3) not very useful in applications.

Note that Pauli’s exclusion principle (1.2) implies that the wave function Ψ_N vanishes on the diagonal set

$$(1.4) \quad \Delta := \{(\mathbf{x}_1, \dots, \mathbf{x}_N) \in (\mathbb{R}^d)^N : \mathbf{x}_i = \mathbf{x}_j \text{ for some } i \neq j\},$$

namely there is zero probability for two quantum particles to occupy a common single position in the configuration space.

In this paper, we want to address the following

QUESTION 1.1. — *Does the Lieb–Thirring inequality (1.1) remain valid if the anti-symmetry assumption (1.2) is replaced by the weaker condition $\Psi_N|_{\Delta} = 0$?*

We will show that the answer is *yes* if and only if $2s > d$. In fact, $2s > d$ is the optimal condition for the vanishing assumption $\Psi_N|_{\Delta} = 0$ to be non-trivial (heuristically this follows from Sobolev’s embedding $H^s(\mathbb{R}^d) \subseteq C(\mathbb{R}^d)$ for $2s > d$). The precise statement of our result and its consequences will be presented in the next section.

2. Main results

Recall that for every $s > 0$ (not necessarily an integer) the operator $(-\Delta)^s$ on $L^2(\mathbb{R}^d)$ is defined as the multiplication operator $|\mathbf{p}|^{2s}$ in Fourier space, namely

$$[(-\Delta)^s f]^\wedge(\mathbf{p}) = |\mathbf{p}|^{2s} \hat{f}(\mathbf{p}), \quad \hat{f}(\mathbf{p}) := \frac{1}{(2\pi)^{d/2}} \int_{\mathbb{R}^d} f(\mathbf{x}) e^{-i\mathbf{p} \cdot \mathbf{x}} d\mathbf{x}.$$

The associated space $H^s(\mathbb{R}^d)$ is a Hilbert space with norm

$$\|u\|_{H^s(\mathbb{R}^d)}^2 := \|u\|_{L^2(\mathbb{R}^d)}^2 + \|u\|_{\dot{H}^s(\mathbb{R}^d)}^2, \quad \|u\|_{\dot{H}^s(\mathbb{R}^d)}^2 := \langle u, (-\Delta)^s u \rangle.$$

The N -particle space $H^s(\mathbb{R}^{dN})$ is defined in the same way. Let us denote the subspace of functions vanishing on the diagonal set Δ in (1.4) by

$$\mathcal{H}^{s,N}(\mathbb{R}^d) := \overline{\left\{ \Psi_N \in C_c^\infty(\mathbb{R}^{dN}) : \Psi_N|_{\Delta} = 0 \right\}}^{H^s(\mathbb{R}^{dN})}.$$

Our main result is

THEOREM 2.1 (Lieb–Thirring inequality for wave functions vanishing on diagonals). *Let $2s > d \geq 1$. Then for every $N \geq 1$ and $\Psi_N \in \mathcal{H}^{s,N}(\mathbb{R}^d)$, with $\|\Psi_N\|_{L^2(\mathbb{R}^{dN})} = 1$, we have*

$$(2.1) \quad \left\langle \Psi_N, \sum_{i=1}^N (-\Delta_{\mathbf{x}_i})^s \Psi_N \right\rangle \geq C \int_{\mathbb{R}^d} \varrho_{\Psi_N}(\mathbf{x})^{1+2s/d} d\mathbf{x}.$$

Here $C = C(d, s) > 0$ is a universal constant independent of N and Ψ_N .

We have some immediate remarks.

- (1) The condition $2s > d$ in Theorem 2.1 is optimal. If $2s \leq d$, then

$$\mathcal{H}^{s,N}(\mathbb{R}^d) = H^s(\mathbb{R}^{dN})$$

by the relatively small size, i.e. the large codimensionality, of the diagonal set (see Appendix B) and thus the Lieb–Thirring inequality fails.

- (2) For $d = 1$ and $s = 1$, it is well known that a *symmetric* wave function which vanishes on the diagonal set is equal to an anti-symmetric wave function up to multiplication by an appropriate sign function [Gir60], and hence (2.1) reduces to the usual Lieb–Thirring inequality [LT76] in this case. However, when $d > 1$ this boson-fermion correspondence is no longer available and our result is new. Furthermore, one may consider *hard-core* bosons defined by the higher-order vanishing around diagonals

$$(2.2) \quad \mathcal{H}_0^{s,N}(\mathbb{R}^d) := \overline{\left\{ \Psi_N \in C_c^\infty(\mathbb{R}^{dN} \setminus \Delta) \right\}}^{H^s(\mathbb{R}^{dN})},$$

and subject to symmetry. For large enough order $2s > d$ there is even for $d = 1$ a non-trivial difference between these spaces, and our result assumes only the weaker vanishing conditions imposed by $\mathcal{H}^{s,N}(\mathbb{R}^d)$ (see Appendix B for some further remarks).

- (3) Theorem 2.1 verifies a conjecture in [LNP16, p. 1362] that the Lieb–Thirring inequality (2.1) holds for all wave functions in the form domain of the interaction potential

$$W_s(\mathbf{x}) := \sum_{1 \leq i < j \leq N} |\mathbf{x}_i - \mathbf{x}_j|^{-2s}, \quad \mathbf{x} = (\mathbf{x}_1, \dots, \mathbf{x}_N) \in (\mathbb{R}^d)^N.$$

In fact, we have (again, see Appendix B for details)

$$(2.3) \quad \left\{ \Psi_N \in H^s(\mathbb{R}^{dN}) : \int_{\mathbb{R}^{dN}} W_s(\mathbf{x}) |\Psi_N(\mathbf{x})|^2 d\mathbf{x} < \infty \right\} \subseteq \mathcal{H}_0^{s,N}(\mathbb{R}^d) \subseteq \mathcal{H}^{s,N}(\mathbb{R}^d),$$

by the singular nature of the potential at the diagonals. We may think of the potential W_s as defining (by Friedrichs extension) a one-parameter family of non-negative and scale-covariant (scaling homogeneously to degree $-2s$) interacting N -body Hamiltonian operators

$$H_\beta := \sum_{j=1}^N (-\Delta_{\mathbf{x}_j})^s + \beta W_s, \quad \beta \geq 0.$$

(In the case $\beta = 0$, then H_β is still defined on the quadratic form domain of W_s .) One may then ask about the best constant $C(\beta) \geq 0$ in the bound

$$(2.4) \quad \langle \Psi_N, H_\beta \Psi_N \rangle \geq C(\beta) \int_{\mathbb{R}^d} \varrho_{\Psi_N}(\mathbf{x})^{1+2s/d} d\mathbf{x},$$

i.e. a Lieb–Thirring inequality (generalized uncertainty principle) for H_β . The case $\beta > 0$ was treated in [LPS15, LNP16], while our setting here concerns the limit $\beta \rightarrow 0$ of zero-range/contact interaction. A crucial difference is the strength of the interaction term, which is of order βN^2 and thus provides a large repulsive energy for fixed $\beta > 0$, while for $\beta \ll 1/N$ it ought to be much weaker than the kinetic term. Nevertheless, for $2s > d$ the potential W_s is singular enough to impose the vanishing condition at Δ , and Theorem 2.1 thus implies that the limiting constant is positive $C(0) > 0$. Furthermore, for $2s = d$ we have $C(0) = 0$ which follows from a straightforward argument through contradiction using the density of $C_c^\infty(\mathbb{R}^{dN} \setminus \Delta)$ in $H^s(\mathbb{R}^{dN})$; again see our first remark and Appendix B. Together with the case $2s < d$ that was treated in [LNP16] and for which also $C(0) = 0$, this settles the question concerning which H_β satisfy the Lieb–Thirring inequality (2.4) with $C(\beta) > 0$.

- (4) The original proof of the Lieb–Thirring inequality [LT75, LT76] is based on the following operator bound

$$(2.5) \quad 0 \leq \gamma_{\Psi_N}^{(1)} \leq 1$$

which is a consequence of Pauli’s exclusion principle (1.2). Here $\gamma_{\Psi_N}^{(1)}$ is the one-body density matrix of Ψ_N , a trace-class operator on $L^2(\mathbb{R}^d)$ with kernel

$$\gamma_{\Psi_N}^{(1)}(\mathbf{x}; \mathbf{x}') = \sum_{j=1}^N \int_{\mathbb{R}^{d(N-1)}} \Psi_N(\mathbf{x}_1, \dots, \mathbf{x}_j = \mathbf{x}, \dots, \mathbf{x}_N) \overline{\Psi_N(\mathbf{x}_1, \dots, \mathbf{x}_j = \mathbf{x}', \dots, \mathbf{x}_N)} \prod_{k \neq j} d\mathbf{x}_k.$$

However, unlike the full anti-symmetry condition (1.2), the vanishing condition $\Psi_N|_\Delta = 0$ alone is not known to be sufficient to ensure the operator inequality (2.5), and therefore the original proof in [LT75, LT76] as well as subsequent proofs based on (2.5) (e.g. Rumin’s method [Rum11]) do not apply.

Our result is in fact more general than as previously formulated. More precisely, define for any $k \geq 2$ the diagonal set of k -particle coincidences

$$(2.6) \quad \Delta_k := \left\{ (\mathbf{x}_1, \dots, \mathbf{x}_N) \in (\mathbb{R}^d)^N : \mathbf{x}_{j_1} = \dots = \mathbf{x}_{j_k} \text{ for distinct indices } j_1, \dots, j_k \right\},$$

and the corresponding space of N -particle wave functions with a vanishing condition on Δ_k

$$\mathcal{H}_k^{s,N}(\mathbb{R}^d) := \overline{\left\{ \Psi_N \in C_c^\infty(\mathbb{R}^{dN}) : \Psi_N|_{\Delta_k} = 0 \right\}}^{H^s(\mathbb{R}^{dN})}.$$

We have

THEOREM 2.2 (Lieb–Thirring inequality for wave functions vanishing on k -diagonals). — *Let $d \geq 1$, $k \geq 2$ and $2s > d(k-1)$. Then for every $N \geq 1$ and every $\Psi_N \in \mathcal{H}_k^{s,N}(\mathbb{R}^d)$, with $\|\Psi_N\|_{L^2(\mathbb{R}^{dN})} = 1$, we have*

$$(2.7) \quad \left\langle \Psi_N, \sum_{i=1}^N (-\Delta_{\mathbf{x}_i})^s \Psi_N \right\rangle \geq C \int_{\mathbb{R}^d} \varrho_{\Psi_N}(\mathbf{x})^{1+2s/d} d\mathbf{x}.$$

Here $C = C(d, s, k) > 0$ is a universal constant independent of N and Ψ_N .

The proof of Theorem 2.2 occupies the rest of the paper. Our proof is based on a general strategy of deriving Lieb–Thirring inequalities for wave functions satisfying some partial exclusion properties, which was proposed by Lundholm and Solovej in [LS13a] and developed further in [FS12, LL18, LNP16, LPS15, LS13b, LS14, LS18, Lun18, Nam18]. We will quickly review this strategy in Section 3 for the reader’s convenience, following the simplification by Lundholm, Nam and Portmann [LNP16].

The main new ingredient is a local version of the exclusion principle using the vanishing condition on the diagonal set. In Section 4, we will discuss a very useful reduction of the desired local exclusion to simply the positivity of a local energy using the scale-covariance of the kinetic operator $(-\Delta)^s$. This step refines and generalizes a recent bootstrap argument for the energy of ideal anyons by Lundholm and Seiringer [LS18]. In Section 5, the remaining crucial fact that the local energy eventually becomes positive with increasing particle number will be settled by means of a new many-particle Poincaré inequality. Some standard and non-standard results on relevant function spaces are collected in the appendices for completeness.

We stress that our method will also work for any other deformations of the Laplacian which retain similar positivity and scale-covariance properties, including other types of point interactions as well as particles subject to intermediate statistics (ideal anyons) in one and two dimensions.

3. General strategy of deriving Lieb–Thirring inequalities

In the following we will summarize the proof of the usual Lieb–Thirring inequality (1.1) for fermionic wave functions, mainly following the simplified representation in [LNP16]. The starting point is the following obvious localization formula: if $\{\Omega\}$ is a collection of disjoint subsets of \mathbb{R}^d , then

$$(3.1) \quad (-\Delta)_{|\mathbb{R}^d}^s \geq \sum_{\Omega} (-\Delta)_{|\Omega}^s,$$

where the Neumann localization $(-\Delta)_{|\Omega}^s$ is defined via the quadratic form (Sobolev seminorm)

$$\begin{aligned} & \langle u, (-\Delta)_{|\Omega}^s u \rangle \\ &= \|u\|_{\dot{H}^s(\Omega)}^2 := \begin{cases} \sum_{|\alpha|=m} \frac{m!}{\alpha!} \int_{\Omega} |D^{\alpha} u|^2 & \text{if } s = m, \\ c_{d,\sigma} \sum_{|\alpha|=m} \frac{m!}{\alpha!} \iint_{\Omega \times \Omega} \frac{|D^{\alpha} u(\mathbf{x}) - D^{\alpha} u(\mathbf{y})|^2}{|\mathbf{x} - \mathbf{y}|^{d+2\sigma}} d\mathbf{x} d\mathbf{y} & \text{if } s = m + \sigma, \end{cases} \end{aligned}$$

for all $u \in H^s(\mathbb{R}^d)$, with $m \in \mathbb{N}_0$, $\alpha \in \mathbb{N}_0^d$ multi-indices, D^{α} corresponding derivatives, and

$$0 < \sigma < 1, \quad c_{d,\sigma} := \frac{2^{2\sigma-1} \Gamma((d+2\sigma)/2)}{\pi^{d/2} |\Gamma(-\sigma)|}.$$

Consequently, for any N -body wave function $\Psi_N \in H^s(\mathbb{R}^{dN})$ we have

$$(3.2) \quad \mathcal{E}_{\mathbb{R}^d}[\Psi_N] \geq \sum_{\Omega} \mathcal{E}_{\Omega}[\Psi_N],$$

where the expected local energy on Ω is

$$(3.3) \quad \mathcal{E}_{\Omega}[\Psi_N] := \left\langle \Psi_N, \sum_{j=1}^N (-\Delta_{\mathbf{x}_j})_{|\Omega}^s \Psi_N \right\rangle = \sum_{j=1}^N \int_{\mathbb{R}^{d(N-1)}} \|\Psi_N\|_{\dot{H}_{\mathbf{x}_j}^s(\Omega)}^2 \prod_{\ell \neq j} d\mathbf{x}_{\ell}$$

Next, we have the following three key tools [LNP16, Lemmas 8, 11, 12].

LEMMA 3.1 (Local uncertainty). — *Let $d \geq 1$ and $s > 0$. Let Ψ_N be a wave function in $H^s(\mathbb{R}^{dN})$ for arbitrary $N \geq 1$ and let Q be an arbitrary cube in \mathbb{R}^d . Then*

$$(3.4) \quad \mathcal{E}_Q[\Psi_N] \geq \frac{1}{C} \frac{\int_Q \varrho_{\Psi_N}^{1+2s/d}}{\left(\int_Q \varrho_{\Psi_N}\right)^{2s/d}} - \frac{C}{|Q|^{2s/d}} \int_Q \varrho_{\Psi_N}.$$

Hereafter, $C = C(d, s) > 0$ denotes a universal constant (independent of N , Ψ_N and Q).

Lemma 3.1 can be interpreted as a local version of the lower bound (1.3) (the negative term appears due to the lack of Dirichlet boundary condition).

LEMMA 3.2 (Local exclusion for fermions). — *Let $d \geq 1$ and $s > 0$. Let Ψ_N be a fermionic wave function in $H^s(\mathbb{R}^{dN})$ satisfying (1.2) for $N \geq 2$ and let Q be an arbitrary cube in \mathbb{R}^d . Then*

$$(3.5) \quad \mathcal{E}_Q[\Psi_N] \geq C|Q|^{-2s/d} \left[\int_Q \varrho_{\Psi_N}(\mathbf{x}) d\mathbf{x} - q \right]_+,$$

where $q := \#\{\text{multi-indices } \alpha \in \mathbb{N}_0^d : 0 \leq |\alpha| < s\}$.

In the non-relativistic case $s = 1$, Lemma 3.2 simply states that as soon as there is more than one particle on Q the energy must be strictly positive, and furthermore that it grows at least linearly with the number of particles. Such a weak formulation of the exclusion principle was used by Dyson and Lenard in their first proof of the stability of matter [DL67], while its general applicability in the above format was noted by Lundholm and Solovej in [LS13a, LS13b].

LEMMA 3.3 (Covering Lemma). — *Let $0 \leq f \in L^1(\mathbb{R}^d)$ be a function with compact support such that $\int_{\mathbb{R}^d} f \geq \Lambda > 0$. Then the support of f can be covered by a collection of disjoint cubes $\{Q\}$ in \mathbb{R}^d such that*

$$(3.6) \quad \int_Q f \leq \Lambda, \quad \forall Q$$

and

$$(3.7) \quad \sum_Q \frac{1}{|Q|^\alpha} \left(\left[\int_Q f - q \right]_+ - b \int_Q f \right) \geq 0$$

for all $\alpha > 0$ and $0 \leq q < \Lambda 2^{-d}$, where

$$b := \left(1 - \frac{2^d q}{\Lambda} \right) \frac{2^{d\alpha} - 1}{2^{d\alpha} + 2^d - 2} > 0.$$

Conclusion of (1.1).

Let q be as in Lemma 3.2 and let $\Lambda = 2^d q + 1$. If $N \leq \Lambda$, then (1.1) follows immediately from (1.3), whose proof is similar to (indeed simpler than) that of Lemma 3.1. If $N > \Lambda$, then we can apply Lemma 3.3 with $f = \varrho_{\Psi_N}$ (by standard approximation we may reduce to compact support), $\alpha = 2s/d$, and obtain a collection of disjoint cubes $\{Q\}$. Combining with (3.2), (3.4) and (3.5) we obtain

$$\begin{aligned} (\varepsilon + 1)\mathcal{E}_{\mathbb{R}^d}[\Psi_N] &\geq \varepsilon \sum_Q \left[\frac{1}{C_1} \frac{\int_Q \varrho_{\Psi_N}^{1+2s/d}}{\left(\int_Q \varrho_{\Psi_N} \right)^{2s/d}} - \frac{C_1}{|Q|^{2s/d}} \int_Q \varrho_{\Psi_N} \right] \\ &\quad + \sum_Q C_2 |Q|^{-2s/d} \left[\int_Q \varrho_{\Psi}(x) dx - q \right]_+ \\ &\geq \frac{\varepsilon}{C_1} \frac{\int_{\mathbb{R}^d} \varrho_{\Psi_N}^{1+2s/d}}{\Lambda^{2s/d}} \end{aligned}$$

for any fixed constant $\varepsilon > 0$ satisfying $\varepsilon C_1 \leq C_2 b$. Thus (1.1) holds true.

As we can see from the above strategy, the only place where the anti-symmetry (1.2) plays a role is the local exclusion bound in Lemma 3.2. Extending this result to the weaker condition $\Psi_N|_{\Delta} = 0$ is the main task of our proof below.

4. Reduction of local exclusion

In this section, we prove a very useful observation, that allows to reduce the local exclusion (3.5) to the positivity of the local energy, using the scale-covariance of the kinetic energy. This step is inspired by the recent work of Lundholm and Seiringer [LS18] on the energy of ideal anyons. We formulate it abstractly as follows:

LEMMA 4.1 (Covariant energy bound). — *Assume that to any $n \in \mathbb{N}_0$ and any cube $Q \subset \mathbb{R}^d$ there is associated a non-negative number ('energy') $E_n(Q)$ satisfying the following properties, for some constant $s > 0$:*

- (scale-covariance) $E_n(\lambda Q) = \lambda^{-2s} E_n(Q)$ for all $\lambda > 0$;
- (translation-invariance) $E_n(Q + \mathbf{x}) = E_n(Q)$ for all $\mathbf{x} \in \mathbb{R}^d$;
- (superadditivity) For any collection of disjoint cubes $\{Q_j\}_{j=1}^J$ such that their union is a cube,

$$E_n \left(\bigcup_{j=1}^J Q_j \right) \geq \min_{\{n_j\} \in \mathbb{N}_0^J \text{ s.t. } \sum_j n_j = n} \sum_{j=1}^J E_{n_j}(Q_j);$$

- (a priori positivity) There exists $q \geq 0$ such that $E_n(Q) > 0$ for all $n \geq q$.

Then there exists a constant $C > 0$ independent of n and Q such that

$$(4.1) \quad E_n(Q) \geq C|Q|^{-2s/d} n^{1+2s/d}, \quad \forall n \geq q.$$

Proof. — Note that for $q \leq n \leq N$, (4.1) holds for some $C = C_N > 0$ by the a priori positivity. The main point here is to remove the N -dependence of the constant.

Denote $E_n := E_n(Q_0)$ with $Q_0 = [0, 1]^d$. Assume by induction in N that

$$(4.2) \quad E_n \geq Cn^{1+2s/d}, \quad \forall q \leq n \leq N-1$$

with a uniform constant $C > 0$ and consider $n = N$. Split Q_0 into 2^d subcubes of half side length and obtain by the superadditivity, translation-invariance and scale-covariance

$$(4.3) \quad E_N \geq 2^{2s} \min_{\{n_j\} \text{ s.t. } \sum_j n_j = N} \sum_{j=1}^{2^d} E_{n_j}.$$

Consider a configuration $\{n_j\} \subset \mathbb{N}_0^{2^d}$ such that the minimum in (4.3) is attained. The a priori positivity $E_N > 0$ ensures that none of the n_j can be N (in the same way we deduce that $E_0 = 0$). Assume that there exist exactly M numbers $n_j < q$ with $0 \leq M \leq 2^d$. Then

$$\sum_{n_j \geq q} 1 = 2^d - M \quad \text{and} \quad \sum_{n_j \geq q} n_j = N - \sum_{n_j < q} n_j \geq N - qM.$$

Therefore, from (4.3), (4.2) and Hölder's inequality we deduce that

$$(4.4) \quad E_N \geq C2^{2s} \sum_{n_j \geq q} n_j^{1+2s/d} \geq C2^{2s} \frac{\left(\sum_{n_j \geq q} n_j \right)^{1+2s/d}}{\left(\sum_{n_j \geq q} 1 \right)^{2s/d}}$$

$$(4.5) \quad \geq CN^{1+2s/d} \frac{(1 - qMN^{-1})^{1+2s/d}}{(1 - M2^{-d})^{2s/d}}$$

with the same constant C as in (4.2). If we take

$$N \geq q2^d \left(1 + \frac{d}{2s} \right),$$

so that also $qMN^{-1} \leq 1$, then by Bernoulli's inequality

$$\left(1 - qMN^{-1} \right)^{1+d/(2s)} \geq 1 - qMN^{-1} \left(1 + \frac{d}{2s} \right) \geq 1 - M2^{-d},$$

and hence (4.4) reduces to

$$(4.6) \quad E_N \geqslant CN^{1+2s/d}$$

with the same constant C as in (4.2).

By induction we obtain (4.6) for all $N \geqslant q$, with a constant C independent of N . This is the desired bound (4.1) for the unit cube Q_0 . The result for the general cube follows from scale-covariance and translation-invariance. \square

Remark 4.2. — It is in fact also possible to allow for $E_n < 0$ for finitely many $n > 0$ in Lemma 4.1, under a small refinement of the assumption of a priori positivity. It is sufficient that there exists $q > 0$ and $c > 1$ such that for all $n \geqslant q$

$$(4.7) \quad E_n > c \frac{d}{2s} 2^{d+2s} E_-, \quad E_- := \max_{0 \leqslant n < q} (-E_n).$$

Namely, with this assumption, the bound in (4.4) may again be used for all $n_j \geqslant q$, and one obtains $E_N \geqslant CN^{1+2s/d} f(M/2^d)$, $C = \min_{q \leqslant n \leqslant N-1} E_n/n^{1+2s/d}$, where the function

$$f(x) := \frac{(1 - q2^d N^{-1}x)^{1+2s/d}}{(1-x)^{2s/d}} - \frac{E_- 2^{d+2s}}{CN^{1+2s/d}} x, \quad x \in [0, 1),$$

is strictly increasing if N is large enough.

We will apply the above general bound to the local ground-state energy among wave functions satisfying the vanishing condition on k -particle diagonals

$$(4.8) \quad E_N(\Omega) := \inf \left\{ \|\Psi_N\|_{\dot{H}^{s,N}(\Omega)}^2 : \Psi_N \in \mathcal{H}_k^{s,N}(\mathbb{R}^d), \|\Psi_N\|_{L^2(\Omega^N)} = 1 \right\},$$

where we have introduced the “completely localized” kinetic functional

$$(4.9) \quad \|\Psi_N\|_{\dot{H}^{s,N}(\Omega)}^2 := \left\langle \mathbf{1}_{\Omega^N} \Psi_N, \sum_{j=1}^N (-\Delta_{\mathbf{x}_j})_{|\Omega}^s \mathbf{1}_{\Omega^N} \Psi_N \right\rangle = \sum_{j=1}^N \int_{\Omega^{N-1}} \|\Psi_N\|_{\dot{H}_{\mathbf{x}_j}^s(\Omega)}^2 \prod_{\ell \neq j} d\mathbf{x}_\ell.$$

Note that $\|\Psi_N\|_{\dot{H}^{s,N}(\Omega)}^2$ is *different* from the functional $\mathcal{E}_\Omega[\Psi_N]$ in (3.3), and its properties will be crucial to deduce the desired local exclusion for $\mathcal{E}_\Omega[\Psi_N]$. The seminorm $\|\cdot\|_{\dot{H}^{s,N}(\Omega)}$ in general contains only some of the terms of the standard homogeneous Sobolev seminorm $\|\cdot\|_{\dot{H}^s(\Omega^N)}$; however, the corresponding norms (i.e. the seminorms plus the L^2 -norm) are actually equivalent modulo N -dependent constants, not only globally on \mathbb{R}^{dN} but also locally on Q^N (see Appendix A).

The superadditivity of the energy $E_N(\Omega)$ follows from the partitioning of the many-body space and by locality respectively non-negativity of any non-local part of the kinetic energy, i.e. (3.1). The method was also used in [LS18, Lemma 4.2] for anyons.

LEMMA 4.3 (Superadditivity of $E_n(\Omega)$). — *Let $\{\Omega_j\}_{j=1}^J$ be a collection of disjoint subsets of \mathbb{R}^d and $\Omega = \cup_j \Omega_j$. Then*

$$(4.10) \quad E_N(\Omega) \geqslant \min_{\{n_j\} \in \mathbb{N}_0^J \text{ s.t. } \sum_j n_j = N} \sum_{j=1}^J E_{n_j}(\Omega_j).$$

Proof. — For any partition $A = \{A_j\}_{j=1}^J$ of $\{1, 2, \dots, N\}$ (i.e. the A_j are disjoint subsets of $\{1, 2, \dots, N\}$ such that $\sum_j |A_j| = N$), we denote by $\mathbb{1}_A$ the characteristic function of the set

$$\left\{(\mathbf{x}_1, \dots, \mathbf{x}_N) \in (\mathbb{R}^d)^N : \mathbf{x}_i \in \Omega_j \Leftrightarrow i \in A_j, \text{ for all } i, j\right\}.$$

Using the operator bound similar to (3.1)

$$(-\Delta_{\mathbf{x}_i})_{|\Omega}^s \geq \sum_{j=1}^J (-\Delta_{\mathbf{x}_i})_{|\Omega_j}^s,$$

the partition of unity

$$(4.11) \quad \mathbb{1}_{\Omega^N} = \sum_A \mathbb{1}_A,$$

and the fact that $\mathbb{1}_A$ commutes with $(-\Delta_{\mathbf{x}_i})_{|\Omega_j}^s$, we can write for any $\Psi \in \mathcal{H}_k^{s,N}(\mathbb{R}^d)$

$$\begin{aligned} \|\Psi\|_{\dot{H}^{s,N}(\Omega)}^2 &= \left\langle \mathbb{1}_{\Omega^N} \Psi, \sum_{i=1}^N (-\Delta_{\mathbf{x}_i})_{|\Omega}^s \mathbb{1}_{\Omega^N} \Psi \right\rangle \geq \sum_{j=1}^J \left\langle \mathbb{1}_{\Omega^N} \Psi, \sum_{i=1}^N (-\Delta_{\mathbf{x}_i})_{|\Omega_j}^s \mathbb{1}_{\Omega^N} \Psi \right\rangle \\ &= \sum_{j=1}^J \sum_A \left\langle \mathbb{1}_A \Psi, \sum_{i=1}^N (-\Delta_{\mathbf{x}_i})_{|\Omega_j}^s \mathbb{1}_A \Psi \right\rangle \\ &= \sum_{j=1}^J \sum_A \int_{\mathbb{R}^{d(N-|A_j|)}} \left\| \Psi(\cdot; \mathbf{x}_{A_j^c}) \right\|_{\dot{H}^{s,|A_j|}(\Omega_j)}^2 \prod_{\ell \neq j} \left[\mathbb{1}_{\Omega_\ell^{|A_\ell|}}(\mathbf{x}_{A_\ell}) d\mathbf{x}_{A_\ell} \right]. \end{aligned}$$

Here we have introduced the shorthand notation

$$(\mathbf{x}_1, \dots, \mathbf{x}_N) = (\mathbf{x}_{A_j}; \mathbf{x}_{A_j^c}), \quad \mathbf{x}_{A_j} = (\mathbf{x}_\ell)_{\ell \in A_j} \in (\mathbb{R}^d)^{|A_j|}.$$

Since $\Psi \in \mathcal{H}_k^{s,N}(\mathbb{R}^d)$, for a.e. $\mathbf{x}_{A_j^c} \in \mathbb{R}^{d(N-|A_j|)}$ the function $\Psi(\cdot; \mathbf{x}_{A_j^c})$ is in $\mathcal{H}_k^{s,|A_j|}(\mathbb{R}^d)$, and hence

$$\begin{aligned} \left\| \Psi(\cdot; \mathbf{x}_{A_j^c}) \right\|_{\dot{H}^{s,|A_j|}(\Omega_j)}^2 &\geq E_{|A_j|}(\Omega_j) \int_{\Omega_j^{|A_j|}} \left| \Psi(\mathbf{x}_{A_j}; \mathbf{x}_{A_j^c}) \right|^2 d\mathbf{x}_{A_j} \\ &= E_{|A_j|}(\Omega_j) \int_{\mathbb{R}^{d|A_j|}} \left| \Psi(\mathbf{x}_{A_j}; \mathbf{x}_{A_j^c}) \right|^2 \mathbb{1}_{\Omega_j^{|A_j|}}(\mathbf{x}_{A_j}) d\mathbf{x}_{A_j}. \end{aligned}$$

Thus in summary

$$\begin{aligned} \|\Psi\|_{\dot{H}^{s,N}(\Omega)}^2 &\geq \sum_{j=1}^J \sum_A E_{|A_j|}(\Omega_j) \int_{\mathbb{R}^{dN}} |\Psi|^2 \prod_{\ell=1}^J \left[\mathbb{1}_{\Omega_\ell^{|A_\ell|}}(\mathbf{x}_{A_\ell}) d\mathbf{x}_{A_\ell} \right] \\ &= \sum_{j=1}^J \sum_A E_{|A_j|}(\Omega_j) \langle \Psi, \mathbb{1}_A \Psi \rangle = \sum_A \left[\sum_{j=1}^J E_{|A_j|}(\Omega_j) \right] \langle \Psi, \mathbb{1}_A \Psi \rangle \\ &\geq \left[\min_{\{n_j\} \in \mathbb{N}_0^J \text{ s.t. } \sum_j n_j = N} \sum_{j=1}^J E_{n_j}(\Omega_j) \right] \sum_A \langle \Psi, \mathbb{1}_A \Psi \rangle \\ &= \left[\min_{\{n_j\} \in \mathbb{N}_0^J \text{ s.t. } \sum_j n_j = N} \sum_{j=1}^J E_{n_j}(\Omega_j) \right] \|\Psi\|_{L^2(\Omega^N)}^2. \end{aligned}$$

Here in the last identity we have used the partition of unity (4.11) again. This implies the desired estimate (4.10). \square

Now we are ready to prove the reduction of the local exclusion.

LEMMA 4.4 (Energy positivity implies local exclusion). — *Assume that there exists a constant $q > 0$ such that for any cube $Q \subset \mathbb{R}^d$,*

$$(4.12) \quad E_N(Q) > 0, \quad \forall N \geq q.$$

Then for all $N \geq 1$ and for all wave functions $\Psi_N \in \mathcal{H}_k^{s,N}(\mathbb{R}^d)$, $\|\Psi_N\|_{L^2(\mathbb{R}^{dN})} = 1$, we have

$$(4.13) \quad \mathcal{E}_Q[\Psi_N] \geq C|Q|^{-2s/d} \left[\int_Q \varrho_{\Psi_N}(\mathbf{x}) d\mathbf{x} - q \right]_+.$$

Here $C > 0$ is a constant independent of N , Ψ_N and Q .

Proof. — Given (4.12), the energy functional $E_n(Q)$ defined in (4.8) verifies all conditions in Lemma 4.1. Therefore, there exists a constant $C > 0$ independent of n and Q such that

$$(4.14) \quad E_n(Q) \geq C|Q|^{-2s/d} n^{1+2s/d} \mathbf{1}_{\{n \geq q\}} \geq C|Q|^{-2s/d} [n - q]_+, \quad \forall n \geq 0.$$

Now we adapt the localization method in the proof of Lemma 4.3 to treat the functional $\mathcal{E}_Q[\Psi_N]$ (instead of $\|\Psi_N\|_{H^{s,N}(Q)}^2$). To be precise, for any subset B of $\{1, \dots, N\}$ we denote by $\mathbf{1}_B$ the characteristic function of the set

$$\{(\mathbf{x}_1, \dots, \mathbf{x}_N) \in (\mathbb{R}^d)^N : \mathbf{x}_i \in Q \Leftrightarrow i \in B, \text{ for all } i\}.$$

For any $\Psi_N \in \mathcal{H}_k^{s,N}(\mathbb{R}^d)$, $\|\Psi_N\|_{L^2(\mathbb{R}^{dN})} = 1$, by inserting the partition of unity

$$(4.15) \quad \mathbf{1}_{\mathbb{R}^{dN}} = \sum_B \mathbf{1}_B$$

we can write

$$\begin{aligned} \mathcal{E}_Q[\Psi_N] &= \left\langle \Psi_N, \sum_{i=1}^N (-\Delta_{\mathbf{x}_i})_Q^s \Psi_N \right\rangle = \sum_B \left\langle \mathbf{1}_B \Psi_N, \sum_{i=1}^N (-\Delta_{\mathbf{x}_i})_Q^s \mathbf{1}_B \Psi_N \right\rangle \\ &= \sum_B \int_{(\mathbb{R}^d \setminus Q)^{N-|B|}} \|\Psi_N(\cdot; \mathbf{x}_{B^c})\|_{H^{s,|B|}(Q)}^2 \mathbf{1}_{(\mathbb{R}^d \setminus Q)^d} d\mathbf{x}_{B^c} \\ (4.16) \quad &\geq \sum_B \int_{(\mathbb{R}^d \setminus Q)^{N-|B|}} E_{|B|}(Q) \|\Psi_N(\cdot; \mathbf{x}_{B^c})\|_{L^2(Q^{|B|})}^2 d\mathbf{x}_{B^c} \\ &= \sum_B E_{|B|}(Q) \langle \Psi_N, \mathbf{1}_B \Psi_N \rangle. \end{aligned}$$

Here we have used the fact that $\mathbf{1}_B$ commutes with $(-\Delta_{\mathbf{x}_i})_Q^s$ and the shorthand notation

$$(\mathbf{x}_1, \dots, \mathbf{x}_N) = (\mathbf{x}_B; \mathbf{x}_{B^c}), \quad \mathbf{x}_B = (\mathbf{x}_\ell)_{\ell \in B} \in (\mathbb{R}^d)^{|B|}.$$

On the other hand, the partition of unity (4.15) implies that

$$\sum_B \langle \Psi_N, \mathbf{1}_B \Psi_N \rangle = \langle \Psi_N, \Psi_N \rangle = 1$$

and

$$\begin{aligned} \sum_B |B| \langle \Psi_N, \mathbb{1}_B \Psi_N \rangle &= \sum_B \left\langle \mathbb{1}_B \Psi_N, \sum_{i=1}^N \mathbb{1}_Q(\mathbf{x}_i) \mathbb{1}_B \Psi_N \right\rangle \\ &= \left\langle \Psi_N, \sum_{i=1}^N \mathbb{1}_Q(\mathbf{x}_i) \Psi_N \right\rangle = \int_Q \varrho_{\Psi_N}. \end{aligned}$$

Thus from (4.16) and (4.14) we conclude that

$$\begin{aligned} \mathcal{E}_Q[\Psi_N] &\geq \frac{C}{|Q|^{2s/d}} \sum_B [|B| - q]_+ \langle \Psi_N, \mathbb{1}_B \Psi_N \rangle \\ &\geq \frac{C}{|Q|^{2s/d}} \left[\sum_B (|B| - q) \langle \Psi_N, \mathbb{1}_B \Psi_N \rangle \right]_+ = \frac{C}{|Q|^{2s/d}} \left[\int_Q \varrho_{\Psi_N} - q \right]_+ \end{aligned}$$

by Jensen's inequality and the convexity of the function $t \mapsto [t]_+$. \square

5. Many-body Poincaré inequality

The crucial fact that the local energy $E_n(\Omega)$ in (4.8) eventually becomes positive with increasing particle number is the content of the following Poincaré inequality:

THEOREM 5.1 (Poincaré inequality for functions vanishing on diagonals). —

Fix an integer $k \geq 2$ and a bounded connected Lipschitz domain $\Omega \subset \mathbb{R}^d$. Assume that $2s > d(k-1)$. For $N \in \mathbb{N}$ large enough ($N \geq \lceil s \rceil^d k$ is sufficient) there exists a positive constant C depending only on s, k, N, Ω so that

$$(5.1) \quad \|u\|_{\dot{H}^{s, N}(\Omega)} \geq C \|u\|_{L^2(\Omega^N)}$$

for all $u \in C^\infty(\Omega^N)$ whose restriction to Δ_k is zero.

Since Theorem 5.1 is of independent interest, we state the result for more general domains although the result for cubes is sufficient for our application.

Conclusion of Theorem 2.2. — From Theorem 5.1 and Lemma 4.4 we obtain the local exclusion bound (4.13). Theorem 2.2 then immediately follows from the proof strategy in Section 3. \square

It remains to prove Theorem 5.1. The central fact used in the proof is that a function for which the left-hand side of (5.1) vanishes must be a polynomial, and that if a polynomial vanishes on too many diagonals it must be zero.

LEMMA 5.2 (Low-degree polynomials vanishing on diagonals are trivial). — *Given $d, k, S \in \mathbb{N}_1$ and $N \geq (S+1)^d k$. Let the dN -variable polynomial $f(\mathbf{x}_1, \dots, \mathbf{x}_N)$, with $\mathbf{x}_i \in \mathbb{R}^d$, satisfy*

- $\deg_{\mathbf{x}_j} f \leq S$ for all $j \in \{1, \dots, N\}$,
- $f(\mathbf{x}_1, \dots, \mathbf{x}_N) = 0$ on Δ_k .

Then $f \equiv 0$.

Proof. — The case $k = 1$ ($\Delta_1 = \mathbb{R}^d$) is trivial. We prove the other cases by induction.

Step 1. — Consider $d = 1$ and $k = 2$. Then $f(x_1, \dots, x_N) = 0$ if $x_i = x_j$ for some $i \neq j$. Consequently, when x_2, \dots, x_N are mutually different, the one variable polynomial $g(x_1) = f(x_1, x_2, \dots, x_N)$ has $\deg g \leq S$ but it has $N - 1$ different roots $x_1 = x_2, \dots, x_1 = x_N$. Therefore, if

$$N - 1 > S$$

(which holds if $N \geq (S + 1)k$) then $g(x_1) \equiv 0$. Thus

$$f(x_1, \dots, x_N) = 0$$

for all $x_1, \dots, x_N \in \mathbb{R}$ satisfying that x_2, \dots, x_N are mutually different. By continuity, we conclude that $f \equiv 0$.

Step 2. — Consider $d = 1$ and $k > 2$. Then $f(x_1, \dots, x_N) = 0$ if at least k points x_i 's coincide. Then if x_k, \dots, x_N are mutually different, the one-variable polynomial

$$g(x_1) = f(x_1, \dots, x_1, x_k, \dots, x_N)$$

has $\deg g \leq S(k - 1)$ but it has $N - k + 1$ different roots $x_1 = x_k, \dots, x_1 = x_N$. Therefore, if

$$N - k + 1 > S(k - 1)$$

(which holds if $N \geq (S + 1)k$) then $g \equiv 0$. Thus

$$f(x_1, \dots, x_1, x_k, \dots, x_N) = 0$$

if x_k, \dots, x_N are mutually different. By continuity, we conclude that

$$f(x_1, \dots, x_1, x_k, \dots, x_N) = 0$$

for all x_1, \dots, x_N . Similarly, by a renumbering, we can show that

$$f(x_1, x_2, \dots, x_N) = 0$$

if at least $(k - 1)$ points x_i 's coincide. By induction in k , we conclude that $f \equiv 0$.

Step 3. — Now consider $d > 1$ and $k \geq 2$. Let us denote

$$\mathbf{x}_i = (y_i, \mathbf{z}_i) \in \mathbb{R} \times \mathbb{R}^{d-1}.$$

Take

$$n = (S + 1)k, \quad N \geq (S + 1)^d k = (S + 1)^{d-1} n.$$

Then for any $\mathbf{z} \in \mathbb{R}^{d-1}$ and $\mathbf{x}_{n+1}, \dots, \mathbf{x}_N \in \mathbb{R}^d$, the polynomial

$$g(y_1, \dots, y_n) = f((y_1, \mathbf{z}), \dots, (y_n, \mathbf{z}), \mathbf{x}_{n+1}, \dots, \mathbf{x}_N)$$

satisfies that $\deg_{y_i} g \leq S$ and $g = 0$ if (at least) k points y_i 's coincide. By the result in the 1D case (with the choice $n = (S + 1)k$) we conclude that $g \equiv 0$. Similarly, we obtain that

$$f(\mathbf{x}_1, \dots, \mathbf{x}_N) = f((y_1, \mathbf{z}_1), \dots, (y_N, \mathbf{z}_N)) = 0$$

if at least n points \mathbf{z}_i 's coincide. By induction in d (i.e. using the induction hypothesis with $d - 1$ and $k = n$, $N \geq (S + 1)^{d-1} n$) we conclude that $f \equiv 0$. \square

We will also need the following technical Lemma 5.3, which essentially states that if a multivariable function is a polynomial in each variable separately, then it is a multivariable polynomial. The proof of this seemingly obvious fact is indeed non-trivial; see Carroll [Car61] for an elegant proof in the two variables case. Here we provide an alternative proof for n variables.

LEMMA 5.3. — Let $f(x_1, \dots, x_n) \in L^1_{\text{loc}}(\mathbb{R}^n)$ satisfy that for any $j = 1, 2, \dots, n$ and for a.e. $(x_1, \dots, x_{j-1}, x_{j+1}, \dots, x_n) \in \mathbb{R}^{n-1}$ the mapping $x_j \mapsto f(x_1, \dots, x_j, \dots, x_n)$ is a polynomial of degree at most M_j . Then f is a polynomial of n variables (x_1, \dots, x_n) of degree at most $M = \sum_{j=1}^n M_j$.

From the proof below, it is clear that we can replace \mathbb{R}^n by a subdomain (e.g. a cube).

Proof.

Step 1. — We use the notation $\mathbf{x} = (x_1, \dots, x_n) \in \mathbb{R}^n$, $\alpha = (\alpha_1, \dots, \alpha_n) \in \mathbb{N}_0^n$, and for every $j = 1, 2, \dots, n$ we write

$$\mathbf{x} = (x_j; \mathbf{x}'_j), \quad \alpha = (\alpha_j; \alpha'_j).$$

By assumption, for a.e. $\mathbf{x}'_j \in \mathbb{R}^{n-1}$, the mapping $x_j \mapsto f(x_j; \mathbf{x}'_j)$ is a polynomial of degree at most M_j . Therefore, for any $\alpha_j > M_j$, $D^{\alpha_j} f(\cdot; \mathbf{x}'_j) = 0$ as distribution on \mathbb{R} , namely

$$(5.2) \quad \int_{\mathbb{R}} f(x_j; \mathbf{x}'_j) D^{\alpha_j} h(x_j) dx_j = 0, \quad \forall h \in C_c^\infty(\mathbb{R}).$$

Consequently, $D^\alpha f = 0$ as distribution in \mathbb{R}^n if $|\alpha| > M$. Indeed, since $|\alpha| > M$ we have $\alpha_j > M_j$ for some j , and hence for any test function $\varphi \in C_c^\infty(\mathbb{R}^n)$ using Fubini's theorem and (5.2) we can write

$$\int_{\mathbb{R}^n} f D^\alpha \varphi d\mathbf{x} = \int_{\mathbb{R}^{n-1}} \left[\int_{\mathbb{R}} f(x_j; \mathbf{x}'_j) D^{\alpha_j} (D^{\alpha'_j} \varphi(x_j; \mathbf{x}'_j)) dx_j \right] d\mathbf{x}'_j = 0.$$

Step 2. — Thus it remains to prove that if $D^\alpha f = 0$ as distribution in \mathbb{R}^n for any $|\alpha| > M$, then f is a polynomial of n variables. We prove this statement by induction in M .

If $M = 0$, then $D_{x_j} f = 0$ as distribution for any $j = 1, 2, \dots, n$, and hence f is constant by [LL01, Theorem 6.1].

Now we prove the statement for $M \geq 1$ using the induction hypothesis for $M - 1$. From

$$D^\alpha f = 0, \quad \forall |\alpha| > M$$

we have for any $j = 1, 2, \dots, n$,

$$D^\alpha (D_{x_j} f) = 0, \quad \forall |\alpha| > M - 1.$$

Thus by the induction hypothesis for $M - 1$, $D_{x_j} f$ is a polynomial of n variables for any $j = 1, 2, \dots, n$. Since $D_{x_j} f \in C(\mathbb{R}^n)$ for all $j = 1, 2, \dots, n$, we obtain that $f \in C^1(\mathbb{R}^n)$ by [LL01, Theorem 6.10] and we have the formula [LL01, Theorem 6.9]

$$f(\mathbf{x}) = f(0) + \int_0^1 \mathbf{x} \cdot (\nabla f)(t\mathbf{x}) dt, \quad \forall \mathbf{x} \in \mathbb{R}^n.$$

The latter formula and the fact that $D_{x_j} f$ is a polynomial of n variables for any $j = 1, 2, \dots, n$ imply that f is a polynomial of n variables. This ends the proof of the Lemma 5.3. \square

Proof of Theorem 5.1. — We argue by contradiction. Assume that (5.1) is false, then there exists a sequence $u_n \in C^\infty(\Omega^N)$ satisfying $\|u_n\|_{L^2} = 1$, $u_n|_{\Delta_k} \equiv 0$, and

$$(5.3) \quad \|u_n\|_{\dot{H}^{s,N}(\Omega)} \rightarrow 0, \quad \text{as } n \rightarrow \infty.$$

In particular, u_n is bounded in the Sobolev space $H^\nu(\Omega^N)$ with $\nu = \min\{s, 1\}$. Indeed, for $d = 1$ this follows from Lemma A.1 and Sobolev's embedding theorem. If $d \geq 2$ then $s > 1$ and the claim follows from Sobolev's embedding theorem combined with that for any Ω the $\dot{H}^1(\Omega^N)$ and $\dot{H}^{1,N}(\Omega)$ seminorms are equivalent. By compactness of the embedding $H^\nu(\Omega^N) \subset L^2(\Omega^N)$, up to a subsequence, u_n converges strongly to a function P in $L^2(\Omega^N)$. Since $\|u_n\|_{L^2(\Omega^N)} = 1$ we have that $\|P\|_{L^2(\Omega^N)} = 1$.

On the other hand, by Poincaré's inequality for $\dot{H}^s(\Omega)$ (combining [LL01, Theorem 8.11] and [HSV13, Lemma 2.2])

$$\|u_n\|_{\dot{H}^{s,N}(\Omega)}^2 = \sum_{j=1}^N \int_{\Omega^{N-1}} \|u_n(x_j; \mathbf{x}')\|_{\dot{H}_{\mathbf{x}_j}^s(\Omega)}^2 d\mathbf{x}' \geq C \sum_{j=1}^N \int_{\Omega^N} |u_n(\mathbf{x}) - P_j^{(n)}(\mathbf{x})|^2 d\mathbf{x},$$

where $P_j^{(n)}(\mathbf{x})$ is a polynomial in \mathbf{x}_j of degree $\leq \lceil s-1 \rceil$. In fact, the polynomial can be written explicitly as

$$(5.4) \quad P_j^{(n)}(\mathbf{x}) = \sum_{|\beta| \leq \lceil s-1 \rceil} \mathbf{x}_j^\beta \left\langle \varphi_\beta(\mathbf{x}_j), u_n(\mathbf{x}_j; \mathbf{x}') \right\rangle_{L_{\mathbf{x}_j}^2(\Omega)}$$

for universal functions $\varphi_\beta \in C^\infty(\Omega)$. Since u_n converges strongly in $L^2(\Omega^N)$, we can conclude that $P_j^{(n)}(\mathbf{x}) \rightarrow P_j(\mathbf{x})$ strongly in $L^2(\Omega^N)$ and the limit is again a polynomial in \mathbf{x}_j of degree $\leq \lceil s-1 \rceil$. The assumption (5.3) allows us to identify the limiting functions and we find that

$$P(\mathbf{x}) = P_j(\mathbf{x}) \text{ in } L^2(\Omega^N), \quad \forall j.$$

Thus the function $P(x)$ is a polynomial in each variable \mathbf{x}_j (of degree $\leq \lceil s-1 \rceil$). By Lemma 5.3, $P(x)$ is a multivariate polynomial whose degree in each \mathbf{x}_j is $\leq \lceil s-1 \rceil$.

We now want to use that $u_n = 0$ on Δ_k to prove that $P = 0$ on Δ_k . Once this is done, then Lemma 5.2 implies that $P \equiv 0$ if $N \geq \lceil s \rceil^d k$. This contradicts that $\|P\|_{L^2(\Omega^N)} = 1$ and hence completes our proof. Note that if we can prove that $P \equiv 0$ in some open subset this is sufficient, in particular we can find some open cube $Q \subseteq \Omega$ and consider instead u_n and P restricted to Q^N .

We consider the diagonal $\mathbf{x}_1 = \mathbf{x}_2 = \dots = \mathbf{x}_k$; the other cases are treated identically. By Lebesgue's differentiation theorem it suffices to prove that

$$(5.5) \quad \lim_{\delta \rightarrow 0} \frac{1}{\delta^{d(k-1)}} \int_{Q^{N-k+1}} \int_{\max_{j \leq k} |\mathbf{x}_1 - \mathbf{x}_j| < \delta} |P(\mathbf{x}_1, \dots, \mathbf{x}_k; \mathbf{x}')| d\mathbf{x} = 0.$$

By Fatou's lemma we have for any $\delta > 0$ that

$$(5.6) \quad \begin{aligned} \int_{Q^{N-k+1}} \int_{\max_{j \leq k} |\mathbf{x}_1 - \mathbf{x}_j| < \delta} |P(\mathbf{x}_1, \dots, \mathbf{x}_k; \mathbf{x}')| d\mathbf{x} \\ \leq \lim_{n \rightarrow \infty} \int_{Q^{N-k+1}} \int_{\max_{j \leq k} |\mathbf{x}_1 - \mathbf{x}_j| < \delta} |u_n(\mathbf{x}_1, \dots, \mathbf{x}_k; \mathbf{x}')| d\mathbf{x}. \end{aligned}$$

Since $u_n = 0$ on Δ_k it holds that

$$(5.7) \quad \begin{aligned} \int_{Q^{N-k+1}} \int_{\max_{j \leq k} |\mathbf{x}_1 - \mathbf{x}_j| < \delta} |u_n(\mathbf{x}_1, \dots, \mathbf{x}_k; \mathbf{x}')| d\mathbf{x} \\ = \int_{Q^{N-k+1}} \int_{\max_{j \leq k} |\mathbf{x}_1 - \mathbf{x}_j| < \delta} |u_n(\mathbf{x}_1, \dots, \mathbf{x}_k; \mathbf{x}') - u_n(\mathbf{x}_1, \dots, \mathbf{x}_1; \mathbf{x}')| d\mathbf{x}. \end{aligned}$$

By Lemma A.1, any $u \in L^2(Q^l)$ with $\|u\|_{\dot{H}^{s,l}(Q)} < \infty$ satisfies that $u \in H^s(Q^l)$ and moreover there is a constant C depending only on Q, l, s such that

$$(5.8) \quad \|u\|_{H^s(Q^l)} \leq C \left(\|u\|_{L^2(Q^l)} + \|u\|_{\dot{H}^{s,l}(Q)} \right).$$

If $2s > dl$, by Sobolev's embedding theorem (see for instance [DNPV12, Theorem 8.2]), there is for any $\gamma \in (0, \min\{1, \frac{2s-dl}{2}\})$ a constant C so that

$$\|u\|_{C^{0,\gamma}(Q^l)} \leq C \|u\|_{H^s(Q^l)}, \quad \text{for all } u \in H^s(Q^l).$$

By assumption $2s > d(k-1)$, and hence we can apply this result to the function

$$(\mathbf{x}_2, \dots, \mathbf{x}_k) \mapsto u_n(\mathbf{x}_1, \dots, \mathbf{x}_k; \mathbf{x}')$$

(whose $\dot{H}^{s,k-1}(Q)$ -seminorm is bounded for a.e. $(\mathbf{x}_1, \mathbf{x}')$). Equation (5.7) then implies that

$$\begin{aligned} & \int_{Q^{N-k+1}} \int_{\max_j \leq k} |\mathbf{x}_1 - \mathbf{x}_j| < \delta} |u_n(\mathbf{x}_1, \dots, \mathbf{x}_k; \mathbf{x}')| d\mathbf{x} \\ & \leq C \int_{Q^{N-k+1}} \int_{\max_j \leq k} |\mathbf{x}_1 - \mathbf{x}_j| < \delta} \|u_n(\mathbf{x}_1, \mathbf{x}''; \mathbf{x}')\|_{H_{\mathbf{x}''}^{s,(Q^{k-1})}} |\mathbf{x}'' - \mathbf{x}_1'|^\gamma d\mathbf{x}'' d\mathbf{x}_1 d\mathbf{x}', \end{aligned}$$

where we set $\mathbf{x}'' = (\mathbf{x}_2, \dots, \mathbf{x}_k)$ and $\mathbf{x}_1' = (\mathbf{x}_1, \dots, \mathbf{x}_1)$. Applying (5.8) and Hölder's inequality yields

$$\begin{aligned} & \int_{Q^{N-k+1}} \int_{\max_j \leq k} |\mathbf{x}_1 - \mathbf{x}_j| < \delta} |u_n(\mathbf{x}_1, \dots, \mathbf{x}_k; \mathbf{x}')| d\mathbf{x} \\ & \leq C \delta^{d(k-1)+\gamma} \left(\|u_n\|_{L^2(Q^N)} + \|u_n\|_{\dot{H}^{s,N}(Q)} \right). \end{aligned}$$

Since $\|u_n\|_{L^2(Q^N)} + \|u_n\|_{\dot{H}^{s,N}(Q)} \leq C$ and $\gamma > 0$, we arrive at (5.5) which completes the proof of Theorem 5.1. \square

We finally note that the many-body nature of the wave functions is crucial for Theorem 5.1 to hold. The following example shows that the requirement that the particle number N is large, in fact typically strictly larger than k , is necessary.

PROPOSITION 5.4 (Counterexample to the k -body case). — *Theorem 5.1 cannot hold for $N < k$, or for $N = k$ if s is integer and*

$$\max\{d, 2\}(k-1) < 2s < (d+k)(k-1).$$

Replacing the condition $u|_{\Delta_k} = 0$ by the stronger condition

$$u \in \mathcal{H}_{0,k}^{s,N}(\mathbb{R}^d) := \overline{\{\Psi \in C_c^\infty(\mathbb{R}^{dN} \setminus \Delta_k)\}}^{H^s(\mathbb{R}^{dN})},$$

or

$$u \in \mathcal{H}_{W,k}^{s,N}(\mathbb{R}^d) := \left\{ \Psi \in H^s(\mathbb{R}^{dN}) : \int_{\mathbb{R}^{dN}} W_{s,k} |\Psi|^2 < \infty \right\},$$

with the k -particle generalization of W_s ,

$$W_{s,k}(\mathbf{x}) := \sum_{\substack{A \subseteq \{1, \dots, N\} \\ |A|=k}} \left(\sum_{\substack{j, l \in A \\ j < l}} |\mathbf{x}_j - \mathbf{x}_l|^2 \right)^{-2s},$$

does not help.

Proof. — If $N < k$ there is no diagonal set Δ_k and we may take the constant function as a counterexample. For $N = k$ we consider the polynomial

$$u(\mathbf{x}_1, \dots, \mathbf{x}_k) := \prod_{1 \leq j < l \leq k} (x_{j,1} - y_{l,1}),$$

for which, by the arithmetic mean-geometric mean inequality and the triangle inequality,

$$\begin{aligned} |u(\mathbf{x})|^2 &\lesssim \left(\sum_{j < l} |x_{j,1} - x_{l,1}|^2 \right)^{\binom{k}{2}} \leq \left(\sum_{j < l} |\mathbf{x}_j - \mathbf{x}_l|^2 \right)^{\binom{k}{2}} \\ &\lesssim \left(\sum_{l \geq 2} |\mathbf{x}_1 - \mathbf{x}_l|^2 \right)^{\binom{k}{2}} =: R^{2\binom{k}{2}}, \end{aligned}$$

where $R \geq 0$ may serve as a radial coordinate on $\mathbb{R}^{d(k-1)}$ relative to \mathbf{x}_1 . Hence, we have that

$$\int_{Q^k} W_{s,k} |u|^2 \lesssim \int_Q \int_{Q^{k-1}} \frac{R^{2\binom{k}{2}}}{R^{2s}} d\mathbf{x}_2 \dots d\mathbf{x}_k d\mathbf{x}_1 \lesssim \int_0^C R^{k(k-1)-2s+d(k-1)-1} dR < \infty,$$

if $d(k-1) < 2s < (d+k)(k-1)$. Thus (analogously to Lemma B.2, and by extension)

$$u \in \mathcal{H}_{W,k}^{s,N}(\mathbb{R}^d) \subseteq \mathcal{H}_{0,k}^{s,N}(\mathbb{R}^d).$$

On the other hand

$$\|u\|_{\dot{H}^{s,k}(\Omega)}^2 = \sum_{j=1}^k \sum_{|\alpha|=m} \frac{m!}{\alpha!} \int_{Q^k} |D_{\mathbf{x}_j}^\alpha u|^2 = 0,$$

if $s = m > k - 1$. □

A particular case included in the above is $d = 3$, $s = 2$, $k = 2$, with the function $u(\mathbf{x}, \mathbf{y}) := x_1 - y_1$.

Appendix A. Equivalence of Sobolev spaces

In this appendix we discuss the N -particle space

$$H^{s,N}(\Omega) := \left\{ u \in L^2(\Omega^N) : \|u\|_{\dot{H}^{s,N}(\Omega)} < \infty \right\}$$

and its relation to the standard Sobolev space $H^s(\Omega^N)$.

If $\Omega = \mathbb{R}^d$ the equivalence of the seminorms (and consequently the spaces)

$$(A.1) \quad c_{s,N} \|u\|_{\dot{H}^s(\mathbb{R}^{dN})} \leq \|u\|_{\dot{H}^{s,N}(\mathbb{R}^d)} \leq C_{s,N} \|u\|_{\dot{H}^s(\mathbb{R}^{dN})}$$

can be seen via the Fourier transform. However, the constants in the equivalence depend on N and s . In particular, if $s \neq 1$ the equivalence degenerates as N tends to infinity; either $c_{s,N} \rightarrow 0$ or $C_{s,N} \rightarrow \infty$. Specifically, the sharp constants in (A.1) are given by

$$c_{s,N} = \min \left\{ 1, N^{(1-s)/2} \right\} \quad \text{and} \quad C_{s,N} = \max \left\{ 1, N^{(1-s)/2} \right\}.$$

Thus it is a slightly subtle question of what happens to these spaces in the many-body limit. An even more subtle question is what happens to the local versions of these spaces, i.e. when \mathbb{R}^d is replaced by $\Omega \subsetneq \mathbb{R}^d$. For us, the following equivalence of the spaces in the case of cubes will suffice:

LEMMA A.1. — *Let $u \in L^2(Q^N)$, $Q = [0, 1]^d$. There exist positive constants c, C depending only on d, s, N so that*

$$c \left(\|u\|_{L^2(Q^N)} + \|u\|_{\dot{H}^{s,N}(Q)} \right) \leq \|u\|_{H^s(Q^N)} \leq C \left(\|u\|_{L^2(Q^N)} + \|u\|_{\dot{H}^{s,N}(Q)} \right).$$

Lemma A.1 is an immediate consequence of the equivalence (A.1) of the two seminorms on \mathbb{R}^{dN} and the following extension Lemma A.2:

LEMMA A.2. — *Let $u \in L^2(Q^N)$, $Q = [0, 1]^d$, and assume that $\|u\|_{L^2(Q^N)} + \|u\|_{\dot{H}^{s,N}(Q)} < \infty$. There exists a function $\tilde{u} \in L^2(\mathbb{R}^{dN})$ with compact support satisfying*

$$\tilde{u}|_{Q^N} = u, \quad \text{and} \quad \|\tilde{u}\|_{L^2(\mathbb{R}^{dN})} + \|\tilde{u}\|_{\dot{H}^{s,N}(\mathbb{R}^d)} \leq C \left(\|u\|_{L^2(Q^N)} + \|u\|_{\dot{H}^{s,N}(Q)} \right),$$

where C is a constant depending only on s, d and N .

Proof. — We shall prove the Lemma A.2 by using higher-order reflection through one side of the hypercube Q at a time. To this end we recall that if $v \in C^n([0, 1])$, for some $n \geq 0$, we can construct an explicit extension $\tilde{v} \in C^n((-\infty, 1])$ satisfying $\tilde{v}(x) = 0$ when $x < -\delta$. Namely, set

$$\tilde{v}(x) = \begin{cases} v(x), & \text{if } x \in [0, 1], \\ \varphi(x) \sum_{j=1}^{n+1} \lambda_j v(-x/j), & \text{if } x < 0, \end{cases}$$

where $\varphi \in C^\infty((-\infty, 0])$ such that $\varphi(x) \equiv 0$ for $x < -\delta$ and $\varphi(x) \equiv 1$ in $[-\delta/2, 0]$. What remains is to verify that we can choose the λ_j 's so that $\tilde{v} \in C^n$. But if we differentiate \tilde{v} for x away from zero we see that the system of equations that we need the λ_j to satisfy to get continuity of the derivatives across $x = 0$ is

$$\left[(-j)^{1-i} \right]_{i,j=1}^{n+1} \begin{pmatrix} \lambda_1 \\ \vdots \\ \lambda_{n+1} \end{pmatrix} = \begin{pmatrix} 1 \\ \vdots \\ 1 \end{pmatrix}.$$

But the determinant of this matrix is non-zero (it is a Vandermonde matrix) and hence there exists a unique solution $(\lambda_1, \dots, \lambda_{n+1})$.

We shall now prove that we can use this one-dimensional extension repeatedly to construct an extension of u to \mathbb{R}^{dN} . The idea is to use the one-dimensional result one coordinate at a time and show that the new function in each step has the quantity corresponding to the $\dot{H}^{s,N}$ -seminorm controlled by that of u .

Without loss we can assume that $u \in C^n(Q^N)$ (the construction is stable under approximation), where we take $n = \lceil s \rceil$. Consider $u(x_1; \mathbf{x}')$, $x_1 \in [0, 1]$ and $\mathbf{x}' \in [0, 1]^{dN-1}$. And apply the above Lemma A.2 for each fixed \mathbf{x}' , that is, we define v_1 by

$$v_1(x_1; \mathbf{x}') = \begin{cases} u(x_1; \mathbf{x}'), & \text{if } x_1 \in [0, 1], \\ \varphi(x_1) \sum_{j=1}^{n+1} \lambda_j u(-x_1/j; \mathbf{x}'), & \text{if } x_1 \in [-1, 0]. \end{cases}$$

It is a simple calculation to use Sobolev's embedding theorem to prove that we can bound the L^p -norm of l^{th} order derivatives of v_1 by the corresponding one for u if $l \leq n$. We need to prove that also the fractional order seminorm is preserved. That is, we wish to show that, with $s = m + \sigma$ and $Q' = [-1, 1] \times [0, 1]^{d-1}$,

$$\begin{aligned}
 (A.2) \quad & \int_{Q^{N-1}} \iint_{Q' \times Q'} \frac{|D_{\mathbf{x}_1}^\alpha v_1(\mathbf{x}_1; \mathbf{x}') - D_{\mathbf{y}_1}^\alpha v_1(\mathbf{y}_1; \mathbf{x}')|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
 & + \sum_{i=2}^N \int_{Q' \times Q^{N-2}} \iint_{Q \times Q} \frac{|D_{\mathbf{x}_i}^\alpha v_1(\mathbf{x}_i; \mathbf{x}') - D_{\mathbf{y}_i}^\alpha v_1(\mathbf{y}_i; \mathbf{x}')|^2}{|\mathbf{x}_i - \mathbf{y}_i|^{d+2\sigma}} d\mathbf{x}_i d\mathbf{y}_i d\mathbf{x}' \\
 & \leq C \left(\|u\|_{\dot{H}^{s,N}(Q)}^2 + \|u\|_{L^2(Q^N)}^2 \right),
 \end{aligned}$$

for all multi-indices $|\alpha| = m$. If we can prove this inequality, then by repeating the procedure to extend v_1 to $x_1 > 1$ the same proof gives that we can bound the corresponding $\dot{H}^{s,N}$ quantity in terms of that of v_1 , and hence u . By repeating the procedure for each coordinate at a time we, after $2dN$ reflections, find a function $\tilde{u} \in L^2(\mathbb{R}^{dN})$ satisfying the claims of the lemma. Thus all that remains is to prove (A.2).

We start with the first term which is also the most difficult:

$$\begin{aligned}
 (A.3) \quad & \int_{Q^{N-1}} \iint_{Q' \times Q'} \frac{|D_{\mathbf{x}_1}^\alpha v_1(\mathbf{x}_1; \mathbf{x}') - D_{\mathbf{y}_1}^\alpha v_1(\mathbf{y}_1; \mathbf{x}')|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
 & = \int_{Q^{N-1}} \left[\iint_{Q \times Q} \frac{|D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - D_{\mathbf{y}_1}^\alpha u(\mathbf{y}_1; \mathbf{x}')|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 \right. \\
 & \quad + 2 \iint_{Q \times (Q' \setminus Q)} \frac{|D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - \sum_{j=1}^{n+1} \lambda_j D_{\mathbf{y}_1}^\alpha (\varphi(y_{1,1}) u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}'))|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 \\
 & \quad + \iint_{(Q' \setminus Q) \times (Q' \setminus Q)} |\mathbf{x}_1 - \mathbf{y}_1|^{-d-2\sigma} \left| \sum_{j=1}^{n+1} \lambda_j \left(D_{\mathbf{x}_1}^\alpha [\varphi(x_{1,1}) u(-x_{1,1}/j, \mathbf{x}'_1; \mathbf{x}')] \right. \right. \\
 & \quad \left. \left. - D_{\mathbf{y}_1}^\alpha [\varphi(y_{1,1}) u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}')] \right) \right|^2 d\mathbf{x}_1 d\mathbf{y}_1 \Big] d\mathbf{x}'.
 \end{aligned}$$

Clearly the integral over $Q \times Q$ is bounded by $\|u\|_{\dot{H}^{s,N}(Q)}$. We treat the two remaining terms separately. In order to bound the integral over $Q \times (Q' \setminus Q)$ we write

$$\begin{aligned}
 Q_1 &= \{\mathbf{x} \in Q' \setminus Q : x_1 > -\delta/2\}, \\
 Q_2 &= \{\mathbf{x} \in Q' \setminus Q : x_1 \leq -\delta/2\}.
 \end{aligned}$$

Thus we can bound the second integral in (A.3) as follows:

$$\begin{aligned}
& \int_{Q^{N-1}} \iint_{Q \times (Q' \setminus Q)} \frac{\left| D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - \sum_{j=1}^{n+1} \lambda_j D_{\mathbf{y}_1}^\alpha (\varphi(y_{1,1}) u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}')) \right|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
&= \int_{Q^{N-1}} \iint_{Q \times Q_1} \frac{\left| D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - \sum_{j=1}^{n+1} \lambda_j D_{\mathbf{y}_1}^\alpha (u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}')) \right|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
&+ \int_{Q^{N-1}} \iint_{Q \times Q_2} \frac{\left| D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - \sum_{j=1}^{n+1} \lambda_j D_{\mathbf{y}_1}^\alpha (\varphi(y_{1,1}) u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}')) \right|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
&\leq \int_{Q^{N-1}} \iint_{Q \times Q_1} \frac{\left| D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - \sum_{j=1}^{n+1} \lambda_j (-j)^{-\alpha_1} D_{\mathbf{y}_1}^\alpha u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}') \right|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
&+ \frac{C}{\delta^{d+2\sigma}} \int_{Q^{N-1}} \iint_{Q \times Q_2} \left| D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - \sum_{j=1}^{n+1} \lambda_j D_{\mathbf{y}_1}^\alpha (\varphi(y_{1,1}) u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}')) \right|^2 d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}'.
\end{aligned}$$

Using the triangle inequality and Sobolev's embedding theorem one finds that the second term is $\lesssim \|u\|_{L^2(Q^N)}^2 + \|u\|_{\dot{H}^{s,N}(Q)}^2$. Since $\sum_j \lambda_j (-j)^{-\alpha_1} = 1$ for any $\alpha_1 \leq m+1$, one obtains for the first integral

$$\begin{aligned}
& \int_{Q^{N-1}} \iint_{Q \times Q_1} \frac{\left| D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - \sum_{j=1}^{n+1} \lambda_j (-j)^{-\alpha_1} D_{\mathbf{y}_1}^\alpha u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}') \right|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
&= \int_{Q^{N-1}} \iint_{Q \times Q_1} \frac{\left| \sum_{j=1}^{n+1} \lambda_j (-j)^{-\alpha_1} \left(D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - D_{\mathbf{y}_1}^\alpha u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}') \right) \right|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
&\leq C \sum_{j=1}^{n+1} \int_{Q^{N-1}} \iint_{Q \times Q_1} \frac{\left| D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - D_{\mathbf{y}_1}^\alpha u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}') \right|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
&\leq C \sum_{j=1}^{n+1} \int_{Q^{N-1}} \iint_{Q \times Q} \frac{\left| D_{\mathbf{x}_1}^\alpha u(\mathbf{x}_1; \mathbf{x}') - D_{\mathbf{y}_1}^\alpha u(\mathbf{y}_1; \mathbf{x}') \right|^2}{(|\mathbf{x}'_1 - \mathbf{y}'_1|^2 + (x_{1,1} + jy_{1,1})^2)^{d/2+\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
&\leq C \|u\|_{\dot{H}^{s,N}(Q)}^2.
\end{aligned}$$

In the last step we used the inequality $(x + jy)^2 \geq (x - y)^2$ for $x, y \geq 0$ and $j \geq 1$.

For the last integral in (A.3) we have

$$\begin{aligned}
& \int_{Q^{N-1}} \iint_{(Q' \setminus Q) \times (Q' \setminus Q)} |\mathbf{x}_1 - \mathbf{y}_1|^{-d-2\sigma} \left| \sum_{j=1}^{n+1} \lambda_j \left(D_{\mathbf{x}_1}^\alpha [\varphi(x_{1,1}) u(-x_{1,1}/j, \mathbf{x}'_1; \mathbf{x}')] \right. \right. \\
& \quad \left. \left. - D_{\mathbf{y}_1}^\alpha [\varphi(y_{1,1}) u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}')] \right) \right|^2 d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
&= \int_{Q^{N-1}} \iint_{(Q' \setminus Q) \times (Q' \setminus Q)} |\mathbf{x}_1 - \mathbf{y}_1|^{-d-2\sigma} \left| \sum_{j=1}^{n+1} \sum_{\gamma+\beta=\alpha_1} \lambda_j (-j)^{-\beta} \right. \\
& \quad \left. \times \left(\varphi^{(\gamma)}(x_{1,1}) D_{\mathbf{x}_1}^{\alpha'} u(-x_{1,1}/j, \mathbf{x}'_1; \mathbf{x}') - \varphi^{(\gamma)}(y_{1,1}) D_{\mathbf{y}_1}^{\alpha'} u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}') \right) \right|^2 d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}',
\end{aligned}$$

where we set α' as the multi-index α but with α_1 exchanged for β . By the triangle inequality and the fact that $\sum_j \lambda_j (-j)^{-\beta} = 1$ the integral is smaller than

$$\begin{aligned}
& C \sum_{j=1}^{n+1} \int_{Q^{N-1}} \iint_{(Q' \setminus Q) \times (Q' \setminus Q)} |\mathbf{x}_1 - \mathbf{y}_1|^{-d-2\sigma} \left| \sum_{\gamma+\beta=\alpha_1} \left(\varphi^{(\gamma)}(x_{1,1}) D_{\mathbf{x}_1}^{\alpha'} u(-x_{1,1}/j, \mathbf{x}'_1; \mathbf{x}') \right. \right. \\
& \quad \left. \left. - \varphi^{(\gamma)}(y_{1,1}) D_{\mathbf{y}_1}^{\alpha'} u(-y_{1,1}/j, \mathbf{y}'_1; \mathbf{x}') \right) \right|^2 d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
& \leq C \sum_{j=1}^{n+1} \int_{Q^{N-1}} \iint_{Q \times Q} (|\mathbf{x}'_1 - \mathbf{y}'_1|^2 + j^2(x_{1,1} - y_{1,1})^2)^{-d/2-\sigma} \\
& \quad \times \left| \sum_{\gamma+\beta=\alpha_1} \left(\varphi^{(\gamma)}(-jx_{1,1}) D_{\mathbf{x}_1}^{\alpha'} u(\mathbf{x}_1; \mathbf{x}') - \varphi^{(\gamma)}(-jy_{1,1}) D_{\mathbf{y}_1}^{\alpha'} u(\mathbf{y}_1; \mathbf{x}') \right) \right|^2 d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
& \leq C \sum_{j=1}^{n+1} \int_{Q^{N-1}} \iint_{Q \times Q} \frac{|D_{\mathbf{x}_1}^\alpha [\varphi(-jx_{1,1}) u(\mathbf{x}_1; \mathbf{x}')] - D_{\mathbf{y}_1}^\alpha [\varphi(-jy_{1,1}) u(\mathbf{y}_1; \mathbf{x}')]|^2}{|\mathbf{x}_1 - \mathbf{y}_1|^{d+2\sigma}} d\mathbf{x}_1 d\mathbf{y}_1 d\mathbf{x}' \\
& \leq C \|u\|_{\dot{H}^{s, N}(Q)}^2,
\end{aligned}$$

where we used that $\|\psi u\|_{\dot{H}^s(Q)} \leq C_\psi \|u\|_{\dot{H}^s(Q)}$ for any $\psi \in C^\infty(Q)$.

To show that the remaining terms in (A.2) are $\lesssim \|u\|_{L^2}^2 + \|u\|_{\dot{H}^{s, N}}^2$ one can proceed in an almost identical manner. The main difference is that in these terms the differentiation is with respect other variables than the variable in which the extension has been made, and the splitting of the integrals is slightly different. However, in the end this only simplifies each step of the proof of Lemma A.2. \square

Appendix B. Spaces of contact interaction

We consider in the following only 2-particle diagonals \triangle , for simplicity, however analogous statements can be made for the case of k -particle diagonals.

Define for $N \geq 2$ the restricted N -particle spaces

$$\begin{aligned}\mathcal{H}_W^{s,N}(\mathbb{R}^d) &:= \left\{ \Psi \in H^s(\mathbb{R}^{dN}) : \int_{\mathbb{R}^{dN}} W_s |\Psi|^2 < \infty \right\}, \\ \mathcal{H}_0^{s,N}(\mathbb{R}^d) &:= \overline{\left\{ \Psi \in C_c^\infty(\mathbb{R}^{dN} \setminus \Delta) \right\}}^{H^s(\mathbb{R}^{dN})}, \\ \mathcal{H}^{s,N}(\mathbb{R}^d) &:= \overline{\left\{ \Psi \in C_c^\infty(\mathbb{R}^{dN}) : \Psi|_\Delta = 0 \right\}}^{H^s(\mathbb{R}^{dN})}.\end{aligned}$$

Then we have for all $s > 0$ the chain of inclusions

$$\mathcal{H}_W^{s,N}(\mathbb{R}^d) \subseteq \mathcal{H}_0^{s,N}(\mathbb{R}^d) \subseteq \mathcal{H}^{s,N}(\mathbb{R}^d) \subseteq H^s(\mathbb{R}^{dN}).$$

The latter two inclusions are trivial while the first one will be proved below. Moreover, for $2s < d$ all four spaces are equal by the Hardy–Rellich inequality (see e.g. [Yaf99]):

$$\int_{\mathbb{R}^{dN}} |\mathbf{x}_1 - \mathbf{x}_2|^{-2s} |\Psi(\mathbf{x}_1; \mathbf{x}')|^2 d\mathbf{x}_1 d\mathbf{x}' \leq C \int_{\mathbb{R}^{d(N-1)}} \|\Psi\|_{\dot{H}_{\mathbf{x}_1}^s(\mathbb{R}^d)}^2 d\mathbf{x}' \leq C \|\Psi\|_{H^s(\mathbb{R}^{dN})}^2.$$

In the critical case $2s = d$ we still have $\mathcal{H}_0^{s,N}(\mathbb{R}^d) = \mathcal{H}^{s,N}(\mathbb{R}^d) = H^s(\mathbb{R}^{dN})$, as is also shown below, but a strict inclusion $\mathcal{H}_W^{s,N}(\mathbb{R}^d) \subsetneq \mathcal{H}_0^{s,N}(\mathbb{R}^d)$, as illustrated by

$$\Psi(\mathbf{x}) = e^{-|\mathbf{x}|^2}$$

which is in $H^s(\mathbb{R}^{dN})$ but not in $\mathcal{H}_W^{s,N}(\mathbb{R}^d)$ due to the non-integrability of W_s . For $2s > d$ and $s - d/2 \notin \mathbb{Z}$ it again holds by the Hardy–Rellich inequality that $\mathcal{H}_W^{s,N}(\mathbb{R}^d) = \mathcal{H}_0^{s,N}(\mathbb{R}^d)$, while not necessarily $\mathcal{H}_0^{s,N}(\mathbb{R}^d) = \mathcal{H}^{s,N}(\mathbb{R}^d)$, as with the example

$$\Psi(x_1, x_2) = (x_1 - x_2)e^{-|x|^2}$$

which is in $\mathcal{H}^{s,2}(\mathbb{R}^d)$ but not in $\mathcal{H}_W^{s,2}(\mathbb{R}^d)$ for $s = 2$ and $d = 1$.

Let $\chi_\varepsilon^{(*)}(\mathbf{x}) := \prod_{1 \leq j < k \leq N} \varphi_\varepsilon^{(*)}(\mathbf{x}_j - \mathbf{x}_k)$ where $\varphi_\varepsilon(\mathbf{x}) = \varphi(|\mathbf{x}|/\varepsilon)$ and $\varphi_\varepsilon^*(\mathbf{x}) = \varphi^*(\varepsilon \ln |\mathbf{x}|)$. We take $\varphi^{(*)}$ as smooth functions from \mathbb{R} to $[0, 1]$ such that $\varphi(x) = 0$ for $x \leq 1$, $\varphi(x) = 1$ for $x \geq 2$, and $\varphi^*(x) = 0$ for $x \leq -2$, $\varphi^*(x) = 1$ for $x \geq -1$.

LEMMA B.1. — *Let $\Omega \subset \mathbb{R}^d$ be open and bounded. For all $s = m + \sigma > 0$, $d \geq 1$ and $N \geq 1$ it holds as $\varepsilon \rightarrow 0$ that*

$$\|\chi_\varepsilon\|_{\dot{H}^{s,N}(\Omega)} \leq C\varepsilon^{d/2-s}, \quad \|D^\alpha \chi_\varepsilon\|_{\dot{H}^{\sigma,N}(\Omega)} \leq C\varepsilon^{d/2-|\alpha|-\sigma},$$

while for $2s = d$

$$\|\chi_\varepsilon^*\|_{\dot{H}^{s,N}(\Omega)} \leq C\varepsilon^{1/2}, \quad \|D^\alpha \chi_\varepsilon^*\|_{\dot{H}^{\sigma,N}(\Omega)} \leq C\varepsilon^{1/2}$$

for $|\alpha| \leq d/2 - \sigma$.

Proof. — For $\alpha \neq 0$ there are in $D_{\mathbf{x}_j}^\alpha \chi_\varepsilon$ a total of $|\alpha|$ derivatives of functions $\varphi_\varepsilon(\mathbf{x}_j - \mathbf{x}_k)$, $k \neq j$, and remaining factors involving the other particles. These factors are uniformly bounded while each derivative yields an additional factor $1/\varepsilon$, while reducing the support in \mathbf{x}_j to $B_{2\varepsilon}(\mathbf{x}_k) \setminus B_\varepsilon(\mathbf{x}_k)$. Furthermore, we thus have

$$\begin{aligned}|D^\alpha \varphi_\varepsilon(\mathbf{x})| &\leq C\varepsilon^{-|\alpha|} \mathbf{1}_{B_{2\varepsilon}(0) \setminus B_\varepsilon(0)}, \\ |D^\alpha \varphi_\varepsilon(\mathbf{x}) - D^\alpha \varphi_\varepsilon(\mathbf{y})| &\leq C\varepsilon^{-|\alpha|-1} |\mathbf{x} - \mathbf{y}| \mathbf{1}_{\mathbf{x}, \mathbf{y} \in B_{2\varepsilon}(0) \setminus B_\varepsilon(0)},\end{aligned}$$

and for $B(j, \varepsilon) = \cup_{k \neq j} B_{2\varepsilon}(\mathbf{x}_k) \setminus \cup_{k \neq j} B_\varepsilon(\mathbf{x}_k)$,

$$|\chi_\varepsilon(\mathbf{x}_j; \mathbf{x}') - \chi_\varepsilon(\mathbf{y}_j; \mathbf{x}')| \leq C\varepsilon^{-1} |\mathbf{x}_j - \mathbf{y}_j| \mathbb{1}_{\mathbf{x}_j, \mathbf{y}_j \in B(j, \varepsilon)},$$

and

$$|D^\alpha \chi_\varepsilon(\mathbf{x}_j; \mathbf{x}') - D^\alpha \chi_\varepsilon(\mathbf{y}_j; \mathbf{x}')| \leq C\varepsilon^{-|\alpha|-1} |\mathbf{x}_j - \mathbf{y}_j| \mathbb{1}_{\mathbf{x}_j, \mathbf{y}_j \in B(j, \varepsilon)}.$$

Hence, $\|D^\alpha \chi_\varepsilon\|_{L^2_{\mathbf{x}_j}(\Omega)}^2 \lesssim \varepsilon^{-2|\alpha|+d}$, and for any $0 < \sigma < 1$

$$\begin{aligned} \|D^\alpha \chi_\varepsilon\|_{H^{\sigma}_{\mathbf{x}_j}(\Omega)}^2 &= \iint_{\Omega \times \Omega} \frac{|D^\alpha \chi_\varepsilon(\mathbf{x}; \mathbf{x}') - D^\alpha \chi_\varepsilon(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^{d+2\sigma}} d\mathbf{x} d\mathbf{y} \\ &\lesssim \varepsilon^{-2|\alpha|-2} \sum_{k \neq j} \iint_{B_{2\varepsilon}(\mathbf{x}_k) \times B_{2\varepsilon}(\mathbf{x}_k)} |\mathbf{x} - \mathbf{y}|^{-d-2\sigma+2} d\mathbf{x} d\mathbf{y} \lesssim \varepsilon^{-2|\alpha|-2\sigma+d}, \end{aligned}$$

so that $\|\chi_\varepsilon\|_{\dot{H}^{s, N}(\Omega)}^2 \lesssim \varepsilon^{-2s+d}$.

Similarly, for χ_ε^* we consider $B(j, \varepsilon) = \cup_{k \neq j} B_{e^{-1/\varepsilon}}(\mathbf{x}_k) \setminus \cup_{k \neq j} B_{e^{-2/\varepsilon}}(\mathbf{x}_k)$ and

$$(B.1) \quad |D^\alpha \varphi_\varepsilon^*(\mathbf{x})| = |D^\alpha \varphi^*(\varepsilon \ln |\mathbf{x}|)| \leq C\varepsilon |\mathbf{x}|^{-|\alpha|} \mathbb{1}_{B_{e^{-1/\varepsilon}}(0) \setminus B_{e^{-2/\varepsilon}}(0)}.$$

In χ_ε^* this could involve different points \mathbf{x}_k but the worst case is if they are the same,

$$\|D^\alpha \chi_\varepsilon^*\|_{L^2_{\mathbf{x}_j}(\Omega)}^2 \lesssim \varepsilon^2 \int_{B(j, \varepsilon)} |\mathbf{x}_j - \mathbf{x}_k|^{-2|\alpha|} d\mathbf{x}_j \lesssim \begin{cases} \varepsilon^2 & \text{for } 0 < 2|\alpha| < d, \\ \varepsilon^2 \int_{-2\varepsilon^{-1}}^{-\varepsilon^{-1}} ds = \varepsilon & \text{for } 2|\alpha| = d. \end{cases}$$

This covers the even-dimensional critical case $d = 2m$, $m \in \mathbb{N}_1$.

In the odd-dimensional critical case $d = 2m + 2\sigma$, $\sigma = 1/2$, we observe that

$$\begin{aligned} \|D^\alpha \chi_\varepsilon\|_{H^{\sigma}_{\mathbf{x}_j}(\Omega)}^2 &= \iint_{\Omega \times \Omega} \frac{|D^\alpha \chi_\varepsilon(\mathbf{x}; \mathbf{x}') - D^\alpha \chi_\varepsilon(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^{d+1}} d\mathbf{x} d\mathbf{y} \\ &\lesssim \varepsilon^{-2|\alpha|-2} \sum_{k \neq j} \iint_{B_{2\varepsilon}(\mathbf{x}_k) \times B_{2\varepsilon}(\mathbf{x}_k)} |\mathbf{x} - \mathbf{y}|^{-d+1} d\mathbf{x} d\mathbf{y} \lesssim \varepsilon^{-2|\alpha|-1+d}, \end{aligned}$$

which is not enough for $2|\alpha| = d - 1$. Instead we shall use χ_ε^* .

For the case $2|\alpha| = d - 1$ things are a bit less straightforward. We start with the case $d = 1$ which is the easiest. Here our approach differs slightly due to the fact that in this case $|\alpha| = 0$.

Let $U_1 = \cap_{k \neq j} B_{e^{-1/\varepsilon}}(\mathbf{x}_k)^c$, $U_2 = \cup_{k \neq j} B_{e^{-2/\varepsilon}}(\mathbf{x}_k)$ and $U = \Omega \setminus (U_1 \cup U_2)$.

We estimate the seminorm $\|\chi_\varepsilon^*\|_{\dot{H}^s_{\mathbf{x}_j}(\Omega)}$. By construction of χ_ε^* we have that

$$|\chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}')| \leq 1, \quad \forall \mathbf{x}, \mathbf{y} \in \Omega.$$

Moreover, the difference is zero whenever $(\mathbf{x}, \mathbf{y}) \in U_1^2 \cup U_2^2$.

For \mathbf{x} and \mathbf{y} close we need to estimate this quantity more precisely. By Taylor's theorem we can estimate

$$\begin{aligned} & \left| \chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}') \right| \\ &= \left| \int_0^1 \sum_{k \neq j}^N \left(\prod_{i \notin \{k, j\}} \varphi_\varepsilon^*(\mathbf{x} - \mathbf{x}_i) \right) (\varphi^*)' \left(\varepsilon (\ln |\mathbf{x} - \mathbf{x}_k| + t(\ln |\mathbf{y} - \mathbf{x}_k| - \ln |\mathbf{x} - \mathbf{x}_k|)) \right) \right. \\ & \quad \left. \times \varepsilon (\ln |\mathbf{y} - \mathbf{x}_k| - \ln |\mathbf{x} - \mathbf{x}_k|) dt \right| \leq C\varepsilon \sum_{k \neq j} \left| \ln |\mathbf{x} - \mathbf{x}_k| - \ln |\mathbf{y} - \mathbf{x}_k| \right|. \end{aligned}$$

By symmetry in \mathbf{x}, \mathbf{y} we find

$$\begin{aligned} (B.2) \quad \|\chi_\varepsilon^*\|_{\dot{H}_{\mathbf{x}_j}^s(\Omega)}^2 &= \iint_{\Omega \times \Omega} \frac{|\chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^2} d\mathbf{x} d\mathbf{y} \\ &\leq 2 \iint_{U \times \Omega} \frac{|\chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^2} d\mathbf{x} d\mathbf{y} + 2 \iint_{U_1 \times U_2} \frac{1}{|\mathbf{x} - \mathbf{y}|^2} d\mathbf{x} d\mathbf{y}. \end{aligned}$$

The latter term is fairly easy to estimate:

$$\iint_{U_1 \times U_2} \frac{1}{|\mathbf{x} - \mathbf{y}|^2} d\mathbf{x} d\mathbf{y} \leq 2|U_1|e^{-2/\varepsilon} \int_{-e^{-2/\varepsilon}}^{e^{-2/\varepsilon}} \frac{1}{(e^{-1/\varepsilon} + r)^2} dr \leq Ce^{-2/\varepsilon}.$$

We return to the remaining term of (B.2):

$$\begin{aligned} & \iint_{U \times \Omega} \frac{|\chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^2} d\mathbf{x} d\mathbf{y} \\ & \leq C\varepsilon^2 \sum_{k \neq j} \iint_{U \times \Omega} \frac{(\ln |\mathbf{x} - \mathbf{x}_k| - \ln |\mathbf{y} - \mathbf{x}_k|)^2}{|\mathbf{x} - \mathbf{y}|^2} d\mathbf{x} d\mathbf{y} \\ & = C\varepsilon^2 \sum_{k \neq j} \iint_{(U - \mathbf{x}_k) \times (\Omega - \mathbf{x}_k)} \frac{1}{|\mathbf{x}|^2} \frac{\ln^2 \left| \frac{\mathbf{y}}{\mathbf{x}} \right|}{\left(1 - \left| \frac{\mathbf{y}}{\mathbf{x}} \right| \right)^2} d\mathbf{y} d\mathbf{x} \\ & \leq C\varepsilon^2 \sum_{k \neq j} \int_{U - \mathbf{x}_k} \frac{1}{|\mathbf{x}|} \int_0^\infty \frac{\ln^2 z}{(1 - z)^2} dz d\mathbf{x}. \end{aligned}$$

The inner integral is convergent and hence we are left with

$$\varepsilon^2 \sum_{k \neq j} \int_{U - \mathbf{x}_k} \frac{1}{|\mathbf{x}|} d\mathbf{x} \leq C\varepsilon^2 \int_{e^{-2/\varepsilon}}^{e^{-1/\varepsilon}} z^{-1} dz = C\varepsilon.$$

When $2|\alpha| = d - 1$ and $d > 1$ the estimates for the difference quotient are a bit more technical. Similarly to above, Taylor's theorem combined with (B.1) yields

$$\begin{aligned} & \left| D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}') \right| = \left| \sum_{|\beta|=1} \int_0^1 D_{\mathbf{x}_j}^{\alpha+\beta} \chi_\varepsilon^*(\mathbf{x} + t(\mathbf{y} - \mathbf{x}); \mathbf{x}') (\mathbf{y} - \mathbf{x})^\beta dt \right| \\ & \leq C\varepsilon |\mathbf{x} - \mathbf{y}| \sum_{k \neq j} \int_0^1 \frac{\mathbb{1}_{B_{e^{-2/\varepsilon}}^c(\mathbf{x}_k)}(\mathbf{x} + t(\mathbf{y} - \mathbf{x}))}{|\mathbf{x} - t(\mathbf{y} - \mathbf{x}) - \mathbf{x}_k|^{|\alpha|+1}} dt. \end{aligned}$$

We estimate the integral

$$\int_0^1 |\mathbf{x} - t(\mathbf{y} - \mathbf{x}) - \mathbf{x}_k|^{-|\alpha|-1} dt.$$

Choosing coordinates in a plane containing \mathbf{x}_k, \mathbf{x} and \mathbf{y} such that $\mathbf{x}_k = (0, 0)$, $\mathbf{x} = (r_1, 0)$ and $\mathbf{y} = (r_2 \cos(\theta), r_2 \sin(\theta))$ with $\theta \in [0, \pi)$ we can write this integral as

$$\begin{aligned} & \int_0^1 \left(((1-t)r_1 - tr_2 \cos \theta)^2 + t^2 r_2^2 \sin^2 \theta \right)^{-\frac{|\alpha|+1}{2}} dt \\ &= \int_0^1 \left(\left(1 - \frac{tr_2}{(1-t)r_1} \cos \theta \right)^2 + \frac{t^2 r_2^2}{(1-t)^2 r_1^2} \sin^2 \theta \right)^{-\frac{|\alpha|+1}{2}} dt \\ &= \frac{1}{r_1 r_2^{|\alpha|}} \int_0^\infty \frac{(s + r_2/r_1)^{|\alpha|-1}}{((1+s \cos \theta)^2 + s^2 \sin^2 \theta)^{\frac{|\alpha|+1}{2}}} ds \\ &\leq \frac{1}{r_1 r_2^{|\alpha|}} \int_0^\infty \frac{(s+1)^{|\alpha|-1}}{((1-s)^2 + 2s(1+\cos \theta))^{\frac{|\alpha|+1}{2}}} ds \\ &=: \frac{g(\theta)}{r_1 r_2^{|\alpha|}} \end{aligned}$$

The integral $g(\theta)$ tends to infinity in the limit $\theta \rightarrow \pi$. However, this corresponds to \mathbf{x} and \mathbf{y} being far apart relative to their distance to the \mathbf{x}_k .

When θ is far from 0 we shall instead use the following bound which follows directly from the supremum bound in (B.1)

$$(B.3) \quad \left| D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}') \right| \leq C\varepsilon \sum_{k \neq j} \left[\frac{\mathbb{1}_{B_{e^{-2/\varepsilon}}(\mathbf{x}_k)}^c(\mathbf{x})}{|\mathbf{x} - \mathbf{x}_k|^{|\alpha|}} + \frac{\mathbb{1}_{B_{e^{-2/\varepsilon}}(\mathbf{x}_k)}^c(\mathbf{y})}{|\mathbf{y} - \mathbf{x}_k|^{|\alpha|}} \right]$$

together with the fact that

$$(B.4) \quad |\mathbf{x} - \mathbf{y}| \geq \sin(\theta/2) \max\{|\mathbf{x} - \mathbf{x}_k|, |\mathbf{y} - \mathbf{x}_k|\},$$

where θ is the angle between the vectors $\mathbf{y} - \mathbf{x}_k$ and $\mathbf{x} - \mathbf{x}_k$. Note that the bound in (B.3) does not capture the continuity of $D^\alpha \chi_\varepsilon^*$ and hence cannot be sufficiently accurate for our purposes when $|\mathbf{x} - \mathbf{y}|$ is small.

We are now ready to start estimating the H^s -seminorm of χ_ε^* . Using the same notation as in the $d = 1$ case

$$\begin{aligned} \|\chi_\varepsilon^*\|_{\dot{H}_{\mathbf{x}_j}^s(\Omega)}^2 &= \iint_{\Omega \times \Omega} \frac{|D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^{d+1}} d\mathbf{x} d\mathbf{y} \\ &\leq 2 \iint_{U \times \Omega} \frac{|D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^{d+1}} d\mathbf{x} d\mathbf{y}, \end{aligned}$$

where we used that $|D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{x}; \mathbf{x}')| = 0$ for $\mathbf{x} \in U_1 \cup U_2$, since $|\alpha| \geq 1$.

To bound the integral we use the estimates derived earlier. Recalling that in the case under consideration $|\alpha| = \frac{d-1}{2}$ the derived bounds tells us that

$$\begin{aligned} & \iint_{U \times \Omega} \frac{|D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{x}; \mathbf{x}') - D_{\mathbf{x}_j}^\alpha \chi_\varepsilon^*(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^{d+1}} d\mathbf{x} d\mathbf{y} \\ & \leq C\varepsilon^2 \sum_{k \neq j} \iint_{U \times \Omega} \frac{\min \left\{ \frac{g(\theta_k)^2}{|\mathbf{x} - \mathbf{x}_k|^{d-1} |\mathbf{y} - \mathbf{x}_k|^2}, \frac{g(\theta_k)^2}{|\mathbf{x} - \mathbf{x}_k|^2 |\mathbf{y} - \mathbf{x}_k|^{d-1}}, \frac{|\mathbf{x} - \mathbf{y}|^{-2}}{|\mathbf{x} - \mathbf{x}_k|^{d-1}} + \frac{|\mathbf{x} - \mathbf{y}|^{-2}}{|\mathbf{y} - \mathbf{x}_k|^{d-1}} \right\}}{|\mathbf{x} - \mathbf{y}|^{d-1}} d\mathbf{x} d\mathbf{y}, \end{aligned}$$

here θ_k denotes the angle between the vectors $\mathbf{x} - \mathbf{x}_k$ and $\mathbf{y} - \mathbf{x}_k$. For each fixed \mathbf{x} we rewrite the integral over Ω in spherical coordinates around \mathbf{x}_k , oriented so that \mathbf{x} is located at the south pole. With $R = |\mathbf{x} - \mathbf{x}_k|$, $r = |\mathbf{y} - \mathbf{x}_k|$ and θ_k as before, the integral becomes

$$\begin{aligned} & \iint_{U \times \Omega} \frac{\min \left\{ \frac{g(\theta_k)^2}{R^{d-1} r^2}, \frac{g(\theta_k)^2}{R^2 r^{d-1}}, \frac{|\mathbf{x} - \mathbf{y}|^{-2}}{R^{d-1}} + \frac{|\mathbf{x} - \mathbf{y}|^{-2}}{r^{d-1}} \right\}}{|\mathbf{x} - \mathbf{y}|^{d-1}} d\mathbf{x} d\mathbf{y} \\ & \leq \int_U \int_0^\infty \int_0^\pi \int_{\mathbb{S}^{d-2}} \frac{\min \left\{ \frac{g(\theta_k)^2}{R^{d-1} r^2}, \frac{g(\theta_k)^2}{R^2 r^{d-1}}, \frac{|\mathbf{x} - \mathbf{y}|^{-2}}{R^{d-1}} + \frac{|\mathbf{x} - \mathbf{y}|^{-2}}{r^{d-1}} \right\}}{((R - r \cos \theta_k)^2 + r^2 \sin^2 \theta_k)^{(d-1)/2}} r^{d-1} \sin^{d-2} \theta_k dr d\theta_k dS(\hat{\theta}) d\mathbf{x} \\ & = C \int_U \int_0^\infty \int_0^\pi \frac{\min \left\{ \frac{g(\theta_k)^2}{R^{d-1} r^2}, \frac{g(\theta_k)^2}{R^2 r^{d-1}}, \frac{|\mathbf{x} - \mathbf{y}|^{-2}}{R^{d-1}} + \frac{|\mathbf{x} - \mathbf{y}|^{-2}}{r^{d-1}} \right\}}{((R - r \cos \theta_k)^2 + r^2 \sin^2 \theta_k)^{(d-1)/2}} r^{d-1} \sin^{d-2} \theta_k dr d\theta_k d\mathbf{x}. \end{aligned}$$

For $\theta \in [\pi/2, \pi]$ we use the bounds in (B.3), (B.4):

$$\begin{aligned} & \int_U \int_0^\infty \int_{\pi/2}^\pi \frac{R^{-d+1} + r^{-d+1}}{\sin^{d+1}(\theta_k/2) \max\{R, r\}^{d+1}} r^{d-1} \sin^{d-2} \theta_k dr d\theta_k d\mathbf{x} \\ & = \int_U \int_R^\infty \int_{\pi/2}^\pi \frac{R^{-d+1} + r^{-d+1}}{\sin^{d+1}(\theta_k/2) r^2} \sin^{d-2}(\theta_k) dr d\theta_k d\mathbf{x} \\ & \quad + \int_U \int_0^R \int_{\pi/2}^\pi \frac{R^{-d+1} + r^{-d+1}}{\sin^{d+1}(\theta_k/2) R^{d+1}} r^{d-1} \sin^{d-2}(\theta_k) dr d\theta_k d\mathbf{x} \\ & \leq C \int_U \int_R^\infty \frac{R^{-d+1} + r^{-d+1}}{r^2} dr d\mathbf{x} + C \int_U \int_0^R \frac{R^{-d+1} + r^{-d+1}}{R^{d+1}} r^{d-1} dr d\mathbf{x} \\ & = C \int_U R^{-d} d\mathbf{x} \leq C \int_{e^{-2/\varepsilon}}^{e^{-1/\varepsilon}} R^{-1} dR = C\varepsilon^{-1}. \end{aligned}$$

Thus this part of the integral is $O(\varepsilon^{-1})$.

What remains is to bound the integral when $r \geq 0$ and $\theta_k \in [0, \pi/2)$. To accomplish this we shall use the bound for the difference of the derivatives derived earlier. Note that since $\theta_k < \pi/2$ we can replace the factor $g(\theta_k)$ by a constant without any loss. Using that $|\mathbf{x} - \mathbf{y}|^2 = R^2 + r^2 - 2rR \cos \theta \geq \max\{(R - r)^2, 2rR(1 - \cos \theta_k)\}$ we for

any fixed $\mu \in (0, 1)$ find

$$\begin{aligned}
& \int_U \int_0^\infty \int_0^{\pi/2} \frac{\min \left\{ \frac{1}{R^{d-1}r^2}, \frac{1}{R^2r^{d-1}} \right\}}{((R - r \cos \theta_k)^2 + r^2 \sin^2 \theta_k)^{(d-1)/2}} r^{d-1} \sin^{d-2} \theta_k dr d\theta_k d\mathbf{x} \\
&= \int_U \int_0^R \int_0^{\pi/2} ((R - r \cos \theta_k)^2 + r^2 \sin^2 \theta_k)^{-(d-1)/2} R^{-d+1} r^{d-3} \sin^{d-2} \theta_k dr d\theta_k d\mathbf{x} \\
&\quad + \int_U \int_R^\infty \int_0^{\pi/2} ((R - r \cos \theta_k)^2 + r^2 \sin^2 \theta_k)^{-(d-1)/2} R^{-2} \sin^{d-2} \theta_k dr d\theta_k d\mathbf{x} \\
&\leq \int_U \int_0^R (R - r)^{-\mu} r^{(d-5+\mu)/2} R^{-(3d-3-\mu)/2} dr d\mathbf{x} \int_0^{\pi/2} \frac{\sin^{d-2} \theta_k}{(1 - \cos \theta_k)^{(d-1-\mu)/2}} d\theta_k \\
&\quad + \int_U \int_R^\infty (r - R)^{-\mu} r^{-(d-1-\mu)/2} R^{-(d-3-\mu)/2} dr d\mathbf{x} \int_0^{\pi/2} \frac{\sin^{d-2} \theta_k}{(1 - \cos \theta_k)^{(d-1-\mu)/2}} d\theta_k \\
&\leq C \int_U (R^{-d} + R^{-d+3}) d\mathbf{x} \\
&\leq C \int_{e^{-2/\varepsilon}}^{e^{-1/\varepsilon}} (R^{-1} + R^2) dR = C\varepsilon^{-1}.
\end{aligned}$$

Consequently, also this part of the integral is $O(\varepsilon^{-1})$ which completes the proof. \square

LEMMA B.2. — For all $s > 0$ it holds that $\mathcal{H}_W^{s,N}(\mathbb{R}^d) \subseteq \mathcal{H}_0^{s,N}(\mathbb{R}^d)$.

Proof. — Take $\Psi \in H^s(\mathbb{R}^{dN})$ s.t. $\int W_s |\Psi|^2 < \infty$ and let $\Psi_\varepsilon := \chi_\varepsilon \Psi$. Since Ψ_ε is supported away from $\Delta_\varepsilon := \Delta + B_\varepsilon(0)$ and thus may be approximated in $C_c^\infty(\mathbb{R}^{dN} \setminus \Delta)$, it is sufficient to prove that $\|\Psi - \Psi_\varepsilon\|_{H^s(\mathbb{R}^{dN})} \rightarrow 0$ to conclude the lemma. We have by dominated convergence

$$\|\Psi - \Psi_\varepsilon\|_{L^2(\mathbb{R}^{dN})}^2 \lesssim \int_{\Delta_\varepsilon \cap \mathbb{R}^{dN}} |1 - \chi_\varepsilon|^2 |\Psi|^2 \rightarrow 0,$$

while for $\alpha \neq 0$

$$D_{\mathbf{x}_j}^\alpha ((1 - \chi_\varepsilon) \Psi) = \sum_{0 \leq \beta \leq \alpha} D_{\mathbf{x}_j}^\beta (1 - \chi_\varepsilon) D_{\mathbf{x}_j}^{\alpha-\beta} \Psi,$$

so for $s = m + \sigma$, $|\alpha| = m$, $0 \leq \sigma < 1$ (for $\sigma = 0$ we replace by L^2)

$$\begin{aligned}
\|\Psi - \Psi_\varepsilon\|_{\dot{H}^{s,N}(\mathbb{R}^d)} &\lesssim \sum_{j,\alpha} \left\| (1 - \chi_\varepsilon) D_{\mathbf{x}_j}^\alpha \Psi \right\|_{\dot{H}^{\sigma,N}(\mathbb{R}^d)} \\
&\quad + \sum_{j,\alpha} \sum_{0 < \beta < \alpha} \left\| (D_{\mathbf{x}_j}^\beta \chi_\varepsilon) (D_{\mathbf{x}_j}^{\alpha-\beta} \Psi) \right\|_{\dot{H}^{\sigma,N}(\mathbb{R}^d)} \\
&\quad + \sum_{j,\alpha} \left\| (D_{\mathbf{x}_j}^\alpha \chi_\varepsilon) \Psi \right\|_{\dot{H}^{\sigma,N}(\mathbb{R}^d)}.
\end{aligned}$$

We may estimate as in the proof of Lemma B.1,

$$\|(1 - \chi_\varepsilon) D^\alpha \Psi\|_{\dot{H}^{\sigma,N}(\mathbb{R}^d)}^2 \lesssim \sum_j \sum_{k \neq j} \int_{\mathbb{R}^{d(N-1)}} \|\Psi\|_{H_{\mathbf{x}_j}^s(B_\varepsilon(\mathbf{x}_k))}^2 \rightarrow 0,$$

$$\begin{aligned}
\|\Psi D^\alpha \chi_\varepsilon\|_{\dot{H}_{\mathbf{x}_j}^\sigma(\mathbb{R}^d)}^2 &= \iint_{\mathbb{R}^d \times \mathbb{R}^d} \frac{|\Psi D^\alpha \chi_\varepsilon(\mathbf{x}; \mathbf{x}') - \Psi D^\alpha \chi_\varepsilon(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^{d+2\sigma}} d\mathbf{x} d\mathbf{y} \\
&\lesssim I_\alpha + \varepsilon^{-2|\alpha|-2} \sum_{k,l \neq j} \iint_{B_{2\varepsilon}(\mathbf{x}_k) \times B_{2\varepsilon}(\mathbf{x}_l)} |\Psi(\mathbf{x}; \mathbf{x}')|^2 |\mathbf{x} - \mathbf{y}|^{-d-2\sigma+2} d\mathbf{x} d\mathbf{y} \\
&\lesssim I_\alpha + \varepsilon^{-2|\alpha|-2} \sum_{k \neq j} \int_{B_{2\varepsilon}(\mathbf{x}_k)} |\Psi(\mathbf{x}; \mathbf{x}')|^2 \int_{B_{4\varepsilon}(\mathbf{x})} |\mathbf{x} - \mathbf{y}|^{-d-2\sigma+2} d\mathbf{y} d\mathbf{x} \\
&\lesssim I_\alpha + \varepsilon^{-2|\alpha|-2\sigma} \sum_{k \neq j} \int_{B_{2\varepsilon}(\mathbf{x}_k)} |\Psi(\mathbf{x}; \mathbf{x}')|^2 d\mathbf{x},
\end{aligned}$$

where

$$I_\alpha = \sum_{k \neq j} \int_{B_{2\varepsilon}(\mathbf{x}_k)} |D^\alpha \chi_\varepsilon(\mathbf{x}; \mathbf{x}')|^2 \int_{B_{4\varepsilon}(\mathbf{x}_k)} \frac{|\Psi(\mathbf{x}; \mathbf{x}') - \Psi(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^{d+2\sigma}} d\mathbf{y} d\mathbf{x}.$$

For the highest-order derivatives $2|\alpha| = 2s - 2\sigma$:

$$\|\Psi D^\alpha \chi_\varepsilon\|_{\dot{H}^{\sigma,N}(\Omega)}^2 \lesssim \int_{\Omega^{N-1}} I_\alpha + \varepsilon^{-2s} \int_{\Delta_{2\varepsilon}} |\Psi(\mathbf{x})|^2 d\mathbf{x} \lesssim \int_{\Omega^{N-1}} I_\alpha + \int_{\Delta_{2\varepsilon}} W_s(\mathbf{x}) |\Psi(\mathbf{x})|^2 d\mathbf{x},$$

where the last term tends to zero as $\varepsilon \rightarrow 0$ by dominated convergence.

For I_α we have that

$$\begin{aligned}
I_\alpha &= \sum_{k \neq j} \int_{B_{2\varepsilon}(\mathbf{x}_k)} |D^\alpha \chi_\varepsilon(\mathbf{x}; \mathbf{x}')|^2 \int_{B_{4\varepsilon}(\mathbf{x}_k)} \frac{|\Psi(\mathbf{x}; \mathbf{x}') - \Psi(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^{d+2\sigma}} d\mathbf{y} d\mathbf{x} \\
&\lesssim \varepsilon^{-2|\alpha|} \sum_{k \neq j} \iint_{B_{4\varepsilon}(\mathbf{x}_k) \times B_{4\varepsilon}(\mathbf{x}_k)} \frac{|\Psi(\mathbf{x}; \mathbf{x}') - \Psi(\mathbf{y}; \mathbf{x}')|^2}{|\mathbf{x} - \mathbf{y}|^{d+2\sigma}} d\mathbf{y} d\mathbf{x} \\
&\lesssim \varepsilon^{-2|\alpha|} \sum_{k \neq j} \|\Psi\|_{\dot{H}_{\mathbf{x}_j}^\sigma(B_{4\varepsilon}(\mathbf{x}_k))}^2.
\end{aligned}$$

By interpolation of Sobolev spaces and scaling we have for $C = C(d, \sigma, m) > 0$

$$\|\Psi\|_{\dot{H}_{\mathbf{x}_j}^\sigma(B_{4\varepsilon}(\mathbf{x}_k))}^2 \leq C\varepsilon^{2m} \left(\|\Psi\|_{\dot{H}_{\mathbf{x}_j}^s(B_{4\varepsilon}(\mathbf{x}_k))}^2 + \left\| \Psi \sqrt{W_s} \right\|_{L_{\mathbf{x}_j}^2(B_{4\varepsilon}(\mathbf{x}_k))}^2 \right),$$

and thus by dominated convergence

$$\int_{\mathbb{R}^{d(N-1)}} I_\alpha \lesssim \sum_{k \neq j} \int_{\mathbb{R}^{d(N-1)}} \left(\|\Psi\|_{\dot{H}_{\mathbf{x}_j}^s(B_{4\varepsilon}(\mathbf{x}_k))}^2 + \left\| \Psi \sqrt{W_s} \right\|_{L_{\mathbf{x}_j}^2(B_{4\varepsilon}(\mathbf{x}_k))}^2 \right) \rightarrow 0,$$

for $|\alpha| = s - \sigma$. Similarly, for the lower-order mixed terms

$$\begin{aligned}
\|(D_{\mathbf{x}_j}^{\alpha-\beta} \Psi)(D_{\mathbf{x}_j}^\beta \chi_\varepsilon)\|_{\dot{H}_{\mathbf{x}_j}^\sigma(\mathbb{R}^d)} &\lesssim \varepsilon^{-2|\beta|} \sum_{k \neq j} \|D_{\mathbf{x}_j}^{\alpha-\beta} \Psi\|_{\dot{H}_{\mathbf{x}_j}^\sigma(B_{4\varepsilon}(\mathbf{x}_k))}^2 \\
&\quad + \varepsilon^{-2|\beta|-2\sigma} \sum_{k \neq j} \|D_{\mathbf{x}_j}^{\alpha-\beta} \Psi\|_{L_{\mathbf{x}_j}^2(B_{4\varepsilon}(\mathbf{x}_k))}^2 \\
&\lesssim \sum_{k \neq j} \left(\|\Psi\|_{\dot{H}_{\mathbf{x}_j}^s(B_{4\varepsilon}(\mathbf{x}_k))}^2 + \left\| \Psi \sqrt{W_s} \right\|_{L_{\mathbf{x}_j}^2(B_{4\varepsilon}(\mathbf{x}_k))}^2 \right),
\end{aligned}$$

which implies that also $\|(D^{\alpha-\beta} \Psi)(D^\beta \chi_\varepsilon)\|_{\dot{H}^{\sigma,N}(\mathbb{R}^d)} \rightarrow 0$. \square

LEMMA B.3. — For all $0 < 2s \leq d$ it holds that $\mathcal{H}_0^{s,N}(\mathbb{R}^d) = \mathcal{H}^{s,N}(\mathbb{R}^d) = H^s(\mathbb{R}^{dN})$.

Proof. — As mentioned above, combining Lemma B.2 with the Hardy–Rellich inequality implies the claim when $0 < 2s < d$. For $2s = d$ we argue as follows.

It suffices to prove that $C_c^\infty(\mathbb{R}^{dN} \setminus \Delta)$ is dense in H^s , and moreover, using that $C_c^\infty(\mathbb{R}^{dN})$ is dense in H^s , it suffices to prove that if $\Psi \in C_c^\infty(\mathbb{R}^{dN})$ then $\Psi_\varepsilon := \chi_\varepsilon^* \Psi \rightarrow \Psi$ in H^s as $\varepsilon \rightarrow 0$. Clearly

$$\|\Psi - \Psi_\varepsilon\|_{L^2(\mathbb{R}^{dN})}^2 \lesssim \int_{\Delta_\varepsilon \cap \text{supp } \Psi} |1 - \chi_\varepsilon^*|^2 \rightarrow 0.$$

Moreover, by Lemma B.1 and arguing as in the proof of Lemma B.2

$$\|\Psi - \Psi_\varepsilon\|_{H^{s,N}(\mathbb{R}^d)}^2 \lesssim \varepsilon \rightarrow 0. \quad \square$$

The above generalizes the case $d = 2$ and $s = 1$ where it is well known that hard-core bosons have non-extensive energy in the dilute limit [LY01] and thus that a Lieb–Thirring inequality of the type (2.1) cannot hold. See also [Sve81] for generalizations with integer s .

BIBLIOGRAPHY

- [Car61] Francis W. Carroll, *A polynomial in each variable separately is a polynomial*, Am. Math. Mon. **68** (1961), 42–44. ↑264
- [Dau83] Ingrid C. Daubechies, *An uncertainty principle for fermions with generalized kinetic energy*, Commun. Math. Phys. **90** (1983), no. 4, 511–520. ↑252
- [DL67] Freeman J. Dyson and Andrew Lenard, *Stability of matter. I*, J. Math. Phys. **8** (1967), no. 3, 423–434. ↑257
- [DNPV12] Eleonora Di Nezza, Giampiero Palatucci, and Enrico Valdinoci, *Hitchhiker’s guide to the fractional Sobolev spaces*, Bull. Sci. Math. **136** (2012), no. 5, 521–573. ↑267
- [Dys67] Freeman J. Dyson, *Ground-state energy of a finite system of charged particles*, J. Math. Phys. **8** (1967), no. 8, 1538–1545. ↑253
- [FHJN18] Rupert L. Frank, Dirk Hundertmark, Michael Jex, and Phan Thành Nam, *The Lieb–Thirring inequality revisited*, <https://arxiv.org/abs/1808.09017>, to appear in *Journal of the European Mathematical Society*, 2018. ↑252
- [FS12] Rupert L. Frank and Robert Seiringer, *Lieb–Thirring inequality for a model of particles with point interactions*, J. Math. Phys. **53** (2012), no. 9, article no. 095201 (11 pages). ↑256
- [Gir60] Marvin D. Girardeau, *Relationship between systems of impenetrable bosons and fermions in one dimension*, J. Mathematical Phys. **1** (1960), 516–523. ↑254
- [HSV13] Ritva Hurri-Syrjänen and Antti V. Vähäkangas, *On fractional Poincaré inequalities*, J. Anal. Math. **120** (2013), 85–104. ↑266
- [Lap12] Ari Laptev, *Spectral inequalities for Partial Differential Equations and their applications*, Fifth International Congress of Chinese Mathematicians. Part 1, 2, Studies in Advanced Mathematics, vol. 51, American Mathematical Society, 2012, pp. 629–643. ↑252
- [LL01] Elliott H. Lieb and Michael Loss, *Analysis*, second ed., Graduate Studies in Mathematics, vol. 14, American Mathematical Society, 2001. ↑265, 266
- [LL18] Simon Larson and Douglas Lundholm, *Exclusion bounds for extended anyons*, Arch. Ration. Mech. Anal. **227** (2018), no. 1, 309–365. ↑256

- [LNP16] Douglas Lundholm, Phan Thành Nam, and Fabian Portmann, *Fractional Hardy–Lieb–Thirring and related inequalities for interacting systems*, Arch. Ration. Mech. Anal. **219** (2016), no. 3, 1343–1382. ↑253, 254, 255, 256, 257
- [LPS15] Douglas Lundholm, Fabian Portmann, and Jan Philip Solovej, *Lieb–Thirring bounds for interacting Bose gases*, Commun. Math. Phys. **335** (2015), no. 2, 1019–1056. ↑255, 256
- [LS09] Elliott H. Lieb and Robert Seiringer, *The Stability of Matter in Quantum Mechanics*, Cambridge University Press, 2009. ↑252
- [LS13a] Douglas Lundholm and Jan Philip Solovej, *Hardy and Lieb–Thirring inequalities for anyons*, Commun. Math. Phys. **322** (2013), no. 3, 883–908. ↑256, 257
- [LS13b] ———, *Local exclusion principle for identical particles obeying intermediate and fractional statistics*, Phys. Rev. A **88** (2013), no. 6, article no. 062106 (9 pages). ↑256, 257
- [LS14] ———, *Local exclusion and Lieb–Thirring inequalities for intermediate and fractional statistics*, Ann. Henri Poincaré **15** (2014), no. 6, 1061–1107. ↑256
- [LS18] Douglas Lundholm and Robert Seiringer, *Fermionic behavior of ideal anyons*, Lett. Math. Phys. **108** (2018), no. 11, 2523–2541. ↑256, 258, 260
- [LT75] Elliott H. Lieb and Walter E. Thirring, *Bound for the kinetic energy of fermions which proves the stability of matter*, Phys. Rev. Lett. **35** (1975), no. 11, 687–689. ↑252, 255
- [LT76] ———, *Inequalities for the moments of the eigenvalues of the Schrödinger hamiltonian and their relation to Sobolev inequalities*, Studies in Mathematical Physics: Essays in Honor (Elliott H. Lieb, ed.), Princeton Series in Physics, Princeton University Press, 1976, pp. 269–303. ↑252, 254, 255
- [Lun18] Douglas Lundholm, *Methods of modern mathematical physics: Uncertainty and exclusion principles in quantum mechanics*, <https://arxiv.org/abs/1805.03063>, 2018. ↑256
- [LY01] Elliott H. Lieb and Jakob Yngvason, *The ground state energy of a dilute two-dimensional Bose gas*, J. Stat. Phys. **103** (2001), no. 3-4, 509–526. ↑280
- [Nam18] Phan Thành Nam, *Lieb–Thirring inequality with semiclassical constant and gradient error term*, J. Funct. Anal. **274** (2018), no. 6, 1739–1746. ↑256
- [Rum11] Michel Rumin, *Balanced distribution-energy inequalities and related entropy bounds*, Duke Math. J. **160** (2011), no. 3, 567–597. ↑255
- [Sve81] Edward C. Svendsen, *The effect of submanifolds upon essential self-adjointness and deficiency indices*, J. Math. Anal. Appl. **80** (1981), no. 2, 551–565. ↑280
- [Yaf99] Dimitri R. Yafaev, *Sharp Constants in the Hardy–Rellich Inequalities*, J. Funct. Anal. **168** (1999), no. 1, 121–144. ↑273

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